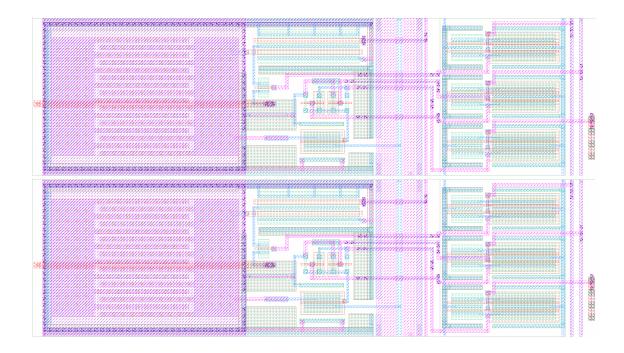
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Fast sampling chip design

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University of Chicago, September 3, 2009

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1 Introduction

The goal of the P-sec timing project is to develop cheap large-area photo-detectors that provide excellent space and time resolution. The key improvement would be a flat structure consisting in an assembly of simple sandwich of layers rather than of discrete parts. Because of its integration, such a structure would allow to measure the time of arrival of relativistic particles with (ultimately) 1 Pico-second resolution. The potential applications would be precise time-of-flight measurement for particle accelerators; Positron-Emission Tomography; large area detectors; and non-proliferation security. The target of the project is to address the main issues of such a new structure, and delivering a working prototype within three year. This means building, assembling and packaging a detector made of all the layers: photo-cathode, micro-channels, and transmission lines anodes; but also developing the readout electronics. [1]

2 Project description

The objective of the Pico-second Timing project is to develop a new family of large-area photo-detectors.

Advances in material science and nano-technologies, along with the recent innovations in microelectronics and data processing, give us the opportunity to apply the basic concept of micro-channel plates to the development of photo-detectors. These micro-channel plates photomultipliers are an evolution of photomultipliers tubes (PMT's).

In the current state of art, PMT's are high-bandwidth, high-gain, low noise with a high quantum efficiency. But, because MCP-PMT's have a much smaller path-length for the photons to electrons amplification, there resulting response times are considerably higher, giving a much better intrinsic time resolution along with all the benefits of standard PMT's.

Also, as MCP-PMT's are an assembly of simple layers, they can naturally cover large area at low cost and provide good space resolution.

Therefore, these detectors could be tailored for a wide variety of applications where photon detection with excellent time and space resolution is required, and will be revolutionary : particle detectors (LHC, RHIC, JPARC, Super-B and the ILC); Positron-Emission Tomography; large area detectors; and non-proliferation security.

3 Detectors in High Energy Physics

3.1 Particle detector

In experimental and applied particle physics and nuclear engineering, a particle detector, is a device used to detect, track, and/or identify highenergy particles, such as those produced by nuclear decay, cosmic radiation, or reactions in a particle accelerator.

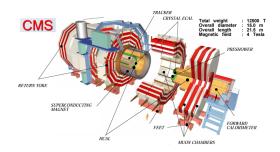


Figure 3.1: The Compact Muon Solenoid (CMS) is an example of a large particle detector.

Modern detectors, combine very often several layers, assembled like an onion, in order to identify and to measure the position and energy of the detected particles.

3.2 High Energy Physics particle measurement

A particle has three related characteristics: mass, velocity, and momentum. A measurement of any two will yield the remaining one. If we can measure the velocity β and momentum p of a particle, we can find its mass m via $p = \gamma m \beta$.

The measurement of the momentum is made by applying a strong and uniform magnetic field parallel to the opposing beam. After the collision, the transverse momentum of each collision product is given by $p_T = qBr$, where qis the charge of the particle, B is the magnitude of the magnetic field, and ris the radius of curvature of its path.

While the measurement of momentum is fairly standard, there are three different measurements that can lead to β : the energy loss in a dense medium, the angle of Čerenkov radiation and time-of-flight.

3.2.1 Energy loss measurement

This technique aims to measure the particle's energy loss in a dense stopping medium, by ionization. This detector is called a calorimeter in the context of particle physics. Most particles enter the calorimeter and initiate a particle shower. The particles' energy is deposited in the calorimeter, collected, and measured (Fig 3.2). In order to achieve a good resolution in energy, calorimeters usually require a significant radial space (Fig 3.1).

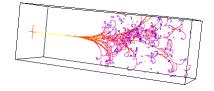


Figure 3.2: Particle shower in calorimeter

3.2.2 Cerenkov angle measurement

The second technique consist in the observation of Čerenkov radiation. When moving trough an optically transparent medium of refractive index n, a charged particle emits radiation if its velocity v is greater than the local speed of light: $\frac{c}{n}$.

The radiation is emitted at an angle θ so that: $\cos(\theta) = \frac{1}{n\beta}$. β can then be measured by placing a plane of photo-detectors somewhere down-path of the refractive medium, resolving the resulting conic section, and calculating θ [2]. This technique requires also a significant radial space to achieve a good

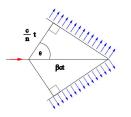


Figure 3.3: Čerenkov light cone

resolution in β .

3.2.3 Time-of flight measurement

In this technique, instead of measuring β directly, one determines the time at which a particle arrives at distance from the interaction point. Instead of fixing t_0 , one uses the reconstructed tracks of the particles from a single vertex to infer a time for the interaction. Each track can be used individually to determine the path length L. From these inputs, one calculates β .

This technique offers several distinct advantages. Unlike a measurement of energy loss in a stopping medium, the resolution of this technique is not limited by the radial space available, but depends on the time resolution of the sensors and readout electronics used. The same consideration allows a time-of-flight detector to be implemented in much less space than a pure Čerenkov detector of equivalent resolution (Fig 3.4).

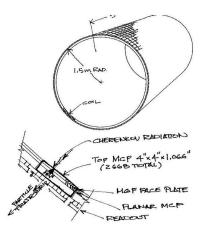


Figure 3.4: Time-of-flight detector structure

3.3 Photo-multiplier tubes

Photomultipliers are constructed from a glass envelope with a high vacuum inside, which houses a photocathode, several dynodes, and an anode (Fig 3.5). Incident photons strike the photocathode material, which is present as a thin deposit on the entry window of the device, with electrons being produced as a consequence of the photoelectric effect. These electrons are directed by the focusing electrode toward the electron multiplier, where electrons are multiplied by the process of secondary emission [3].

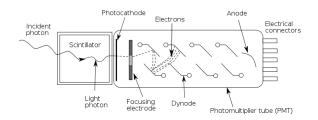


Figure 3.5: Schematic of a photomultiplier tube coupled to a scintillator

The electron multiplier consists of a number of electrodes called dynodes. Each dynode is held at a more positive voltage than the previous one. The electrons leave the photocathode, having the energy of the incoming photon (minus the work function of the photocathode). As the electrons move toward the first dynode, they are accelerated by the electric field and arrive with much greater energy. Upon striking the first dynode, more low energy electrons are emitted, and these electrons in turn are accelerated toward the second dynode. The geometry of the dynode chain is such that a cascade occurs with an ever-increasing number of electrons being produced at each stage. Finally, the electrons reach the anode, where the accumulation of charge results in a sharp current pulse indicating the arrival of a photon at the photocathode [4].

3.4 Micro-channel plates time-of-flight detector

3.4.1 Detector sandwich architecture

Micro-channel plates, as it can be seen in Figure 3.6 are made of four different parts: input window, photo-cathode , micro-channel plate and an anode array.

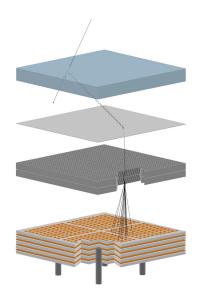


Figure 3.6: Fast Time-of-flight detector

3.4.2 Input window

The input window is the surface by which the photon will enter the detector. It must therefore be transparent to the radiations to be measured. It has also to be sufficiently resistant, as the detector is in vacuum and it will have to hold the atmospheric pressure. For detection of simple glass can be used.

Furthermore, this window can also be used to detect charged particle using the Čerenkov effect. In this case, a relativistic particle will generate a Čerenkov cone of light in the window medium, which will be seen by the detector.

3.4.3 Photo-cathode

Photoelectric effect

When a negatively charged metallic surface is exposed to electromagnetic radiation above a certain threshold wavelength (typically visible light), the light is absorbed and electrons are emitted. The emitted electrons can be referred to as photoelectrons in this context [5]. Although a plain metallic cathode will exhibit photoelectric properties, a custom coating greatly increases the efficiency. A photocathode usually consists of alkali metals with very low work functions.

Unfortunately, as it can be seen in Figure 3.7(b), the quantum efficiency of this conversion is not 100% and varies with the material used. Which means that only a fraction of the incident photons will be converted to electrons and therefore observed. For a good detector, this quantum efficiency must the highest possible in the spectrum range.

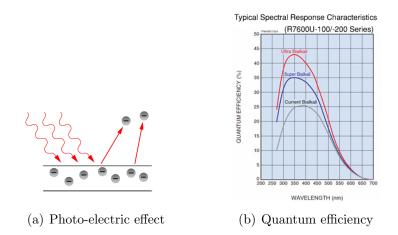


Figure 3.7: (a) Photo-electron production (b) Typical quantum efficiency in our design

Advanced photo-cathodes

In order to improve the quantum efficiency of the photo-cathode, some complex structure can be used (Fig 3.8) [6].

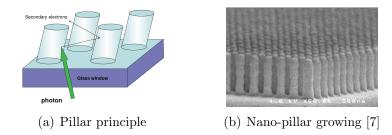


Figure 3.8: (a) Pillar photo-cathode principle (b) Nanoscale realization

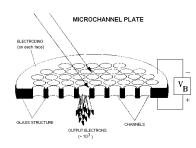
3.4.4 Micro-channel plate

Electron amplification

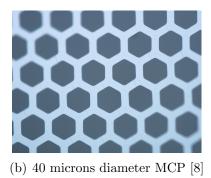
The signal coming from the photo-cathode is usually too small to be easily processed by the read-out electronics. An amplification is therefore required. To amplify the produced electrons, one can use high-voltage biased pores.

Micro-channel plate

A micro-channel plate is a plate made of an assembly of tiny tubes or pores (microchannels) leading from one face to the opposite, densely distributed over the whole surface. The micro-channels are typically 10 micrometers in diameter.



(a) Micro-channel plate principle



Micro-channel principle

Each micro-channel is a continuous-dynode electron multiplier, in which the multiplication takes place under the presence of a strong electric field. A particle or photon that enters one of the channels through a small

Figure 3.9: (a) Micro-channel plate sketch (b) Microscopic view

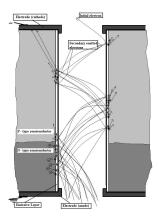


Figure 3.10: Electron amplification in micro-channel plates

orifice is guaranteed to hit the wall of the channel due to the channel being at an angle to the plate and thus the angle of impact. The impact starts a cascade of electrons that propagates through the channel, which amplifies the original signal by several orders of magnitude depending on the electric field strength and the geometry of the micro-channel plate (Fig 3.10) [9].

3.4.5 Anode plate

At the end of the micro-channel plate, the electrical signal is collected on an anode array (Fig 3.4). The dimensions of the array give the space resolution. The more pixels on the anode, the more precisely the location of an event on the detector is determined.

But, with this structure, they are as many output channels as the number of anode pads. That is to say that if there are n by n pads, there will be n^2 channels for the front-end electronics.

A way to reduce this number is to use delay lines (Fig 3.11): pads are connected row by row to a delay line. Whenever a pad receive a pulse coming from the MCP, it will be transmitted to the delay line and propagate to both ends of it. The signal is then collected and digitized at each end [10].

This design reduces by a factor n the number of output channels. Furthermore, the interpolated location of a pulse strike on the line can be more accurate than the size of the pads (this cannot be achieved with the array structure only).

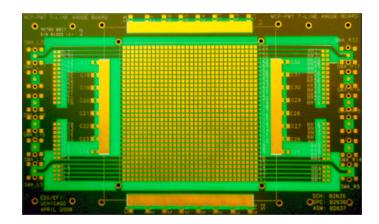


Figure 3.11: Delay lines used to collect the electrical signal

3.4.6 Full sandwich structure of the detector

The assembly of this detector is showed Figure 3.12(a), and the principle sketched on Figure 3.12(b).

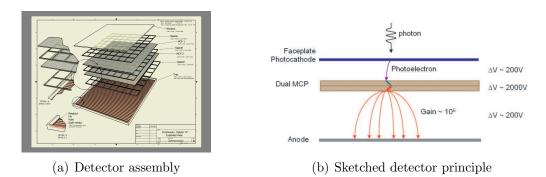


Figure 3.12: (a) Detector assembly sketch (b) Photon to electric signal conversion

4 Signal processing for Pico-second resolution

4.1 Timing techniques

Present photo-detectors such as micro-channel plate photo-multipliers (MCP-PMTs) and silicon photo-multipliers achieve rise-times well below one nano-second. An ideal timing readout electronics would extract the time-ofarrival of the first charge collected, adding nothing to the intrinsic detector resolution. Traditionally the best ultimate performance in terms of timing resolution has been obtained using constant fraction discriminators (CFDs) followed by high precision time digitization. However, these discriminators make use of wide-band delay lines that cannot be integrated easily into silicon integrated circuits, and so large front-end readout systems using CFD's to achieve sub-nsec resolution have are not yet been implemented. Several other well-known techniques in addition to constant-fraction discrimination have long been used for timing extraction of the time-of-arrival of a pulse [11]:

- Single threshold on the leading edge.
- Multiple thresholds on the leading edge, followed by a fit to the edge shape.
- Pulse waveform sampling, digitization and pulse reconstruction.

4.2 Single threshold

Single threshold discriminator is the simplest structure that gives the time-of-arrival of the first charge. The principle, very simple, is described Figure 4.1.

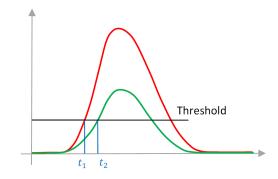


Figure 4.1: Single threshold timing resolution

We can see on Figure 4.1, that the time-of-arrival with this technique is strongly variable if the rise time of the signal is amplitude-dependent. Also, for applications in which one is searching for rare events with anomalous times, the single measured time does not give indications of possible anomalous pulse shapes due to intermittent noise, rare environmental artifacts, and other real but rare annoyances common in real experiments.

4.3 Multiple threshold

The multiple threshold technique samples the leading edge at amplitudes set to several values, for instance at values equally spaced between a minimum and a maximum threshold. The leading edge is then reconstructed from a fit to the times the pulse reaches the thresholds to extract a single time as characteristic of the pulse (Fig 4.2).

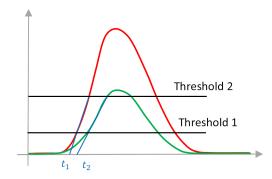
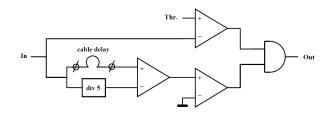


Figure 4.2: Multiple threshold timing resolution

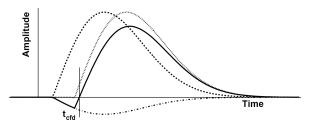
This technique provide that, if the pulse is linearly amplitudedependent, the reconstructed time will not depend on it.

4.4 Constant fraction discriminator

The most commonly used constant fraction discriminator technique forms the difference between attenuated and delayed versions of the original signal, followed by the detection of the zero crossing of the difference signal (Fig 4.3).



(a) Functional diagram of a constant fraction discriminator



(b) Operation of the CFD

Figure 4.3: (b) The input pulse (dashed curve) is delayed (dotted) and added to an attenuated inverted pulse (dash-dot) yielding a bipolar pulse (solid curve). The output of the CFD fires when the bipolar pulse changes polarity which is indicated by time t_{cfd} .

The timing resolution improvement obtained by a constant fraction discriminator is shown Figure 4.4. This technique which takes into account the amplitude dependency of the rising edge, without depending on the pulse shape usually gives very good results. There are three tunable parameters: the delay, the attenuation ratio, and the threshold. These parameters have to be carefully set with respect to the pulse characteristics in order to obtain the best timing resolution.

4.5 Pulse waveform sampling

Waveform sampling stores successive values of the pulse waveform. For precision time-of-arrival measurements, such as considered here, one needs

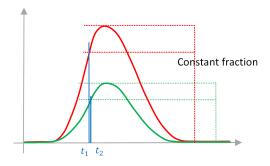


Figure 4.4: Constant fraction timing

to fully sample at least the leading edge over the peak. In order to fulfill the Shannon-Nyquist condition, the sampling period has to be chosen short enough to take into account all frequency components containing timing information, which is that the minimum sampling frequency is set at least at twice the highest frequency in the signal's Fourier spectrum.

After digitization, using the knowledge of the average waveform, pulse reconstruction allows reconstructing the edge or the full pulse with good fidelity. The sampling method is unique among the four methods in providing the pulse amplitude, the integrated charge, and figures of merit on the pulseshape and baseline, important for detecting pile-up or spurious pulses.

4.6 Simulation results of comparison between the four techniques

The four timing techniques presented below have been simulated using Matlab [12]. At the input, the signal used is a typical MCP output signal. To this signal is superimposed white shot noise from the MCP and white thermal noise coming from the electronics (Fig 4.5). For all the techniques, the input bandwith is taken to be 1.5 GHz. For the pulse sampling technique, a sampling rate of 20 GSa/s is used.

The time-of-arrival is determined, for all the timing techniques described before and plotted versus the number of photo-electrons (strength of the signal) Figure 4.6. We can clearly see that the sampling technique is the one providing the best timing resolution out out the four other.

This simulation also shows that the Pico-second precision range is achievable with MCP's signals. Coming from this statement a very fast sam-

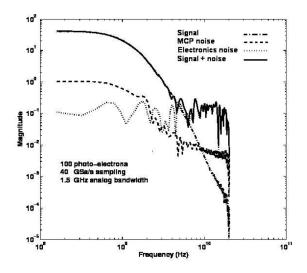


Figure 4.5: Fourier spectra of the noisy signal used in simulation. [11]

pling chip has been designed in order to achieve this Pico-second timing precision.

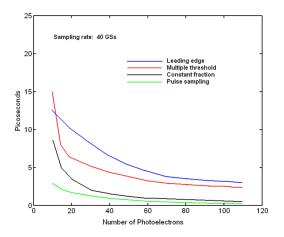


Figure 4.6: Time resolution versus the number of primary photo-electrons. [11]

5 Detector to Chip integration

5.1 Strip lines

As we discuss previously the bottom of the detector are strip lines that are running underneath it, as it can be seen on Figure 5.1. These strips lines are matched to be 50 Ohms. From the output of these lines to the input of the chip, the distortion of the signal must be as small as possible. Futhermore, at the input of the chip, the input structure is coplanar. Putting everything together, we end up with the following specifications:

- 50 Ohms matching.
- Microstrip to coplanar transition.
- Minimizing attenuation of the signal.
- Simple assembly design.

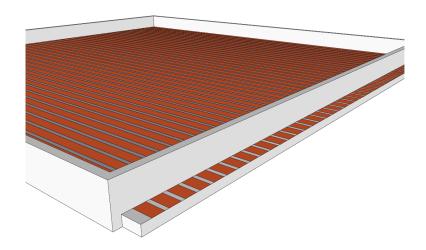


Figure 5.1: Strip line design

5.1.1 Strip line models and impedance

Microstrips lines consists of a conducting strip separated from a ground plane by a dielectric layer (Fig 5.2).

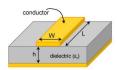


Figure 5.2: Microstrip line principle

Impedance

The impedance of a single strip line is given by the formula [13]:

$$Z_{0m} = \frac{377}{2\pi \left\{ \left(\epsilon_r + 1\right)/2 \right\}^{1/2}} \left[ln \left(\frac{8h}{W}\right) + \frac{1}{8} \left(\frac{W}{2h}\right)^2 - \frac{1}{2} \frac{\epsilon_r - 1}{\epsilon_r + 1} \left\{ ln \frac{\pi}{2} + \frac{1}{\epsilon_r} ln \frac{4}{\pi} \right\} \right]$$

If we choose one type of substrate (here glass) with a specified thickness, the only remaining parameter is the width W of the lines. In our case, a width of 3.75mm has been calculated with a high h of 2mm. The glass used is Borosilicate, with an ϵ_r of 4.

Lines losses

There is four different sources for the losses:

- Losses due to the resistance of the line
- Losses due to the dielectric conductivity
- Losses due to the dipole rotation $tan(\delta)$
- The radiative losses

The more important factors appears to be the losses due to the resistance of the line and the losses due to dipole rotation.

Bandwidth

The bandwidth depends mainly upon the dielectric constant and of the geometry of the strip line, as it can be seen on the following formula [14]:

$$BW_{TL} = \frac{3dB}{2.3 \times tan(\delta) \times \sqrt{\epsilon_r}} \times \frac{1}{d} = \frac{1.3}{tan(\delta) \times \sqrt{\epsilon_r}} \times \frac{1}{d}$$

5.2 Chip

At the input of the chip, all the inputs (Signals and GND's) are at the same metal level. A microstrip line at the input is therefore impossible. A coplanar line will therefore be used at the input of the chip.

Also, the chip can not simply sit on the glass because of the high number of I/O pads (144). Chips are usually connected using PCB (Printed Circuit Board). These boards have different physical and electrical characteristics (h, ϵ_r and metalizations layers), the connection between the glass and the PCB is therefore not straightforward as we want to avoid 50 Ohms breaks throughout the line.

On the PCB, one will also eventually go from microstrip to coplanar lines. The easiest way to do it as presented on Figure 5.3, but it has to be tested previously.

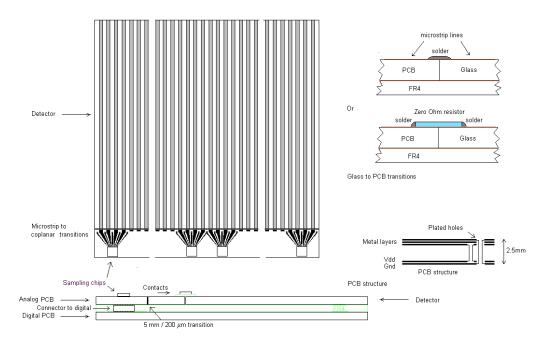


Figure 5.3: Detector & chip assembly

6 Chip structure and characteristics

In order to get the smallest spread in time-of-arrival of the photon at the photocathode, we have, according to the previous study, to acquire the signal coming out the microstrip lines with a rate as high as 40Gs/s, and keeping an analog bandwidth at the input higher than 2GHz. For that we have designed a new integrated sampling circuit using a 130nm CMOS process, that we will present now.

6.1 Fast read-out electronics

As we have seen previously, the required range for the signal acquisition should be higher that 10Gs/s. At this rate, a straight digitization is for now impossible. Therefore, simply having an ADC at the input of the line not doable. The solution is to sample the analog signal at a very high frequency and then digitize it at a slower rate using ADC's. This is achieved using the following structure.

6.2 Structure

6.2.1 General architecture and mechanism

Top design

In order to store the analog value of the signal the structure used is a switched capacitor array. The circuit principle is very simple and mainly consist of four different structures presented on Figure 6.1

- A timing generator
- An array of sampling cells
- An array of ADC's
- A token controlled register readout system

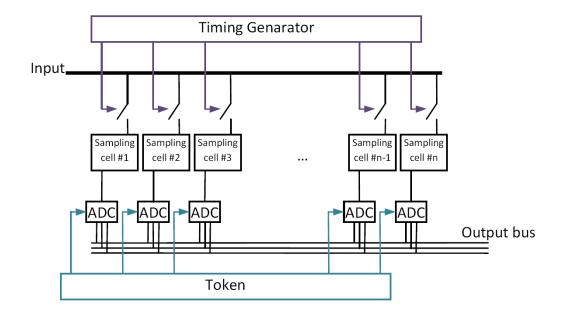


Figure 6.1: Circuit general sketch

Mechanism

The signal is coming from the delay lines and is put at the input of the chip with a 50 Ohms matching.

This signal is sampled at t_0 in the first sampling cell, then $t_0 + \Delta_{sampling}$ in the second cell, and so on... The sampling frequency is therefore given by $\frac{1}{\Delta_{sampling}}$. When all the cells have been written, the signal overwrite the first cells for a new cycle. So, if we use *n* cells, the signal can only be recorded during the last $n \times \Delta_{sampling}$ seconds.

6.2.2 Timing generator

The timing generator is the structure providing control signals that will allow to write in the sampling cells. More precisely, it provides a writing window to every storage cell. During this window, the input switch of the storage cell is closed and the signal stored in it.

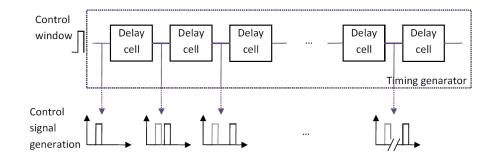


Figure 6.2: Timing generator principle

The timing generator, as it can be seen on Figure 6.2, is essentially made of a chain of delay cells. The sampling window is sent at the input of the timing generator and transmitted trough every delay cells. Each delay cell takes Δ_{delay} to pass it to the following one. If we have *n* delay cells in the timing generator, we have then at the output *n* sampling window, each of which is delay by Δ_{delay} compared to the next one.

The sampling frequency achieved by this way is therefore $f_{sampl} = \frac{1}{\Delta_{delay}}$. So the smaller is the delay, the higher is the sampling frequency, and in our case, the better the timing resolution.

6.2.3 Sampling cells

The sampling cell is the structure where the signal is stored. The sampling cells, being controlled by the timing generator are storing, one after another the input signal, as it can be seen on Figure 6.3.

Once the control signal has closed the switch at the input of one cell, this one is going to reach the input signal value as long as the time window given by the timing generator allows and holds it. The time constant of this process determines the bandwidth of the sampling cell.

The storage cell must have a very high bandwidth at the input, in order to be able to follow the rising edges of the fast incoming pulses. Once the signal is stored it should remains as stable as possible until it has been digitized by the ADC. Indeed, leakages are responsible for droops before the ADC has digitized the cell voltage.

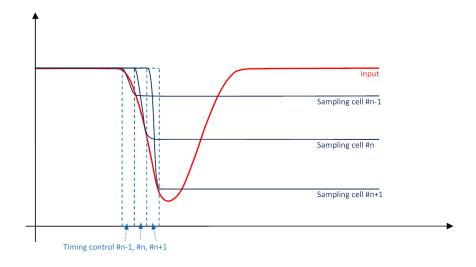


Figure 6.3: Basic sampling principle

The sampling cell will be described in more details in the following parts of this report.

6.2.4 ADC's

The ADC's we choose to use in our design are ramp ADC's. This structure has been choose because of its very good parallelism integration. And with our design of one ADC per storage cell, this structure is therefore the more suitable.

The ADC's are working with the following principle (described on Fig 6.5). A sketch of one ADC can be seen on Figure 6.4.

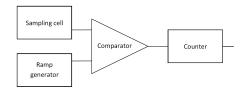


Figure 6.4: Sketch of the ramp ADC

At the beginning of the digitization process, the sampled value is available at the output of the sampling cell. When the digitization process is started, a voltage ramp and the sampling cell stored value are put at the input of a comparator. In the meantime, the counter starts counting. When the ramp value reach the sampling cell value, the comparator fires and stops the counter.

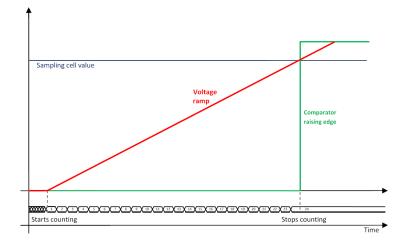


Figure 6.5: Ramp ADC principle

The digital value available at the output of the counter is the digitized value of the sampling cell.

6.2.5 Token controlled register readout system

There is only one output bus to read the data at the output of the chip. As it can be seen on Figure 6.1, all the ADC's are connected to this bus. To avoid conflict during the readout phase, the output of the ADC's are serially put on the output bus using a token passing circuit. The role of token controlled register readout system is therefore to put successively the output of the ADC's on the bus one at a time.

6.3 Specification

Coming from the simulation and the structure, the following specification have been set for the sampling chip:

	Chip characteristics
Sampling rate	10GS/s
Analog Bandwidth	2GHz
Dynamic range	0.7V
Sampling window	adjustable 500ps-2ns
Sampling jitter	10ps
Maximum latency	TBD
Crosstalk	1%
DC Input impedance	50Ω internal
Conversion clock	Adjustable 1-2 GHz internal ring oscillator.
	Minimum conversion time 2us.
Read clock	40 MHz. Readout time (4-channel)
	$4 \times 256 \times 25ns = 25.6\mu s$
Power	40mW/channel
Power supply	1.2V
Process	IBM 8RF-DM (130nm CMOS)

7 Operation of the chip

The purpose of this chip is to record very fast pulses from the MCP detector. The operation is very simple and it has mainly two states. The first one is called the write state, in which the chip is recording the signals coming at the input as described. The second one is call the read state and is when the user is reading the data stored inside the chip.

7.1 Writing

7.1.1 Sampling cell writing

The writing phase is started by launching the timing generator. When started, this one will controlled the successive closing of the input switches of every storage cell of one channel. As we previously described, the input signal is stored, at t_0 , in the first cell, $t_0 + \Delta_{sampling}$, in the second one and so on, until the last one. Once the last cell has been written, the signal overwrite the first cell for a new cycle.

As we also said previously, the storage cells record the last $n \times \Delta_{sampling}$ seconds only. Basically, the cells are then constantly overwritten until a trigger event.

7.1.2 Trigger event

When an event occurs at the detector (pulse for example), a trigger bit is raised. The effect of the trigger signal is the following: when at 0, the signal is let to be written is each cell successively at it has been described. When the trigger raise to 1, it overrides all the timing control signal and open all of the storage cells' input switches. At this time we have then the shape of the signal during the last $n \times \Delta_{sampling}$ seconds stored in the cells.

The trigger signal also start the digitization process for each cell. At the end, we then have the digitized value of the signal during the last $n \times \Delta_{sampling}$ seconds before the trigger available.

7.2 Reading

The reading phase starts after a trigger event, when the signal has been digitized and its values are available from the chip. During the reading phase, the token controlled readout register serially assert every stored value on the output bus. These values are then read externally, using a FPGA for example.

8 Storage cells

8.1 Storage cell principle

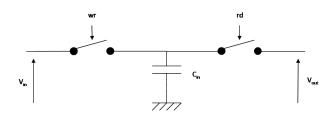


Figure 8.1: Elementary sampling cell

The most simplest sampling cell, is only a capacitor, in which the input voltage will be stored . There are three steps to store the data :

Write state

The first one is the write state, when the write switch is closed and read switch open (Fig 8.2). During this state, the capacitance is being charged by the input voltage V_{in} .

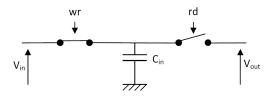


Figure 8.2: Write state

Intermediate state

During this state, both read and write switches are open (Fig 8.3), the value stored should not evolve. We will see later that it is usually not the case due to leakages.

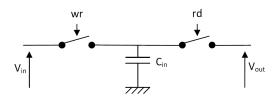


Figure 8.3: Intermediate state

Read state

The third one is the read state, when the write switch is open and the read switch is close (Fig 8.4). The output voltage is then given by the voltage stored by the capacitance during the write state.

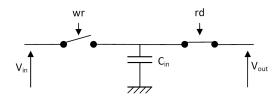


Figure 8.4: Read state

This is the structure always employed to sample an analog signal. The principle is simple, but they are several issues that must be taken care of when designing. The main ones are:

- Write state bandwidth
- Charge leakage
- Charge injection

8.1.1 Write state bandwidth

Basically, the write state consist of charging the input capacitance. When closed the write switch resistance is not zero and can be relatively high (typically around $10k\Omega$). Therefore, one must take care of the cutoff frequency of this RC circuit.

8.1.2 Charge leakage

Between the read and write state, when both read and write switch are open, the charge stored on the input capacitance may evolve due to several effect. This effect will lead to reading an output that could be perceptibly different of the input. The predominant contributions of such a leakage are the following :

Capacitor leakage

When charged, the insulating layer between the two plates of the transistor, theoretically prevent the charge going from one plate to the other. Practically, there is always a leakage current between the two plates discharging slowly the capacitor. This leakage is small in our case.

Switch leakage

When off, the equivalent resistance of the switch is ideally infinite, in reality this resistance has a high value, but not infinite. Therefore, there is a leakage current trough the off switch too.

8.1.3 Charge injection

The charge injection is the small charge transferred to the storage capacitor via the interelectrode capacitance of the switch and the stray capacitance when switching to the hold mode. The offset step is directly proportional to this charge:

Offset error = Incremental Charge/Capacitance = $\frac{\Delta Q}{C}$

It can be reduced somewhat by lightly coupling an appropriate polarity version of the hold signal to the capacitor for a first-order cancellation. The error can also be reduced by increasing the capacitance, but this increases acquisition time and decreases the bandwidth. The charge injection is also amplitude dependent, which makes a total cancellation hard to realize.

8.1.4 Output current level

In the architecture shown Figure 8.1, one can see that we use on the output the same charge as it has been stored at the input. This is not the smartest design as we will be confronted to all the previous factors of error when moving this charge to the output. Therefore, we have to think about an architecture that prevents the charge to flow again during the read phase, responsible for voltage droops.

8.2 Input switch

The input switch is the first stage of our circuit and must therefore be carefully studied in order to preserve the signal transfer to the storage capacitor.

8.2.1 Design

The design of the switch is shown on Figure 8.5.

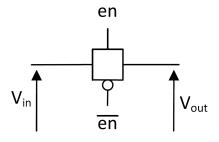


Figure 8.5: Switch design

8.2.2 Large signal analysis

The switch consist of the association of two FET's: one Nfet and one Pfet in parallel controlled by the "en" tension. Indeed, two FET's are necessary is one wants to switch in the full range given by the power supplies. The two important parameters are the R_{ON} and R_{OFF} resistances of the switch:

Switch R_{ON}

As the switch consist in the association of one Pfet and one Nfet in parallel, the R_{ON} will be given by the R_{ON} resistance of each fet.

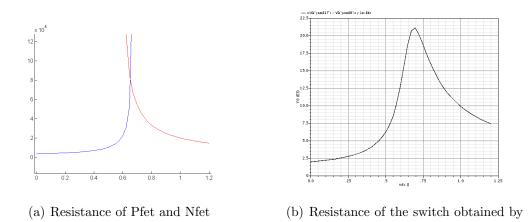
In the ON state of the switch, the voltage across the switch should be small and v_{GS} should be large. Therefore, the fet is assumed to be in the nonsaturation region. And we have:

$$i_D = \frac{KW}{L} \left[\left(v_{GS} - V_T \right) v_{DS} - \frac{v_{DS}^2}{2} \right]$$

The r_{ON} resistance of the switch is then given by:

$$r_{ON} = \frac{1}{\partial i_D / \partial v_{DS}} = \frac{L}{KW \left(V_{GS} - V_T - V_T \right)}$$

This resistance is plotted Figure 8.6(a) for both Pfet (red) and Nfet (blue). The resistance become almost infinite for $V_{GS} < V_T$ for both transistors.



a Spice simulation

The switch resistance is the parallel addition of the resistances of the Pfet and the Nfet, it has been simulated with Spice, and is plotted Figure 8.6(b).

On the Figure 8.6, is plotted the experimental resistance obtained with Spice and the equivalent resistance obtained with Matlab. We can see that the experimental curve fits very well the model of the resistance of a Pfet and an Nfet in parallel.

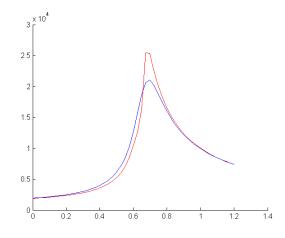


Figure 8.6: Resistance of the switch given by Spice & Matlab

8.2.3 Bandwidth

From the small signal model of the switch given Figure 8.7, we can find the following cutoff frequency induced by the switch:

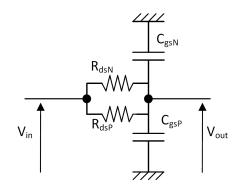


Figure 8.7: Small signal model of the switch

 $V_{out} = V_{in} \left(1 + j (R_{dsP} / / R_{dsN}) (C_{gsP} + C_{gsN}) \omega \right)$

The cutoff frequency of the switch is therefore given by:

$$f_c = \frac{2\pi}{(R_{dsP}//R_{dsN})(C_{gsP} + C_{gsN}))}$$

As the resistance R_{ds} , decreases when the length L of the transistor increases, and the capacitance C_{gs} increases when the area $W \times L$ of the transistor increases, the input bandwidth, therefore cannot exceed a maximum value.

The switch bandwidth obtained after simulation is plotted Figure 8.8. The maximum cutoff frequency is 13GHz at 3dB. Getting a bandwidth higher than 13GHz is therefore not possible for our circuit.

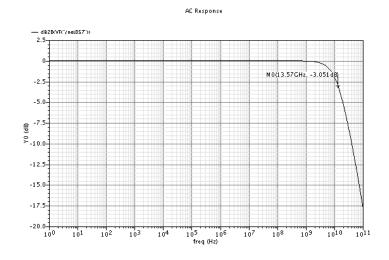


Figure 8.8: Simulated switch bandwidth

8.3 Non-linear storage cell

The first designed cell can be seen Figure 8.9. In this structure, a common source structure is used to read the input capacitance. The advantage of this structure is its very large current gain, allowing to read the deposited charge with almost no current flow, and therefore minimum leakages. But the voltage V_{out} dependence with V_{in} is not linear.

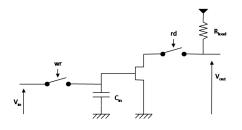


Figure 8.9: Electric structure of the non-linear storage cell

8.3.1 Improved structure

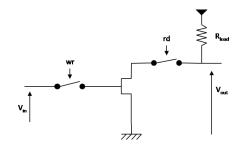


Figure 8.10: Storage cell using the gate capacitance

In the structure Figure 8.10, there is no more physical input capacitance. The capacitance used to store the signal is the gate capacitance of the Nfet. This structure presents two main advantages.

First, the gate capacitance which already exists and was in parallel with the input capacitance (Fig 8.9) is not seen as a parasitic capacitance anymore.

Secondly, as this input capacitance is very small, the input bandwidth will be more larger.

8.3.2 Large-signal analysis

Depending of the input voltage, the transistor can be operating in different regions:

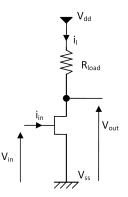


Figure 8.11: Large-signal schematics

Cutoff region

If $v_{GS} - V_T \leq 0$, then $i_D = 0$

Yet: $v_{GS} = v_{in} - V_{SS}$ That is to say: for v_{in} from V_{SS} to $V_{SS} + V_T$; $v_{out} = V_{DD}$.

Saturation region

From the large signal schematics (Fig 8.11) we can derive the following equations:

 $i_L = \frac{\mu C_{ox}}{2} \frac{W}{L} (V_{GS} - V_T)^2$ With:

 $v_{GS} = v_{in} - V_{SS}$

And:

 $v_{out} = V_{DD} - R_L i_L$

So:

 $v_{out} = V_{DD} - R_L \frac{\mu C_{ox}}{2} \frac{W}{L} (v_{in} - V_{SS} - V_T)^2$

The previous equation justify the non-linear behavior of this cell: the output depends quadratically on the input. Such a transfer function does not fit very well with the further digitization. But this structure has the advantage of being very fast.

Triode region

If $0 < v_{DS} \leq (v_{GS} - V_T)$ then $i_L = \frac{\mu C_{ox}}{2} \frac{W}{L} \left[(v_{GS} - V_T) - \frac{v_{DS}}{2} \right] \frac{v_{DS}}{2}$ The transistor will then enter the saturated region when $v_{DS}(sat) = v_{GS} - V_T$. Yet, from the schematics Fig 8.11:

 $v_{DS} = v_{out} - V_{SS}$ $v_{GS} = v_{in} - V_{SS}$ $v_{out} = V_{DD} - R_L i_L$

So:

$$v_{GS} - V_T = v_{DS}(sat)$$

$$\Leftrightarrow v_{in}(sat) - V_{SS} - V_T = v_{out}(sat) - V_{SS}$$

$$\Leftrightarrow v_{in}(sat) - V_{SS} - V_T = V_{DD} - V_{SS} - R_L \frac{\mu C_{ox}}{4} \frac{W}{L} \left(v_{in}(sat) - V_{SS} - V_T\right)^2$$

So, the transistor will enter the triode region when:

$$\Leftrightarrow v_{in}(sat) - V_{SS} - V_T = \frac{4L}{R_L \mu C_{ox} W} \left(\sqrt{1 + \frac{(V_{DD} - V_{SS}) R_L \mu C_{ox} W}{L}} - 1 \right)$$

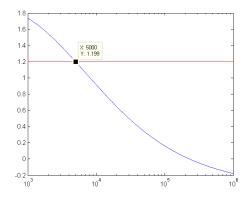


Figure 8.12: Plot of the $v_{in}(sat)$ value versus R_L

Figure 8.12 is plotted $v_{in}(sat)$ versus R_L . We can see that for a resistance value smaller than $5k\Omega$, $v_{in}(sat)$ is always higher than the maximum supply voltage (ie 1.2V). Therefore, for such resistance values, the transistor will remain in the saturation region.

Results

On Figure 8.13, the current generated by the transistor as a function of v_{in} for a resistance value of $10k\Omega$. With this value, we can see that the transistor enters the triode region for high values of v_{in} . In this region, we can clearly see that the output swing of the current is not as good as the saturation region. In order to improve this output swing one as therefore to use a resistance smaller than $5k\Omega$.

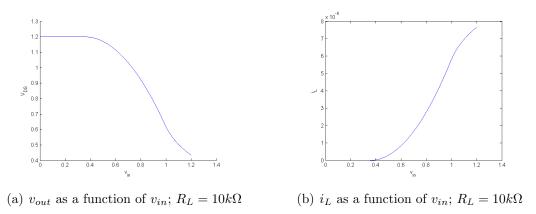


Figure 8.13: $v_{out} \& i_L$ as a function of v_{in} simulated on Matlab

Such a value for a resistor is hard to achieve with a good reproducibility from one cell to an other, and this resistor usually takes a lot of space in a layout, which can be annoying in our case where we foresee a high density integration.

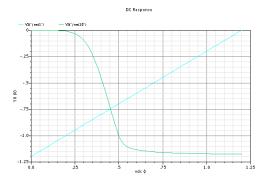


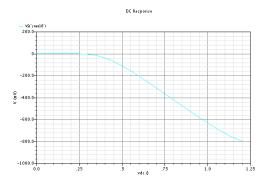
Figure 8.14: DC output of the non-linear cell

These results are in perfect agreement with the result of simulation from Spice. Figure 8.14, we can see the output of the storage cell for a very high value of R_L . The transistor enter deeply in the triode region and we can see (Figure 8.14), that the output swing is then completely flattened for high value of v_{in} .

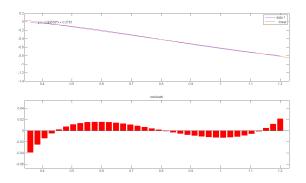
Linearity study

When we use a smaller resistance Fig 8.15(a)(here $12k\Omega$), we can see that the output is much more linear, and that it match well the results obtained from Matlab simulations.

However, using a smaller resistor, leads to a smaller output range: only 600mV here, out of the full range of 1.2V. This means that the full range of the ADC used to digitized the output signal, will not be used. And this will therefore leads to a much smaller precision in the digitized data.



(a) Storage cell with improved output linearity



(b) Results for a linear fit of the output

Figure 8.15: Linearity of the output over a 800mV output range

To conclude, this structure present the advantage of being fairly simple, with a fair linearity.

8.3.3 Small signal analysis

Small signal model of the circuit

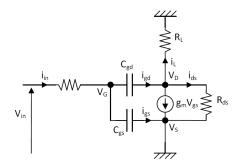


Figure 8.16: Small-signal model of the cell

Figure 8.16 is the small-signal model of the cell. With this model we can determine the transfer function of this cell, and its bandwidth.

Transfer function

In order to simplify this study, we will ignore, on the first order, the C_{gd} capacitance. Hence, from the Figure 8.16, we have :

$$i_{in} = i_{gs}$$

$$v_{GS} = \frac{1}{jC_{GS}\omega}i_{gs}$$

$$v_{in} - v_{gs} = R_S i_{in}$$
So:
$$v_{in} = v_{gs} (1 + jR_S C_{GS}\omega)$$

$$v_{DS} = v_{out} = R_{DS}i_{DS}$$

$$v_{out} = R_L i_L$$

$$i_L + g_m V_{GS} + i_{DS} = 0$$

$$\frac{v_{out}}{R_L} + g_m \left(\frac{v_{in}}{1 + jR_S C_{GS}\omega} + \frac{v_{out}}{R_{DS}}\right)$$

The transfer function of the cell is then the following:

$$\frac{v_{out}}{v_{in}} = -\left(\frac{R_L R_{DS}}{R_L + R_{DS}}\right) \frac{g_m}{1 + jR_S C_{GS}\omega}$$

From it, we can extract the gain of the cell :

$$G = g_m \left(\frac{R_L R_{DS}}{R_L + R_{DS}} \right)$$

Input bandwidth

From the previous study, we can see that the cutoff frequency of the cell is given by $R_S C_{GS}$. To improve the input bandwidth, one as therefore to reduce both R_S and C_{GS} . To do so, for C_{GS} , the dimensions of the input capacitor have to be the chosen as small as possible. But one cannot make them the smallest dimensions of the design, because if so, the capacitance parameters would be too much scattered.

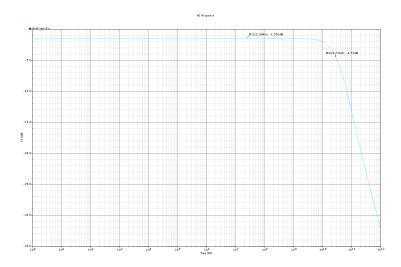


Figure 8.17: Input bandwidth

Figure 8.17, we can see the input bandwidth of the cell. We can see that the bandwidth is really high (around 20GHz), witch allow us a very fast sampling rate. However, this cutoff frequency is higher than the 13GHz cutoff frequency that we determine before. In this case, the switch will therefore the component limiting the bandwidth.

8.3.4 Read & Write state of the cell

The read and write sequence is not as simple as one could think. Indeed, Figure 8.18, one can see that the parasitic capacitances of a transistor are dependent of the state of the transistor. Therefore, in order to read and write the same input, the transistor has to be in the same state during the read and the write phase.

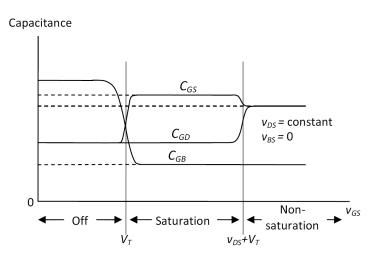


Figure 8.18: Voltage dependence of C_{GS} , C_{GD} , and C_{GB} as a function of V_{GS}

In order to have the transistor in the same state during the read and write phase, the read switch has to be closed during the write state too. The read & write sequencing is then given by the Figures 8.19

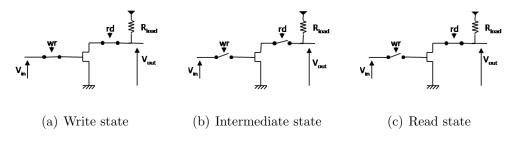


Figure 8.19: Read & write sequencing

8.3.5 Cell issues

Input bandwidth

Due to the non-infinite input bandwidth, the signal tracked during the write phase is not the same as the input; there is an attenuation for the higher frequencies as we can see in Figure 8.20.

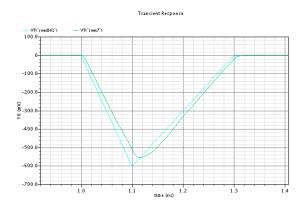


Figure 8.20: Visualization of the difference between the input signal and the signal on the sampling cell

Charge leakage

As we discussed previously, there are two sources of charge leakage: the capacitance and the switch. For the capacitance, the charge leakage is dependent of the area of the Nfet used. In our process, this charge leakage is quite important: $0.7pA/\mu m^2$.

The time necessary to loose 1mV is given by:

$$i = \frac{\partial q}{\partial t} = C \frac{\partial U}{\partial t}$$
Yet:

$$i = 11 f A$$

$$C = 200 a F$$
Then:

$$\frac{\partial U}{\partial t} = 55 V/s$$
So:

$$\Delta t_{1mV} = 18 \mu s$$

There is also the discharge of the capacitance trough the switch. The time constant of this discharge is given by $\tau = R_{OFF}C_{IN}$

Let's calculate the order of magnitude necessary for R_{OFF} to make the leakage trough the switch negligible compared to the leakage trough the Nfet.

The voltage on the gate is following the exponential law:

 $V = V_0 e^{-\frac{t}{\tau}}$

Let's take an input voltage of 600mV, and a time constant to loose 1mV ten times higher than the previous one: $\Delta t_{1mV} = 200 \mu s$

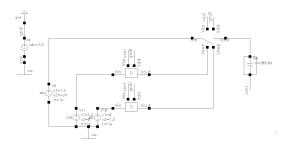
We have then:

$$V_0 - 1mV = V_0 e^{-\frac{t_{1mV}}{\tau}}$$
$$\tau = \frac{-\Delta t_{1mV}}{\ln\left(\frac{V_0 - 1mV}{V_0}\right)} = 0.12$$
$$R_{OFF} = \frac{\tau}{C} = 600T\Omega$$

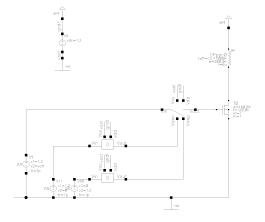
This determined value of R_{OFF} is much more higher than the real achievable value for the switch. It means therefore that the switch will be the most leaking component in our design.

Determination by simulation of the most leaking component

In order to find what is the limiting parameter the following simulation has been done: a circuit with an ideal capacitance (no leakage) of the same value of the transistor capacitance and the real circuit are been simulated Figure 8.21



(a) Switch with ideal input capacitance



(b) Switch with the standard input transistor

Figure 8.21: Comparison of the sampling with an ideal and non-ideal input capacitance

The two circuits of the Figures 8.21 have been simulated under the same conditions: 600mV is set at the input of each cell, then the write switch is open and the output of the both cells are compared Figure 8.22. What we basically see in this figure is the decreasing of the voltage stored on the capacitor due to the leakage. As the two curves are similar, we can obviously say that the leakage trough the capacitance is negligible compared to the leakage of the switch. A special effort has therefore to be made for the design of the switch in order to reduce this leakage. This means increasing R_{OFF} .

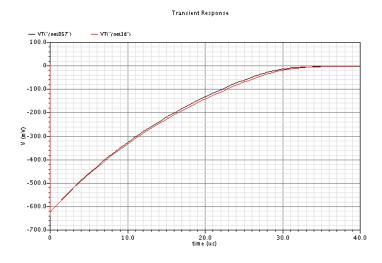


Figure 8.22: Comparison of the stored value with an ideal capacitance (red) and the real circuit (black)

Charge injection

The most constraining parasitic effect in this storage cell during the write phase is the charge injection of the switch. With no compensation, the charge injection is presented Figure 8.23. We can see that whereas the value stored on the input capacitance is the desired one (600mV); when the switch is closed, a charge (negative) is injected making the input value evolve.

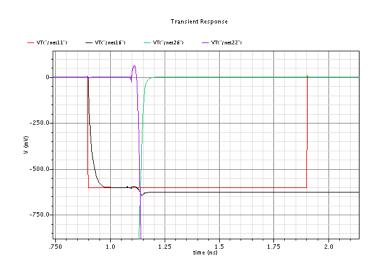


Figure 8.23: Visualisation of the charge injection when closing the write switch

This charge injection can be reduced by an elaborated design of the switch [15], or could be either taken care of by calibrating each cell.

8.3.6 Read state

During the read state the write switch is open and the read switch closed. The output is available and can be sent to the comparator for the digitization. During this phase, the output must stay as stable as possible and as close as possible to the value of the corresponding input. The charge leakage must be as small as possible and the digitization as fast as possible.

Figure 8.24, is the schematic view of the storage cell simulated with Spice. The result of the simulation can be seen Figure 8.25. There are three phases, according to the previous section. On the figure, we can clearly see the input voltage evolving on the gate of the Nfet for different regimes. But in good agreement with the previous section the voltage is restored to the initial value during the read phase.

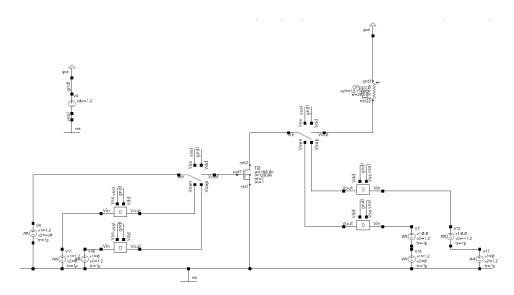


Figure 8.24: Simulated storage cell

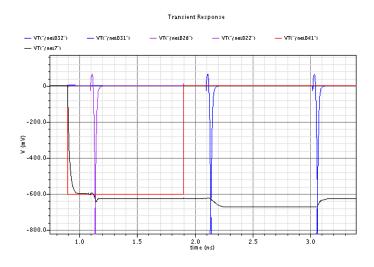


Figure 8.25: Evolution of the output during the read and write states

8.4 Linear storage cell

The other structure that can be found in the literature is the commondrain structure, basically linear. The structure of this cell can be seen on Figure 8.26.

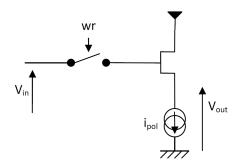


Figure 8.26: Linear storage cell

8.4.1 Large signal analysis

From the Figure 8.26:

Cutoff region

If, $V_{in} < V_T$ then no current can flow trough the transistor and $V_{out} = V_{SS}$

Saturation region

If,
$$V_{in} > V_T$$

 $V_{out} = V_{in} - VGS$
Yet:
 $V_{GS} = V_T + \sqrt{\frac{2i_{pol}}{\mu_0 C_{ox} \frac{W}{L}}}$
So:

$$V_{out} = Vin - V_T - \sqrt{\frac{2i_{pol}}{\mu_0 C_{ox} \frac{W}{L}}}$$

From this formula we can clearly see that the output will be much more linear than with the previous one. Indeed, $V_{out} = Vin + K$.

8.4.2 Linear cell with real current source

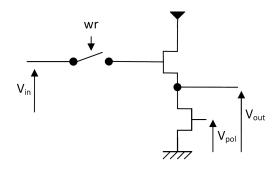


Figure 8.27: Linear storage cell with a real current source

In order to implement the current source (Fig 8.26) of the circuit, we use a Nfet biased by the voltage V_{pol} (Figure 8.27), we have then:

$$i_{pol} = \frac{\mu_0 C_{ox}}{2} \frac{W_{pol}}{L_{pol}} (V_{GS} - V_T)^2$$

So:

So:

$$V_{out} = Vin - V_T - (V_{pol} - V_T) \sqrt{\frac{W}{L} \frac{L_{pol}}{W_{pol}}}$$

DC characteristic

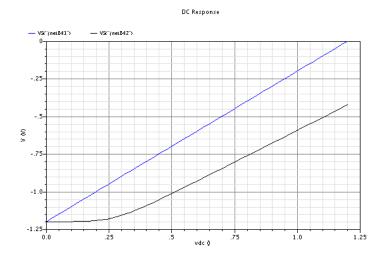


Figure 8.28: DC characteristic of the cell

Figure 8.28 the output (black) vs input (blue) of the cell is plotted. One can see that the linearity of the cell is particularly good. However, having an output using the full range of the DC sources (from GND to VDD), is not achievable.

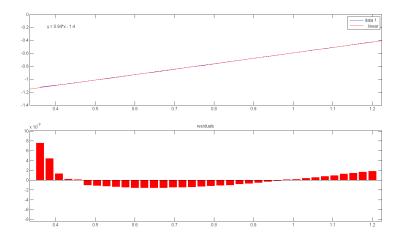


Figure 8.29: Linear fit of the output

We can check this linearity on Figure 8.29. The linearity is improved by a factor 10 compared to the previous cell.

8.4.3 Small signal analysis

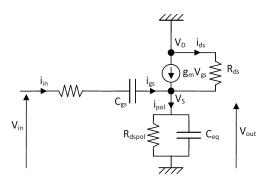


Figure 8.30: Small signal model of the cell

From the Figure 8.30 we have:

$$V_{in} - V_{out} = i_{in} \left(R_{in} + \frac{1}{jC_{GS}\omega} \right)$$
$$V_{out} = r_{DS}i_{DS}$$
$$V_{out} = i_{pol} \left(r_{pol} + \frac{1}{jC_{eq}\omega} \right)$$

 $i_{in} + g_m V_{GS} + i_{DS} = i_{pol}$

At the first order we can neglect i_{DS} and i_{in} compared to the other current level.

Yet:

$$g_m V_{GS} = i_{pol}$$

$$V_{out} = g_m V_{GS} \left(R_{eq} + \frac{1}{jC_{eq}\omega} \right)$$

$$V_{GS} = \frac{1}{jC_{gs}\omega} \frac{V_{in} - V_{out}}{R_{in} + \frac{1}{jC_{gs}\omega}} = \frac{V_{in} - V_{out}}{1 + jR_{in}C_{gs}\omega}$$

So:

$$V_{out} = (V_{in} - V_{out})g_m R_{eq} \frac{1}{1 + jR_{in}C_{gs}\omega} \frac{1}{1 + jR_{eq}C_{eq}\omega}$$

Finally:

$$\frac{V_{out}}{V_{in}} = \frac{1}{1 + \frac{(1+jR_{eq}C_{eq})(1+jR_{in}C_{gs})}{g_m R_{eq}}}$$

From this equation, we can see that in addition to the expected input pole $(R_{in}C_{gs})$, a second pole is added at the output due to the floating source potential of the Nfet. This pole is given by $R_{eq}C_{eq}$. R_{eq} is in fact the R_{DS} of the biasing transistor and C_{eq} is the sum of all the capacitance of the biasing transistor.

8.4.4 Input & Output bandwidth

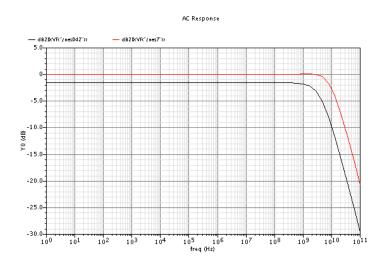
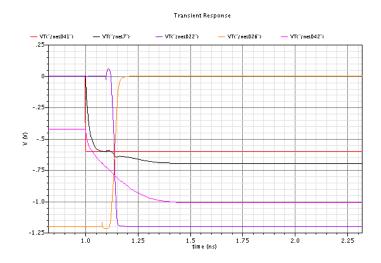


Figure 8.31: Input and output bandwidth of the cell

On Figure 8.31, we can see the input and output bandwidth of the cell. We can check that the input bandwidth is as high as the bandwidth obtained with the previous cell. But the output is slower, due to the second pole.



8.4.5 Read state

Figure 8.32: Evolution of the input and the output during the read state, the red curve is the input, the black one is the signal stored in the input Nfet, the yellow and purple curves are the control signals for the switch, and the pink curve is the signal at the output of the cell.

Figure 8.32, we can see during the read and write states the input (in black) and the output (in pink). As previously stated, the output is evolving much slower than the input. Therefore, the source voltage of the Nfet is evolving during the read state. Yet, in Figure 8.18, we saw that the value of the input capacitance depends on the working condition of the Nfet. If V_S is evolving it will make V_G evolve too.

However, it was found in an improved study of the cell, that this pole can be rejected high enough to allow a correct reading and writing, as it can be seen on Figure 8.33

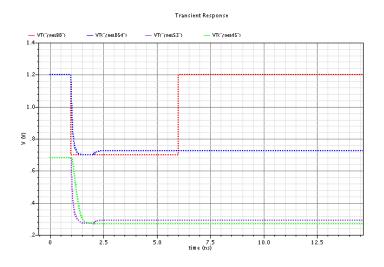


Figure 8.33: Transient signals with second pole rejection

8.5 Conclusion: cell chosen

After studying in detail these several structures of storage cell, we choose to use the linear storage cell, mainly because of it's linearity and also because it's response time can be adjusted by the level of the current source. An other advantage of this cell is the small area it covers when laid out, due to the absence of resistor in it. This will permit an easy integration of a multi-cell structure.

9 Design

Following the previous study a storage cell has been designed for the sampling chip. It has been done using the kit CMRF_8_SF of IBM under Cadence.

9.1 Storage cell

9.1.1 Schematic view

The schematic view of the storage cell is shown Figure 9.1

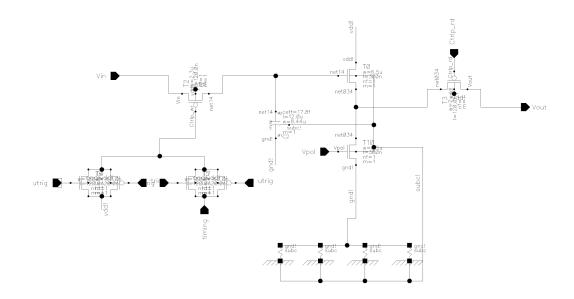


Figure 9.1: Schematic view of the storage cell

The different parts of the cell can be seen on Figure 9.2. The cell consists of five different parts:

• The input switch

- The input capacitance and Nfet
- The current source
- The multiplexor
- The output switch

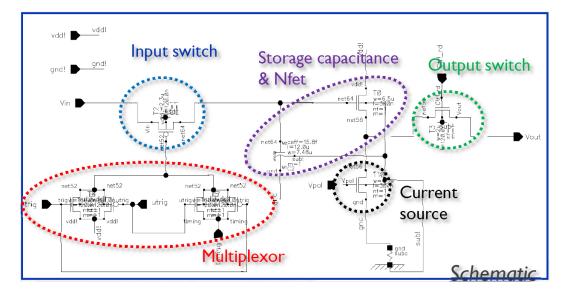


Figure 9.2: Schematic bloc description

9.1.2 The input switch

The present timing generator, does not have a fully differential output. Therefore it is not possible to control a PFet, Nfet switch. We therefore only used a single Pfet switch, allowing a signalrange from 1.2V to .4V at the input. A better switch with a differential output from the timing generator, would cover full-range signals at the input, and would also have a smaller value for the resistance of the switch, and reduce the charge injection. The actual charge injection can be seen on Figure 8.33. It is not negligible at all, but as it will be identical for every cells, it can be seen as a parasitic offset that can be canceled by a calibration of each cell.

In the next chip, this input switch will be improved a lot, mainly having the timing generator output differential. Therefore, charge injection cancellation techniques will be used, the input resistance will be reduced leading to a wider input bandwidth, and the input range will cover the full range allowed by the tecnology.

9.1.3 The input capacitance and Nfet

In this chip we choose to add an additional input capacitor to the parasitic capacitor of the input Nfet, indeed the value of this one was too small 20fF, and therefore too much subject to leakage, noise, and charge injection.

Furthermore, the characteristic of a real capacitor are much reliable than a parasitic one, in particular the capacitance does not depend on the bias voltages at its input (as it is the case for the parasitic transistor of the Nfet) and are also very defined. Unfortunately, the layout of a physical capacitor takes a lot more space than a simple transistor, leading to a more complicated layout and integration.

9.1.4 The current source

The current source is a simple Nfet driven externally, the value of the current set by the current source controls the time response of the circuit (the more current, the fastest is the circuit) but also the drop of voltage between the input and the output. Having this external control of the current source is mandatory in this first chip in order to tune the cell and get the best results.

9.1.5 The multiplexer

The multiplexer is the structure that controls the input signals to be switched. Indeed the switch can be either controlled by the timing signal or the trigger signal when an event has been recorded. This structure is a simple CMOS multiplexer.

9.1.6 The output switch

We did not detail this switch in the previous section, and yet it plays a very important role in the chip as we will see now:

Indeed, when simulating this chip without it, we were aware that the leakages in the cell were very important, and that not taking care of them, would lead to loose a huge amount of our signals even before starting the digitization.

We can see that the subthreshold current, which is the leakage current

of an open switch strongly depends of the state of it:

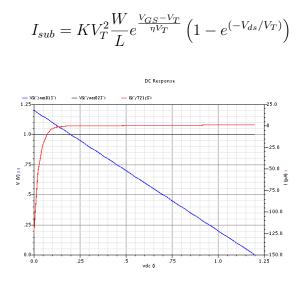


Figure 9.3: Plot of the subthreshold current function of V_{GS}

We can can see on Figure 9.3, that for an input leading to a very small value of V_{GS} , the subthreshold current increase exponentially, leading to an important leakage current. The only way to avoid this important leakage current is to have a switch operating with a different value of V_{GS} . This what is done using the output switch, indeed at the output of the cell, there is a drop of tension due to the current source. Therefore, if a switch is put here the V_{GS} value will be increased, and therefore the leakage current exponentially reduced. The results of this improvements can be seen on Figure 9.4.

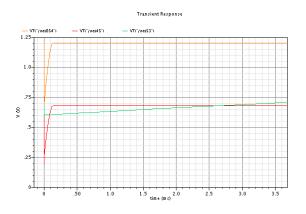


Figure 9.4: Difference of voltage droops at the output of the cell with (green) and without (red) read switch.

The drawback of this additional switch is the fact that it will add extra charge injection.

9.1.7 Layout

The layout of this cell can be seen on Figure 9.5.

Starting from the bottom of the figure, we can first see the storage capacitance, taking most of the space. This capacitor is an inter-digited capacitor, in order to avoid cross-talk from the adjacent capacitor, it has been shielded by a top metal plate and the side wall.

Then we can see on the left the input Nfet, which has been designed close to the current source to obtain a better matching.

On the right, we have first the multiplexer that will select between the timing control signal and the trigger. And then the input switch.

Going towards the top of the cell, we can then see the structure of the input line for the signal. This structure will be discussed along with the integration of the cells to make a channel.

On the very top of the cell, one can see a series of three double inverters. These inverters are in fact reshapers for the fast control signals: timing control, trigger, and output switch control.

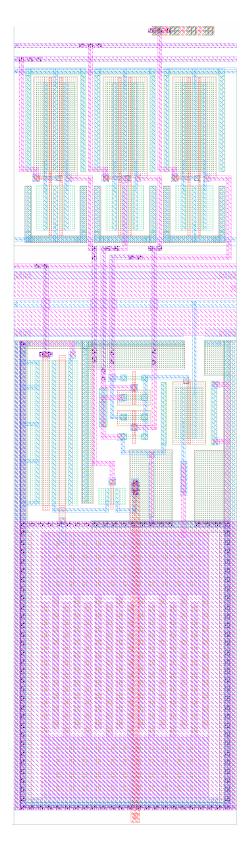


Figure 9.5: Layout of the storage cell

9.2 Storage cell assembly

The further stage for the chip is the integration of 256 storage cells in a row, in order to obtained a full sampling unit (Fig 9.6).

Figure 9.6: Zoom in the layout of one channel

The layout of one channel in our case is not as simple as putting all the cells together, indeed, the signal are coming from a 50 Ohms matched coplanar line at the input. Non matching this to 50 Ohms for the propagation of the signal in he channel would result in reflections to the input of the chip and can lead to degradations of the waveform in the 1-10ps range.

Some RF lines, provided by the IBM design kit, are design to be 50 Ohms matched. Unfortunately, using them stright was not possible, so the layout was made under the constraints dictated by the context.

A section of the input line can be seen on Figure 9.5 between the inverter and the control structure. We can see on this picture the signal trace running in the middle (blue wire) shielded by the pink wires on each side. The top of the signal wire was not covered in order to reduce the capacitance of the line and therefore high input bandwidth.

9.3 Channel

On Figure 9.7 a zoom of 64 digitizing cells of a complete channel is shown. A digitizing cell comprises a storage cell, a comparator, and a counter.

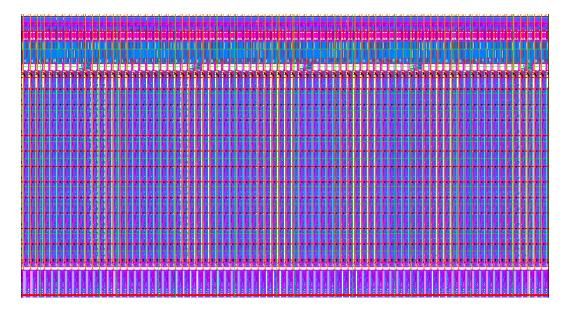
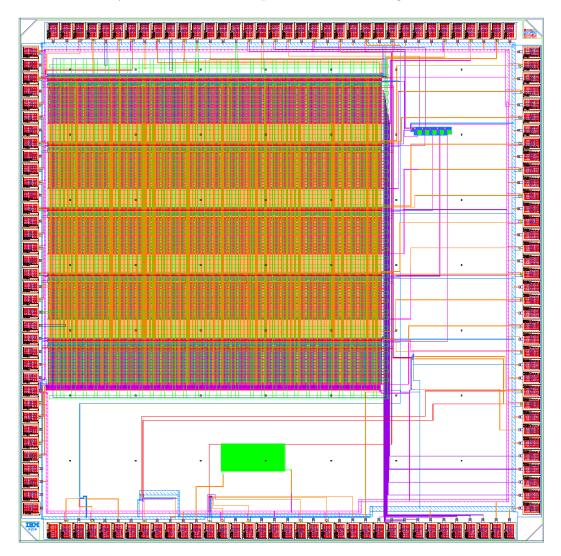


Figure 9.7: Layout of one channel

The storage cells are on the top of the channel, followed by the comparator and the counter. The signal from the timing generator are coming from the top to each cell, and the signals from the token controlled readout register are coming from the bottom to every counter, asserting the data on the output bus laid out above the channel.

9.4 Chip



The layout of the entire chip can be seen on Figure 9.8.

Figure 9.8: Final layout of the chip submitted

On this Figure, we can clearly see the four input channels on the top left. There is a fifth channel that will be used to monitor the signal coming from the timing generator.

At the bottom of the chip some test structures have been placed: a sampling cell, a comparator, a ring oscillator and a counter.

All the structures are tied to the 144 pads of the chip (in red). The connectivity of each pad is shown Figure 9.9.

This chip has been successfully submitted to MOSIS, an academic facility which provides access to fabrication of prototype and low-volume pro-

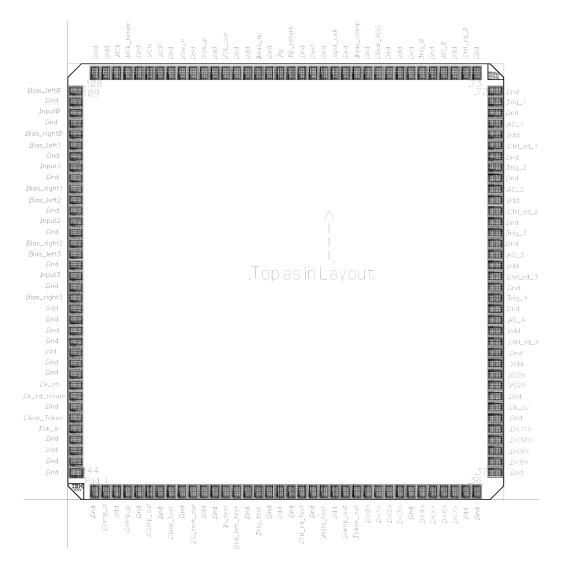


Figure 9.9: Chip pinou

duction quantities of integrated circuits at low cost.

10 Conclusion

During this internship, a new very fast high bandwidth chip has been designed for pico-second time-of-fight applications, using a very small technology (130nm IBM). The internship started by an in-depth study of the picosecond project and the sampling cells present in the literature. Then the sampling process has been studied, simulated, designed and laid out using the Cadence software tool. The cells have being integrated in a channel allowing a full digitization of incoming signals. The internship ended by designing the full chip, and successfully submit it to the MOSIS academic after verification and stream out. This internship was a wonderful experience, it allowed me to acquire a unique experience in CMOS VLSI design, and therefore complete adequately my student formation. Doing this internship in the United States, in particular at the University of Chicago, i got an incredible experience in both technique and linguistic, therefore I would like to thank my professors H. J. Frisch and J. F. Genat, for choosing me and giving me this opportunity, and also for all the support and the attention they provided me.

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