

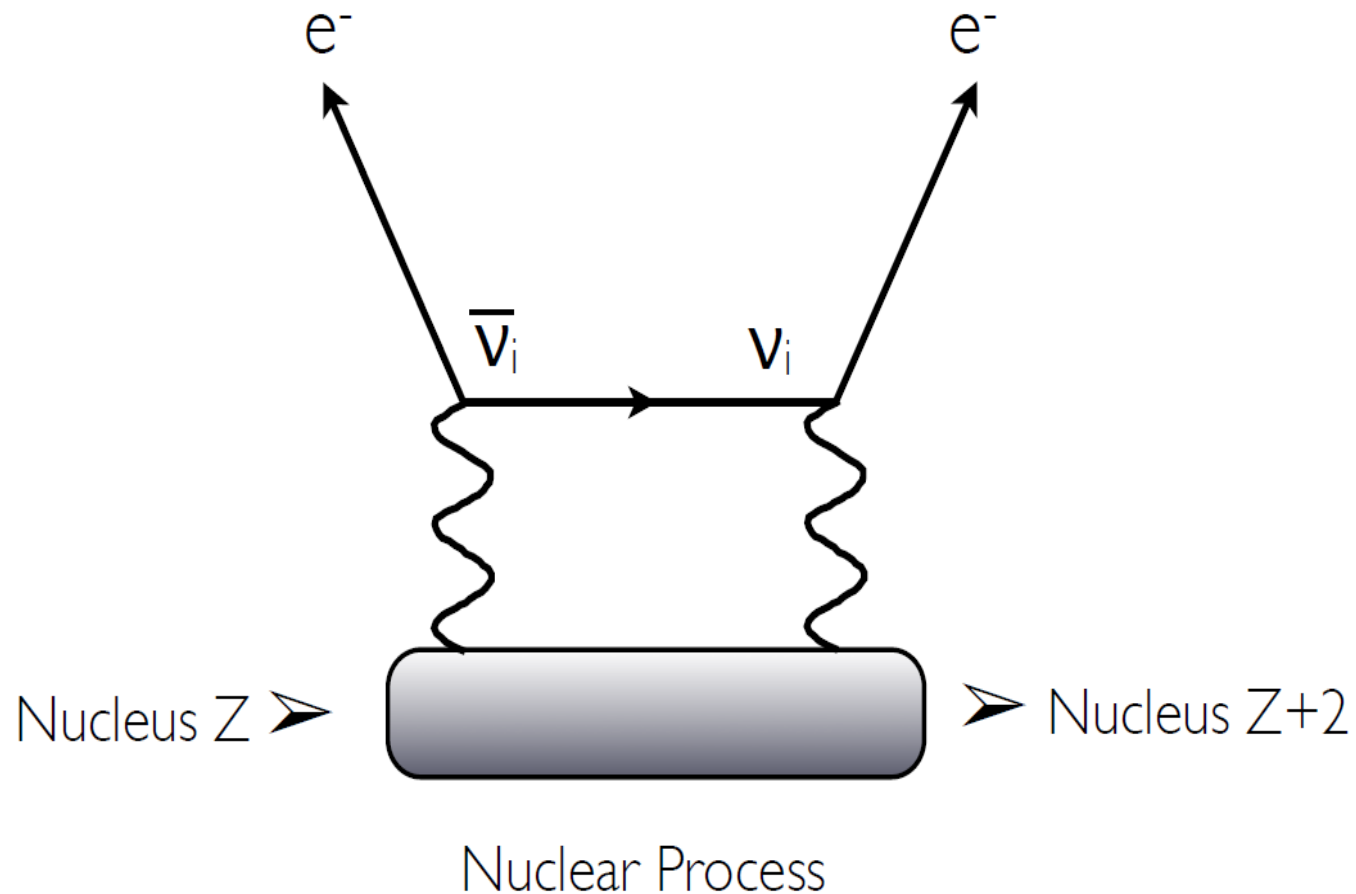
Hermetic Packaging of the LAPPDs and Fast Timing Implications for Neutrino-less Double-Beta Decay Searches

Andrey Elagin
University of Chicago



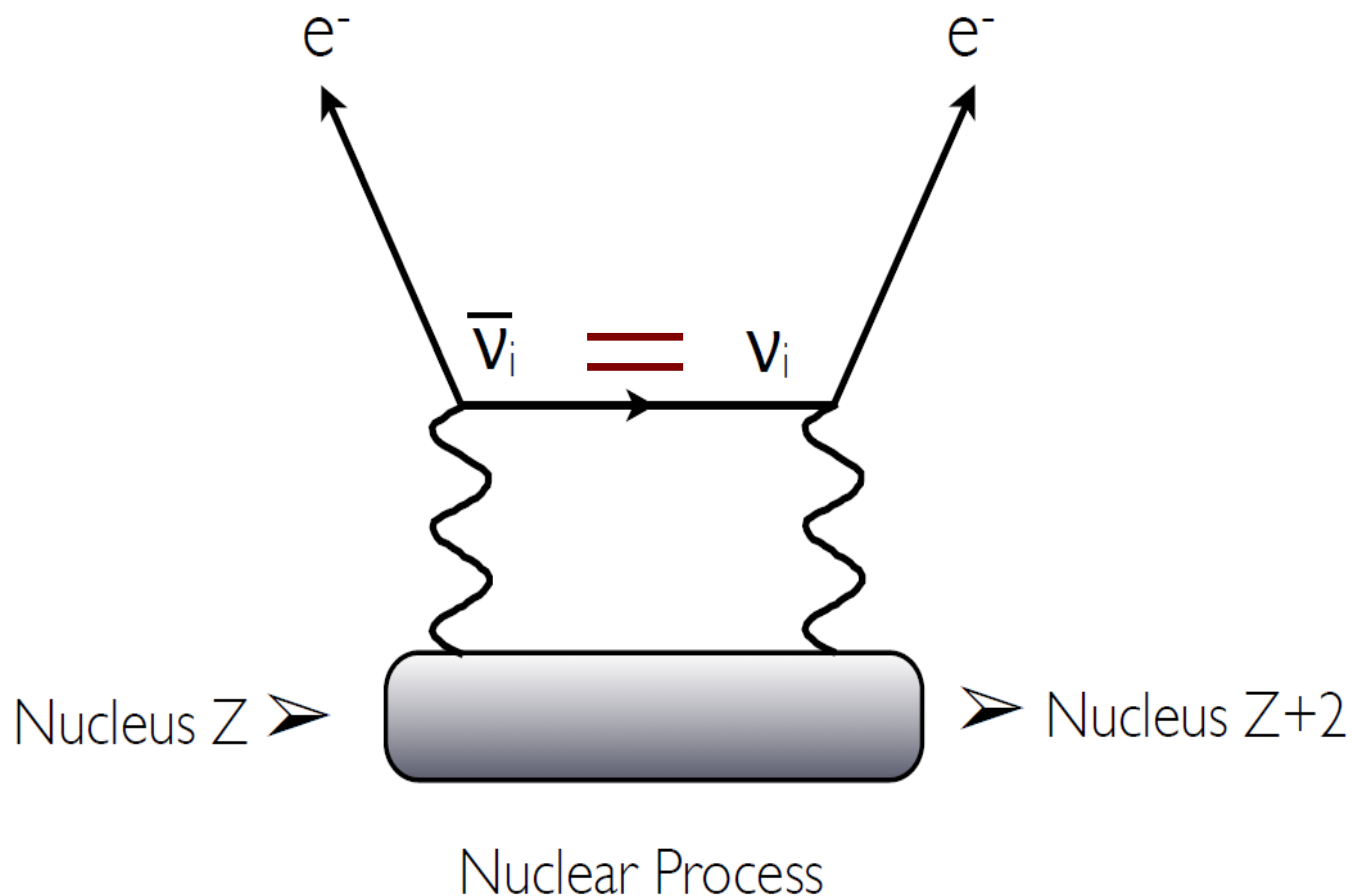
HEP Lunch Seminar 10/28/2013

What is $0\nu\beta\beta$?



Compare to normal beta decay: $Z \rightarrow (Z+1), e^-, \bar{\nu}_e$

Why is it interesting?



If observed, neutrino is a Majorana particle, i.e. own antiparticle.

Signature: two electrons with well defined total kinetic energy (2-4 MeV)

What are the challenges?

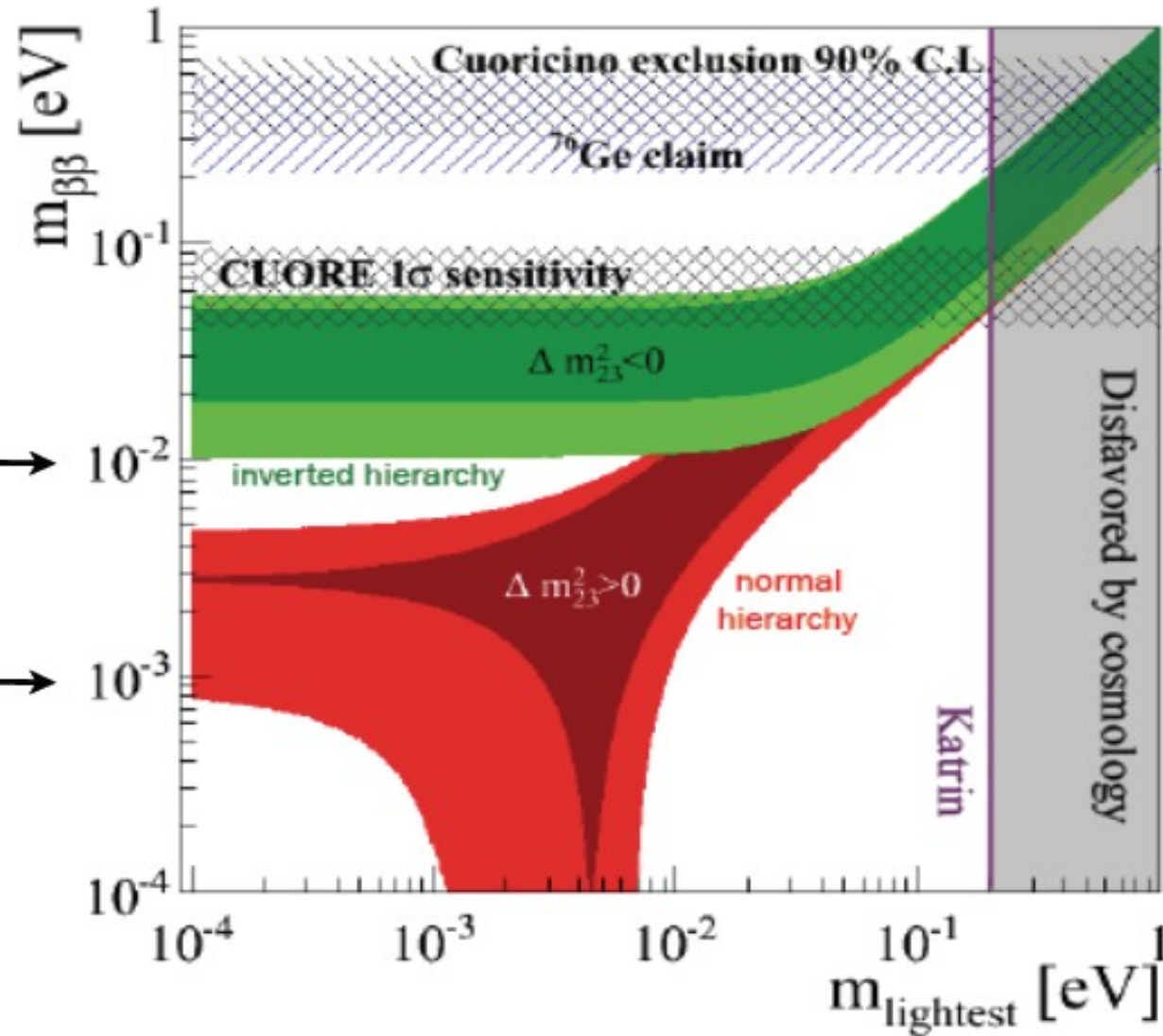
Rare process (e.g. $T_{1/2} [^{136}\text{Xe}] > 10^{25}$ years): **need to get bigger**

$$\Gamma = G |M|^2 |m_{\beta\beta}|^2$$

EXO ~ 32 kg yr \rightarrow

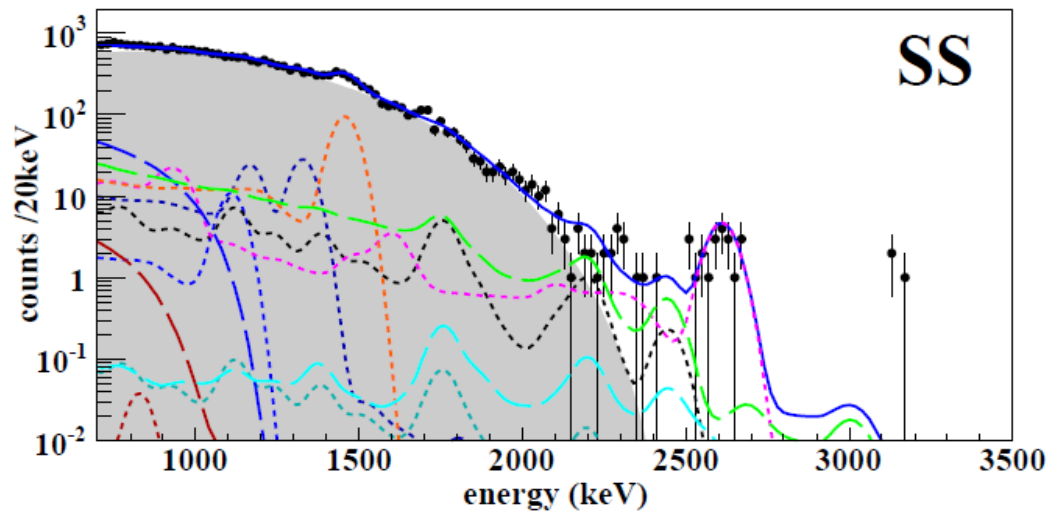
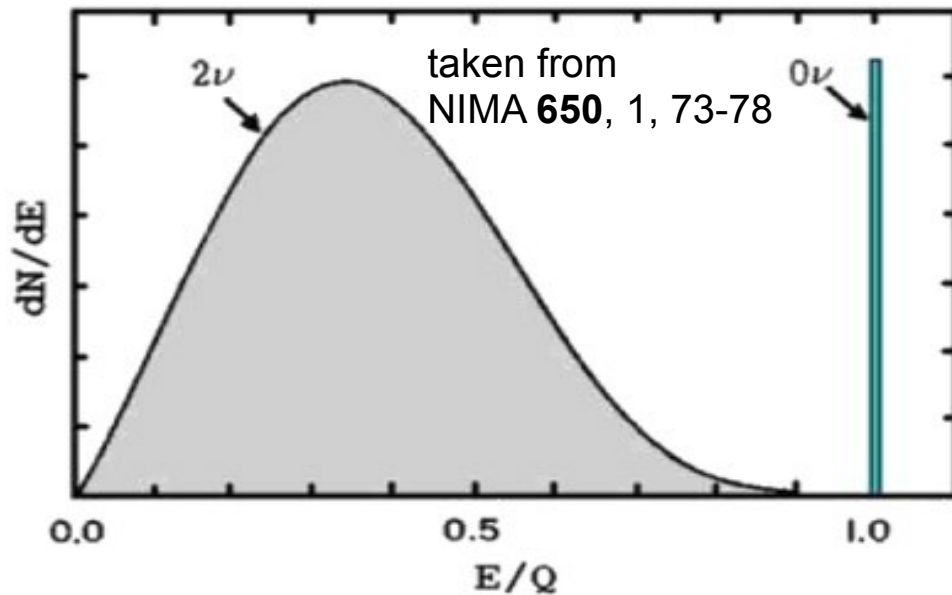
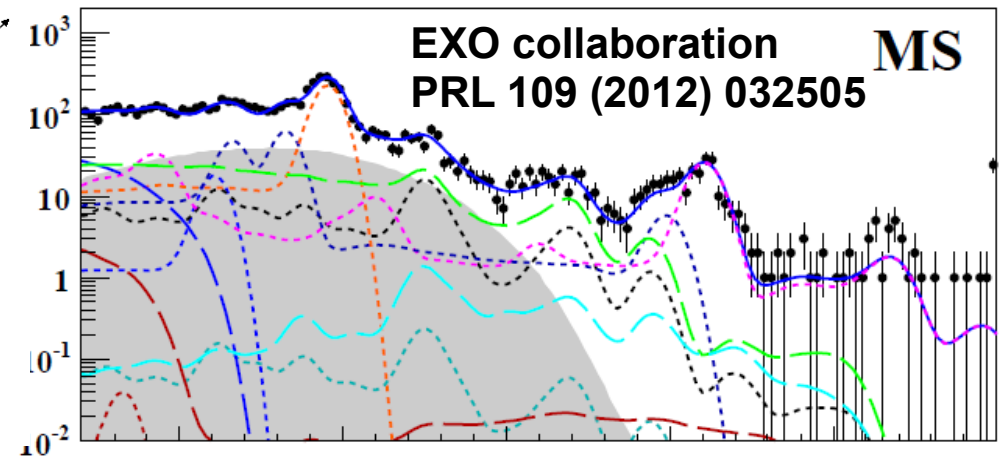
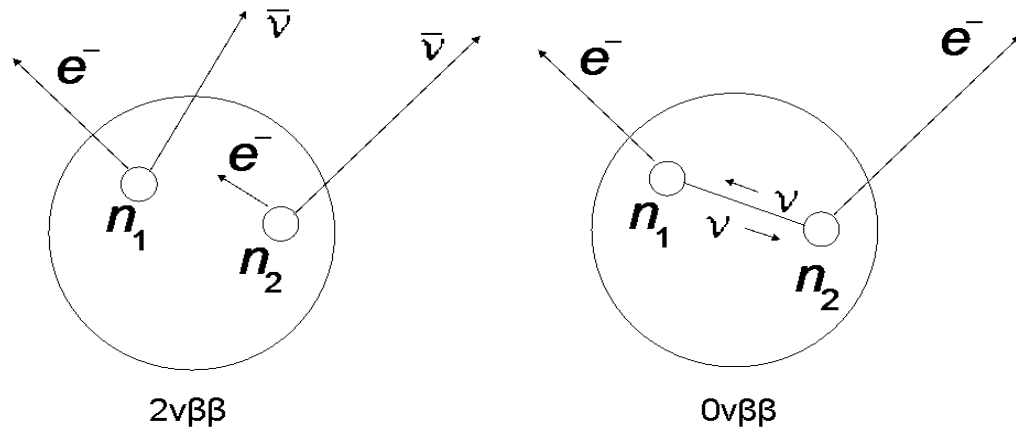
1 ton \rightarrow

10 ton \rightarrow



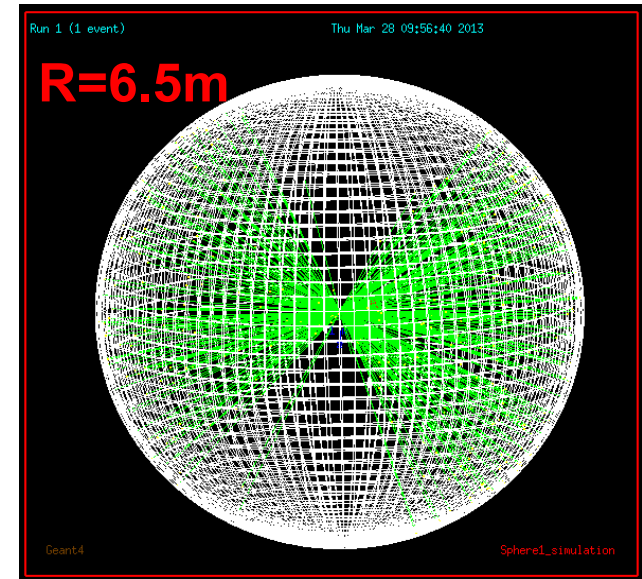
What are the challenges?

Tough backgrounds: **need to get smarter**



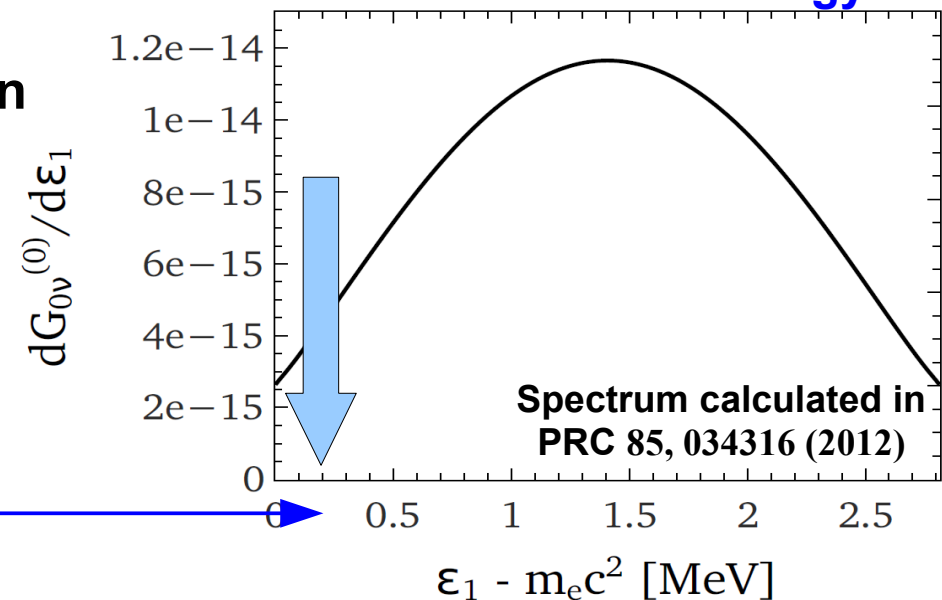
Ideas for $0\nu\beta\beta$ experiments

Simulation of ^{116}Cd $0\nu\beta\beta$ event



- Total energy in signal events is well defined.
- Use scintillation light for energy measurements
- Use event topology to suppress backgrounds
 - signal is two, mostly, “back-to-back” electrons
- Electrons are $\sim 1\text{MeV}$ \rightarrow above Cherenkov threshold
- Use Cherenkov light to extract directionality of the two electrons
 - all light can be used to constrain location of the vertex
 - Cherenkov light arrives early because of longer wavelength and delay of the scintillation process

Electron kinetic energy



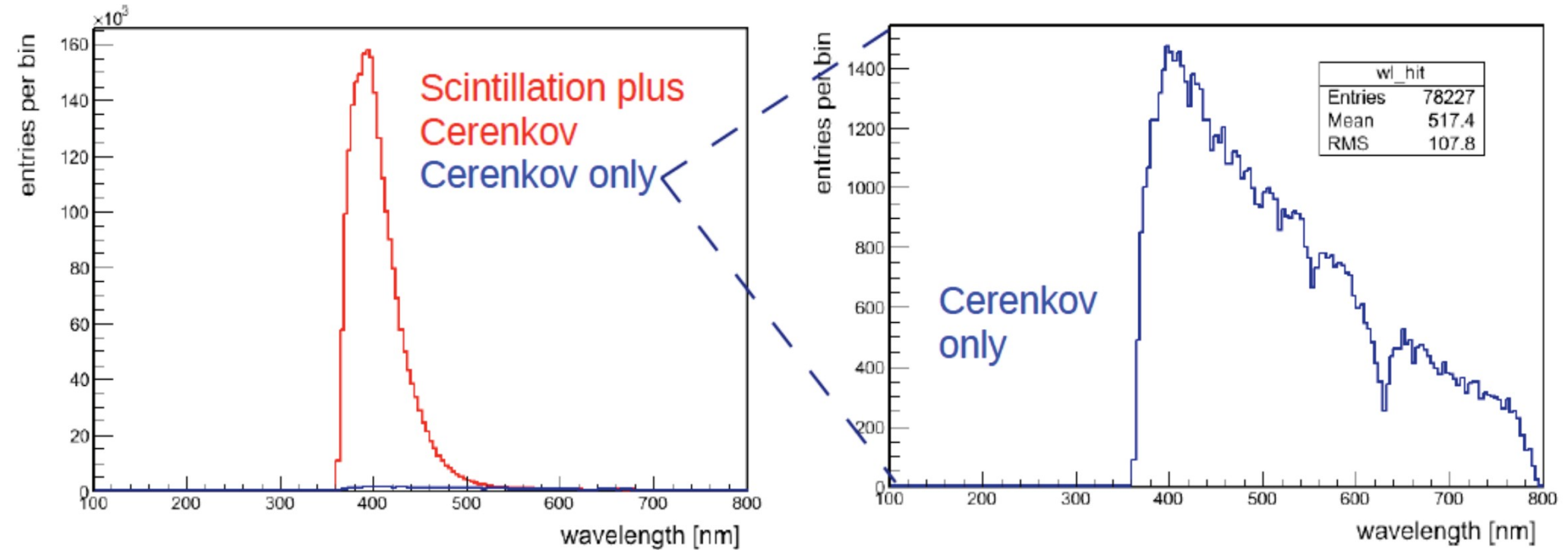
Cherenkov threshold for $n=1.47$

Emission Spectra

Simulation of 5 MeV electrons in KamLAND scintillator

5 MeV is a little higher for $0\nu\beta\beta$ search but much lower than typical energies where cherenkov light is being considered.

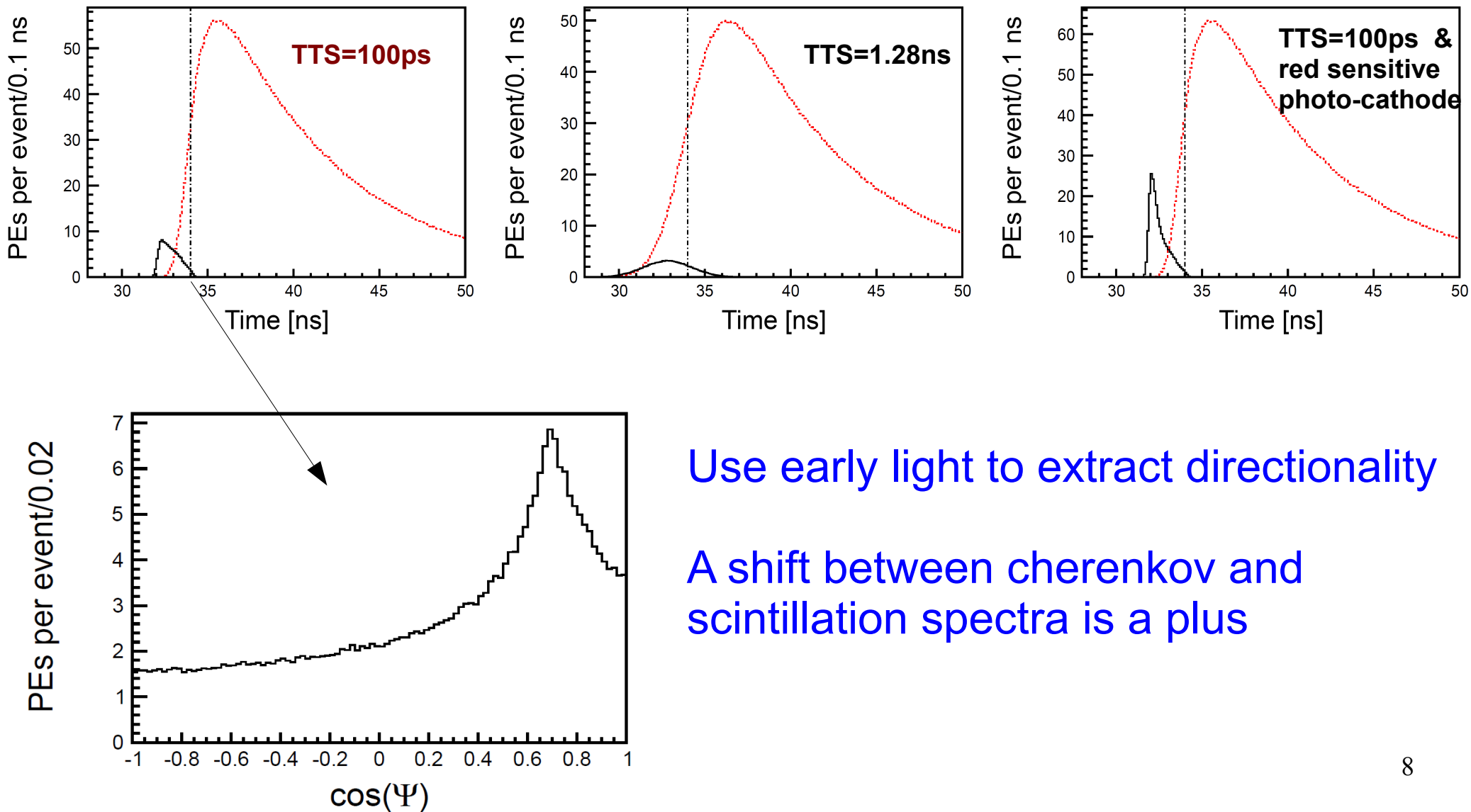
Seems to be a reasonable choice to test separation between Cherenkov and scintillation



All photons below 360nm get absorbed

Cherenkov vs Scintillation

Simulation of 5 MeV electrons in KamLAND scintillator



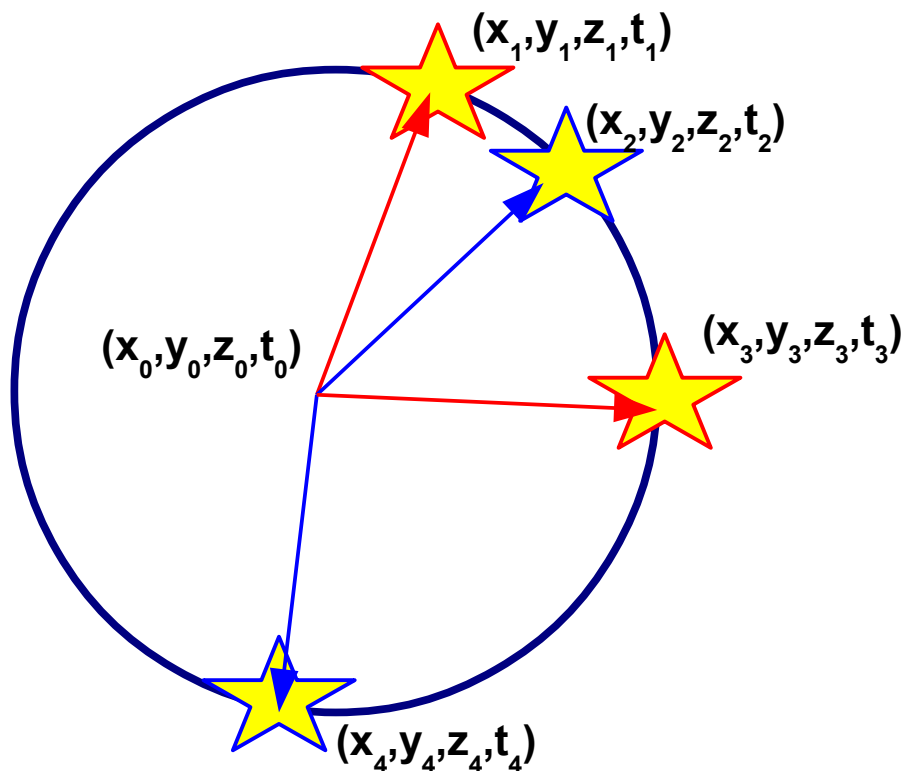
Use early light to extract directionality

A shift between cherenkov and scintillation spectra is a plus

Reconstruction: vertex

Step 1: find vertex (adapted from water cherenkov)

- Assume all light is emitted from a single point (~ 3 cm track in a ~ 6 m detector)
- For light emission from a single point any 4 photons (quadruples) would be sufficient to solve for vertex
- With all „real world“ effects we use 400 randomly chosen quadruples and select the one which fits the best to the full ensemble of all photon hits
- Goodness of the fit is based on the distribution of „point time residuals“ (the difference between actual hit time and predicted time of flight from the vertex)



For $i=1\dots 4$:

$$(x_i - x_0)^2 + (y_i - y_0)^2 + (z_i - z_0)^2 = v(t_i - t_0)^2$$

v is the speed of light in the media

Time residuals $dT = (t_j - t_0)$,

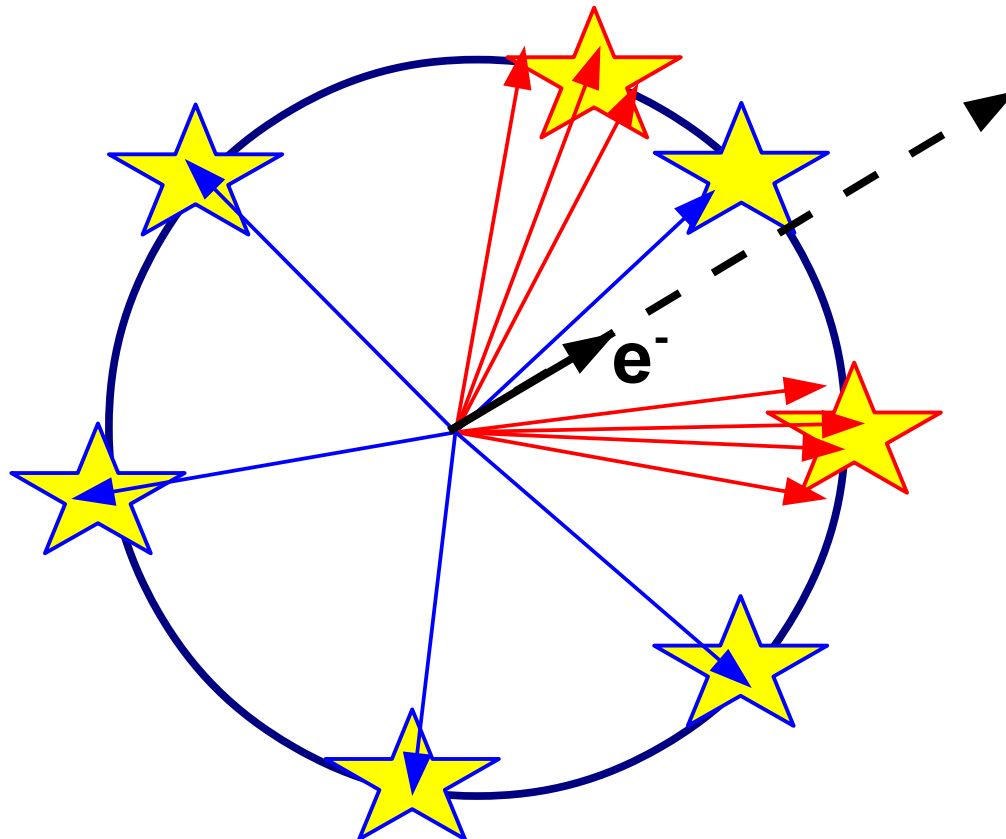
where $j=1\dots N_{\text{all}}$

The most narrow dT distribution is the closest to the true vertex

Reconstruction: direction

Step 2: find direction

- Cherenkov light is directional
- Timing cut enhances the purity of the Cherenkov light
- The centroid of all vectors pointing from the vertex is a good measure of the direction of the track



Directionality of 5 MeV e^-

Measuring Directionality in Double-Beta Decay and Neutrino Interactions with Kiloton-Scale Scintillation Detectors

C. Aberle,¹ A. Elagin,² H. J. Frisch,² M. Wetstein,² and L. Winslow¹

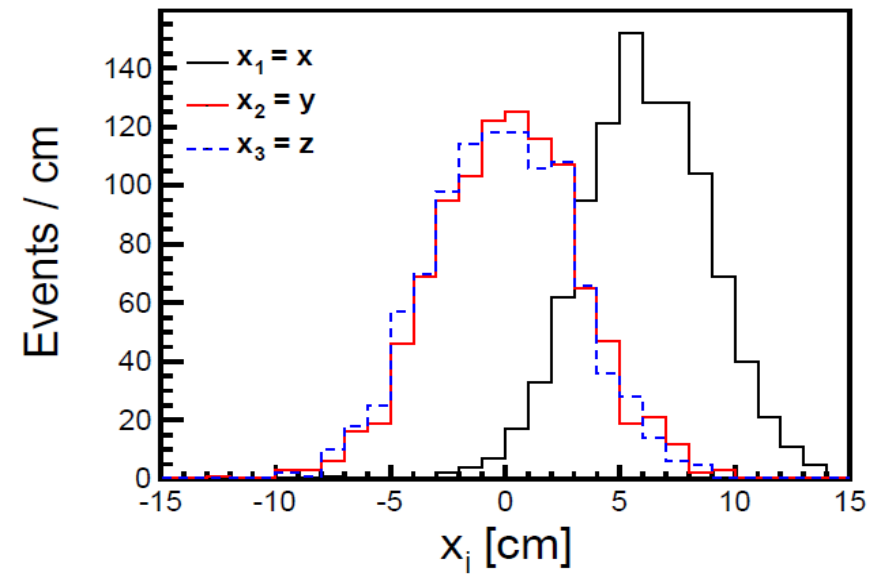
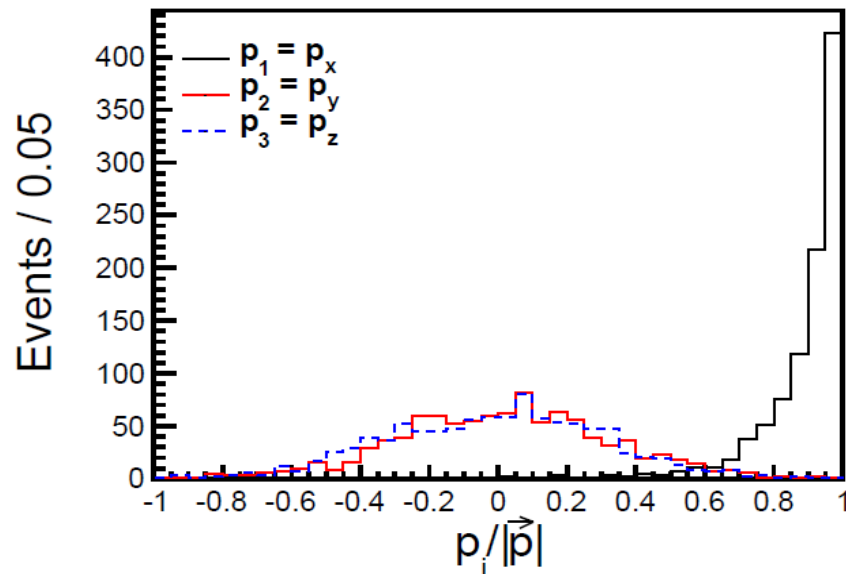
¹University of California Los Angeles, Los Angeles, CA 90095, USA

²University of Chicago, Chicago, IL 60637, USA

(Dated: July 23, 2013)

[arxiv:1307.5813](https://arxiv.org/abs/1307.5813)

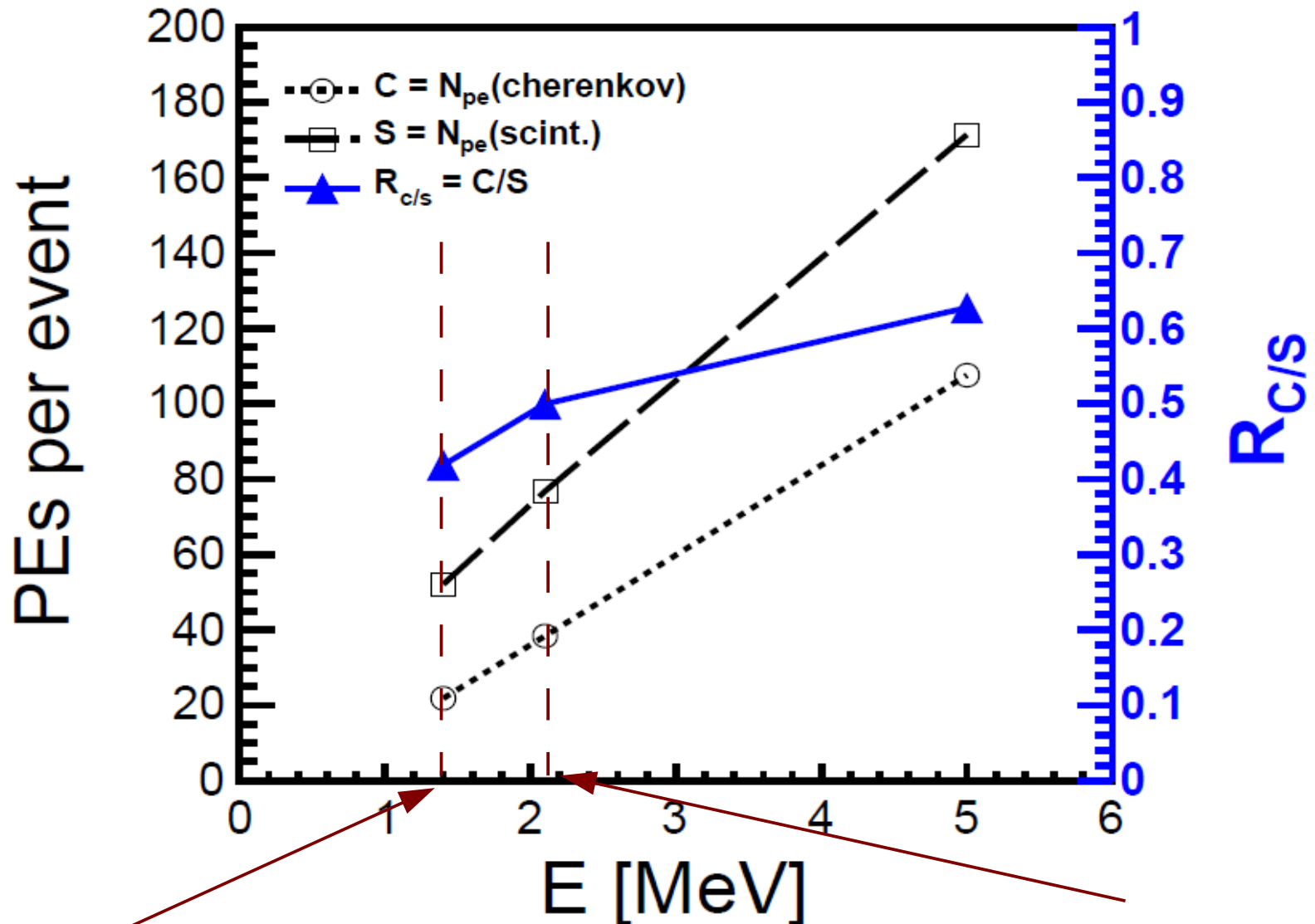
Large liquid-scintillator-based detectors have proven to be exceptionally effective for low energy neutrino measurements due to their good energy resolution and scalability to large volumes. The addition of directional information using Cherenkov light and fast timing would enhance the scientific reach of these detectors, especially for searches for neutrino-less double-beta decay. In this paper, we develop a technique for extracting particle direction using the difference in arrival times for Cherenkov and scintillation light, and evaluate several detector advances in timing, photodetector spectral response, and scintillator emission spectra that could be used to make direction reconstruction a reality in a kiloton-scale detector.



5 MeV electrons are very promising even with a very simple reconstruction

What about Lower Energies?

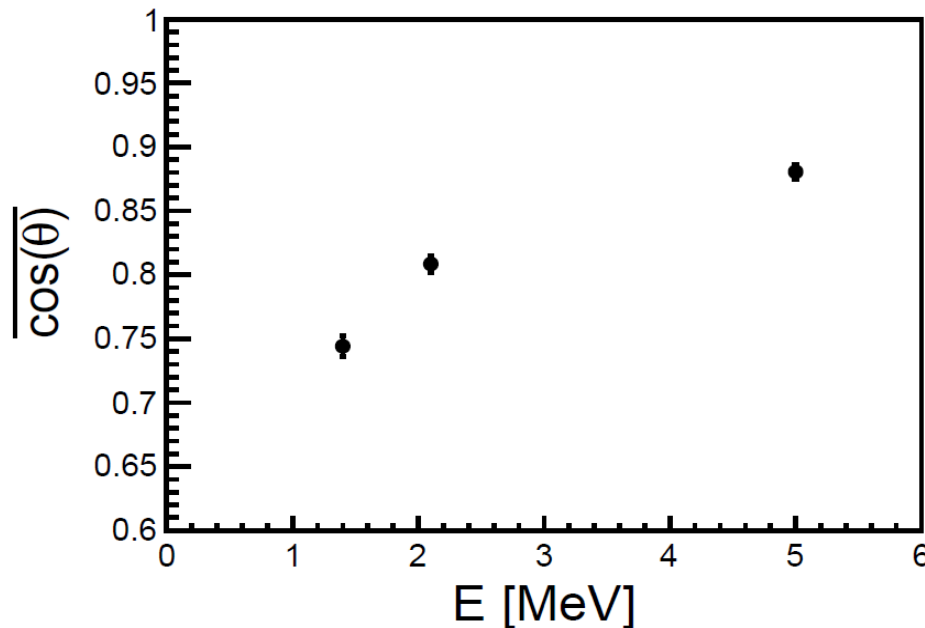
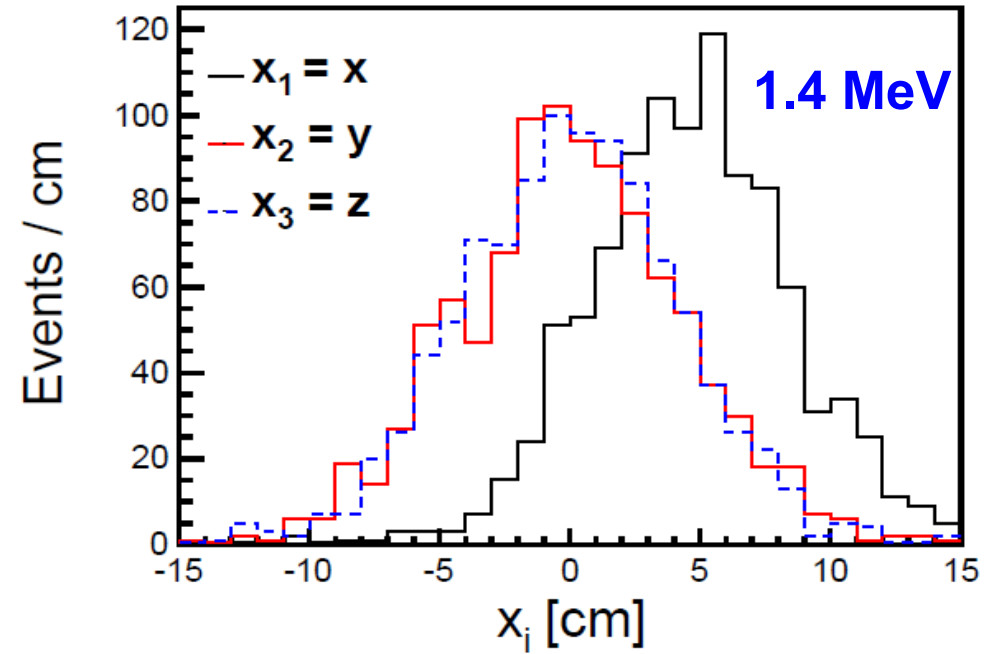
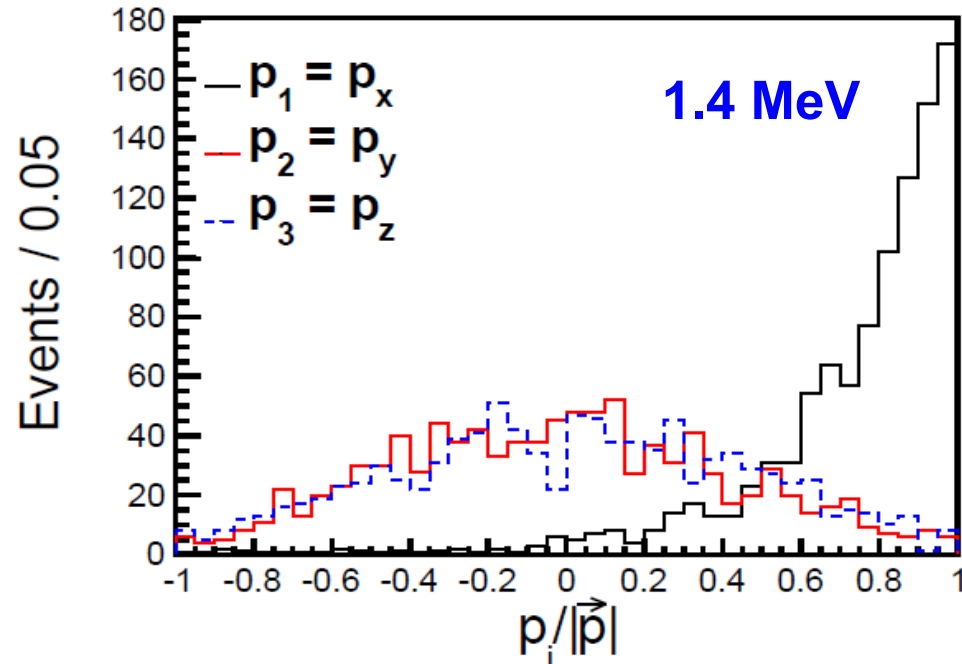
Light yield: Cherenkov vs scintillation



$\frac{1}{2} Q (^{116}\text{Cd}) = 1.4 \text{ MeV}$

$\frac{1}{2} Q (^{48}\text{Ca}) = 2.1 \text{ MeV}$

What about Lower Energies?

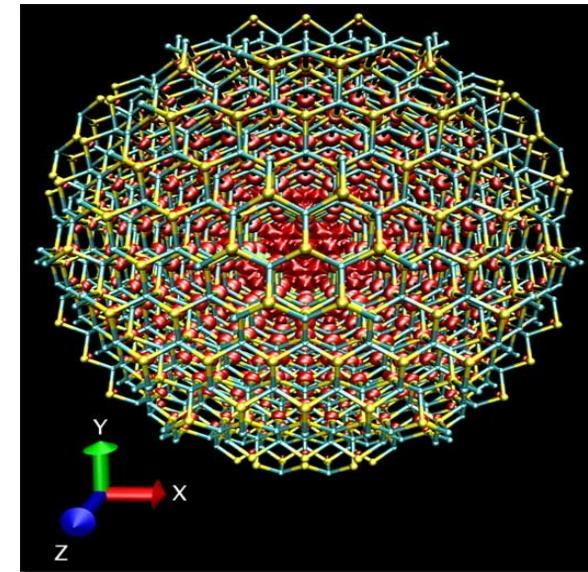


- *With 100ps timing, the vertex is constrained within 4-6 cm and the directional information can be extracted even for ~ 1 MeV electrons*
- *The next major step is to compare „back-to-back“ $0\nu\beta\beta$ events with $2\nu\beta\beta$ background*

Candidate Isotopes for $0\nu\beta\beta$

Isotope	Endpoint	Abundance
^{48}Ca	4.271 MeV	0.0035%
^{150}Nd	3.367 MeV	5.6%
^{96}Zr	3.350 MeV	2.8%
^{100}Mo	3.034 MeV	9.6%
^{82}Se	2.995 MeV	9.2%
^{116}Cd	2.802 MeV	7.5%
^{130}Te	2.533 MeV	34.5%
^{136}Xe	2.479 MeV	8.9%
^{76}Ge	2.039 MeV	7.8%
^{128}Te	0.868 MeV	31.7%

Quantum dot doped scintillators



Common materials are CdS, CdSe, CdTe

Work by UCLA group
C.Aberle, J.J.Li, S.Weiss, and
L.Winslow

This is another lunch topic!

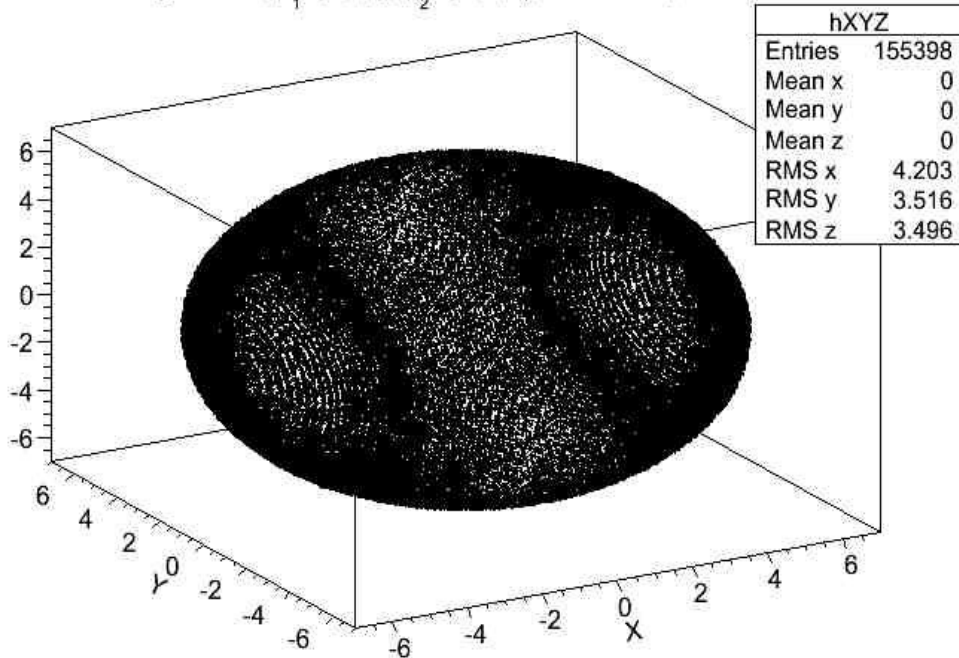
Advantage of quantum dot doping:

- Narrow the scintillation spectrum
- Shift scintillation spectrum to shorter wavelength
- Dope with metals which can undergo $0\nu\beta\beta$

Event Topology

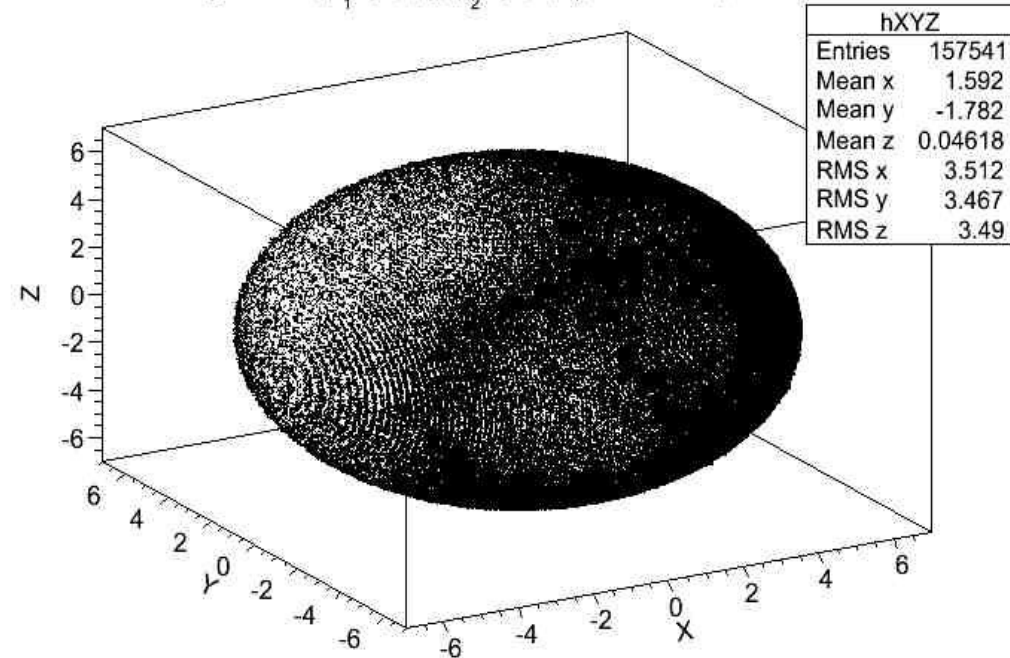
100 „signal-like“ events
(5MeV electrons back-to-back)

Light=CHE, $\vec{p}_1=(1,0,0)$, $\vec{p}_2=(-1,0,0)$, QE=100%, $t<1000\text{ns}$



100 „bkg-like“ events
(5MeV electrons at 90 degree)

Light=CHE, $\vec{p}_1=(1,0,0)$, $\vec{p}_2=(0,-1,0)$, QE=100%, $t<1000\text{ns}$



Cherenkov light only, no time cut, 100% light collection

Spherical Harmonics

Real-value basis:

$$Y_{\ell m} = \begin{cases} \frac{1}{\sqrt{2}} (Y_{\ell}^m + (-1)^m Y_{\ell}^{-m}) = \sqrt{2} N_{(\ell, m)} P_{\ell}^m(\cos \theta) \cos m\varphi & \text{if } m > 0 \\ Y_{\ell}^0 & \text{if } m = 0 \\ \frac{1}{i\sqrt{2}} (Y_{\ell}^{-m} - (-1)^m Y_{\ell}^m) = \sqrt{2} N_{(\ell, |m|)} P_{\ell}^{|m|}(\cos \theta) \sin |m|\varphi & \text{if } m < 0. \end{cases}$$

$$f(\theta, \varphi) = \sum_{\ell=0}^{\infty} \sum_{m=-\ell}^{\ell} f_{\ell m} Y_{\ell m}(\theta, \varphi).$$

$$N_{(\ell, m)} \equiv \sqrt{\frac{(2\ell + 1)(\ell - m)!}{4\pi(\ell + m)!}}.$$

$$f_{\ell}^m = \int_{\Omega} f(\theta, \varphi) Y_{\ell}^{m*}(\theta, \varphi) d\Omega = \int_0^{2\pi} d\varphi \int_0^{\pi} d\theta \sin \theta f(\theta, \varphi) Y_{\ell}^{m*}(\theta, \varphi).$$

L2 norm

„Power“ (rotation invariant)

$$\int_{\Omega} |f(\Omega)|^2 d\Omega = \sum_{\ell=0}^{\infty} S_{ff}(\ell)$$

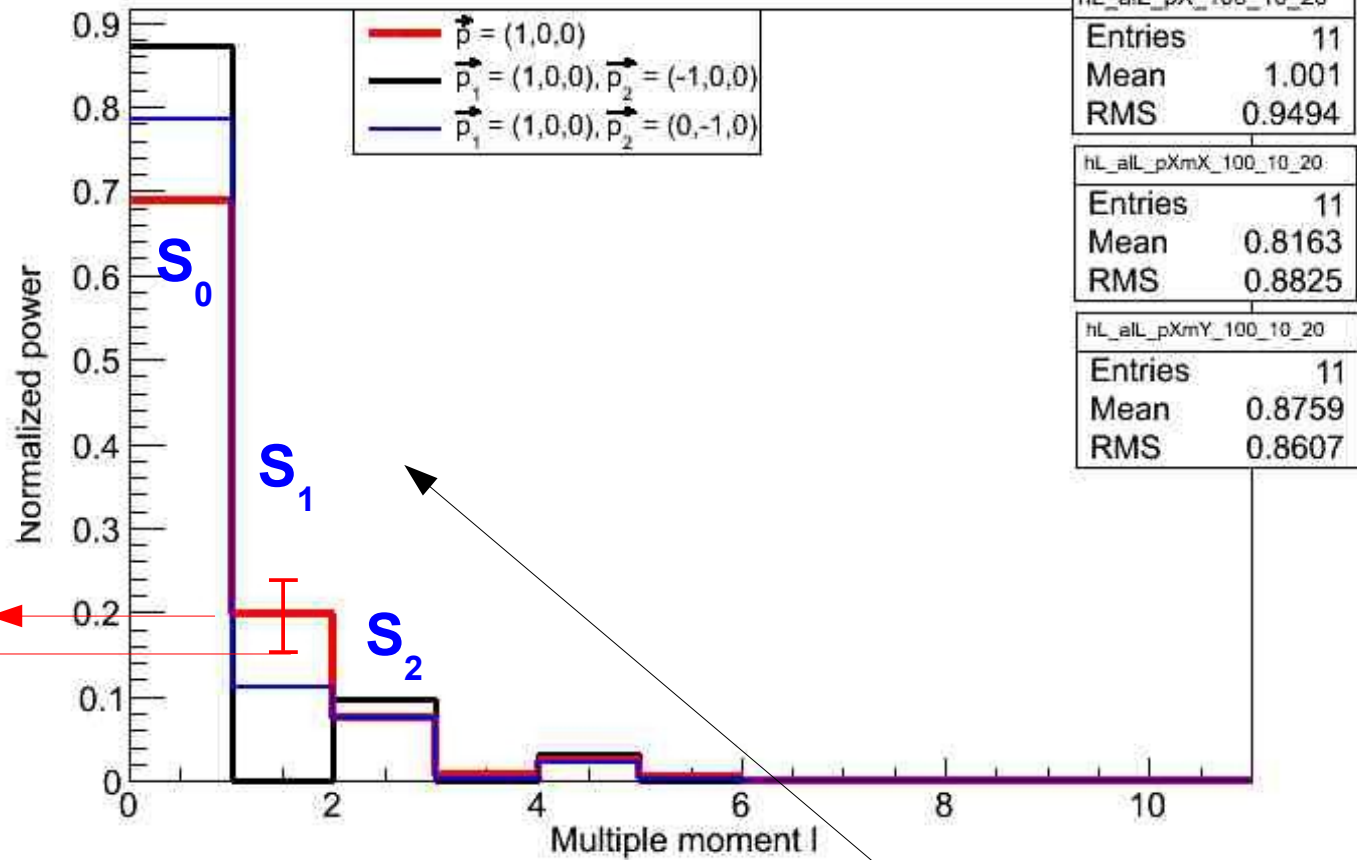
$$S_{ff}(\ell) = \sum_{m=-\ell}^{\ell} |f_{\ell m}|^2$$

Spherical Harmonics for back-to-back vs 90 degree topologies

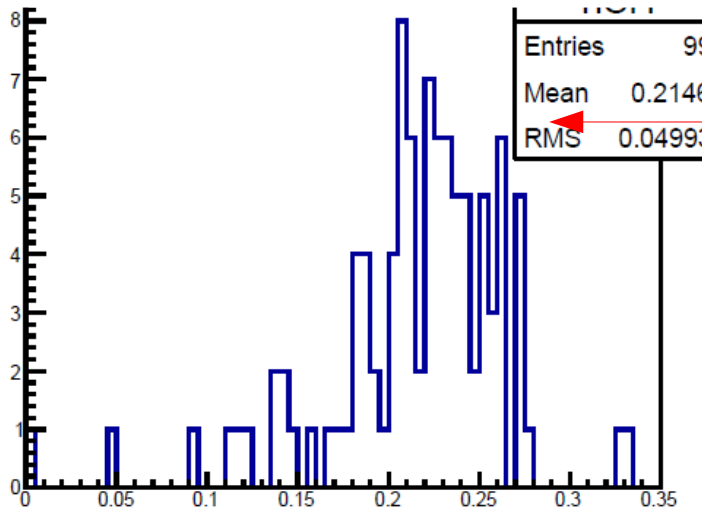
- Look for the difference between **black (back-to-back)** and **blue (90-degree)** lines.
- **Red line** is for comparison with single electron events

Light=ALL, QE=100%, t<34ns

$$S_{ff}(\ell) = \sum_{m=-\ell}^{\ell} |f_{\ell m}|^2$$



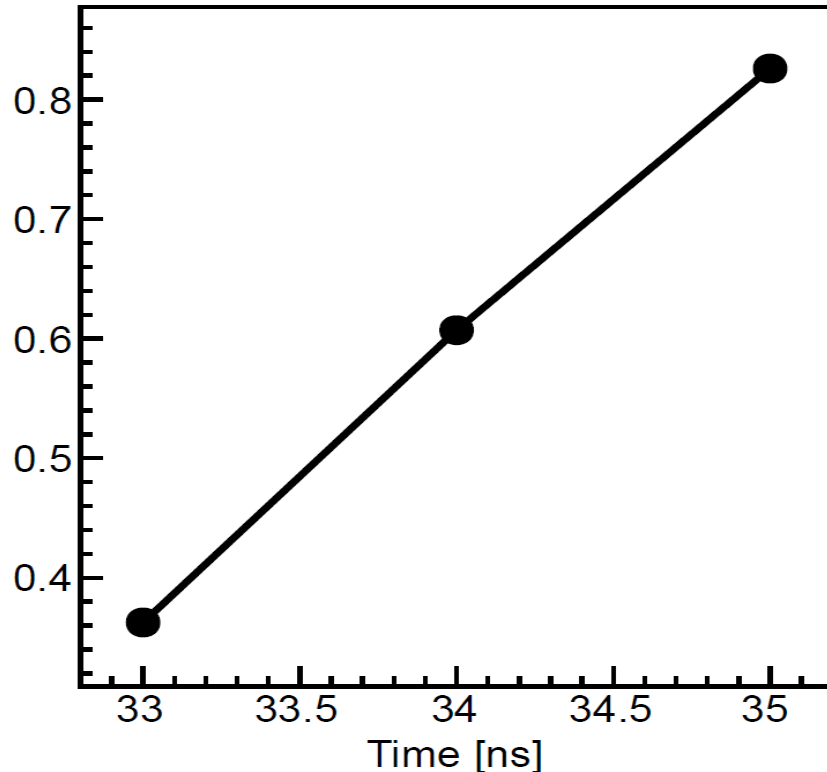
S_1 event-by-event fluctuations



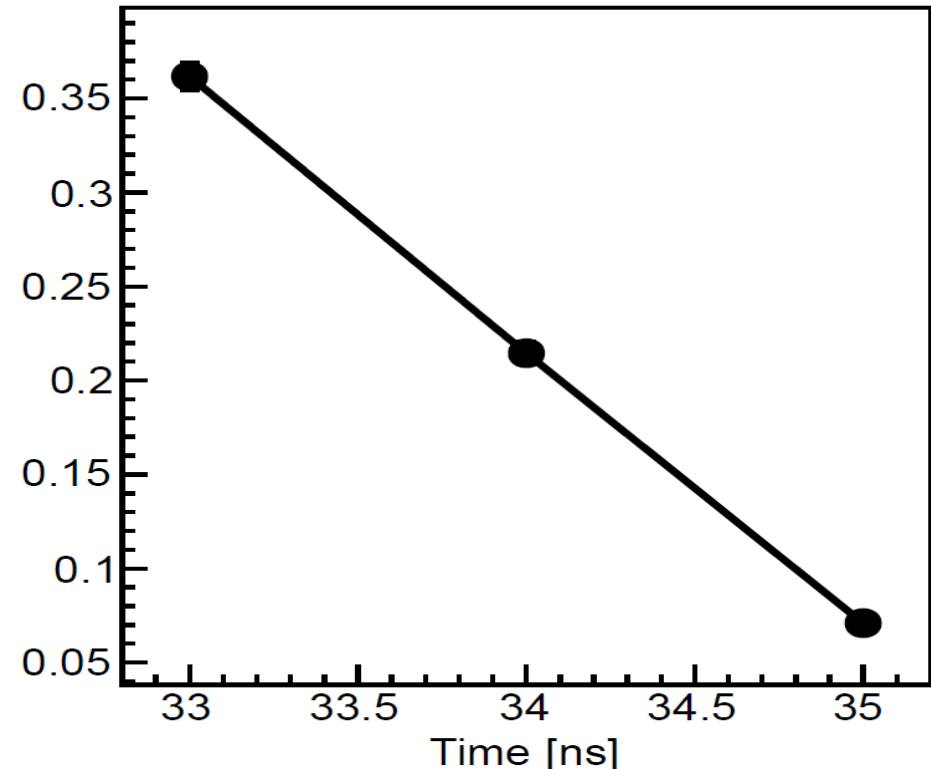
These S_l are calculated for combined hit distribution of 100 events

S_1 as function of time cut

Mean of S_0



Mean of S_1

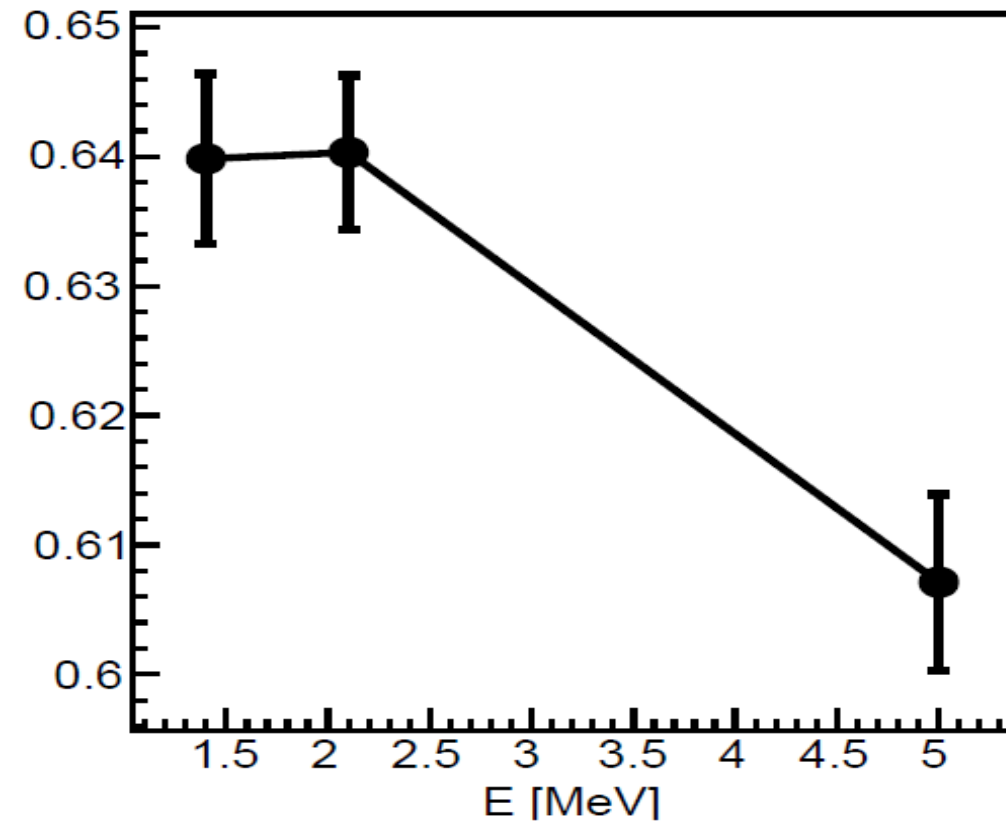


I like this strong dependence on time because

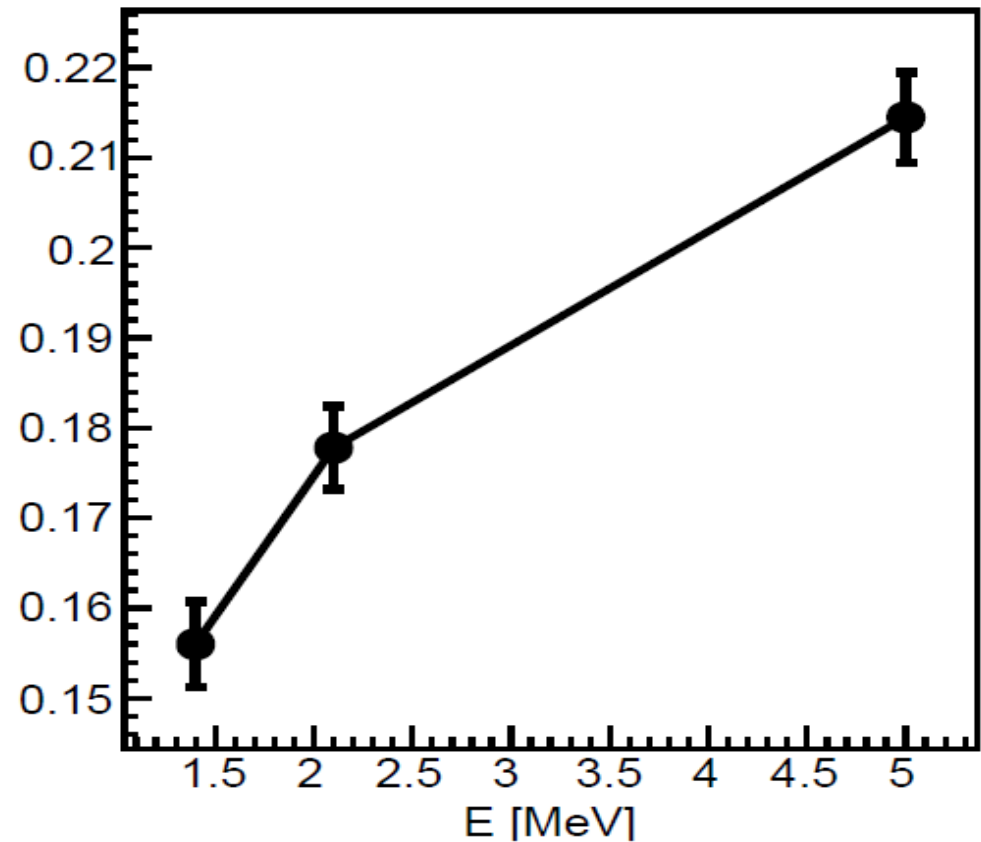
with fast photo-detectors we may be able to follow the time evolution of spherical harmonics and use this information in reconstruction

S_1 as function of electron energy

Mean of S_0



Mean of S_1



Note Y axis range, the dependence isn't too strong

What's Next?

- *Compare simulated $0\nu\beta\beta$ events with $2\nu\beta\beta$*
- *Consider events off center*

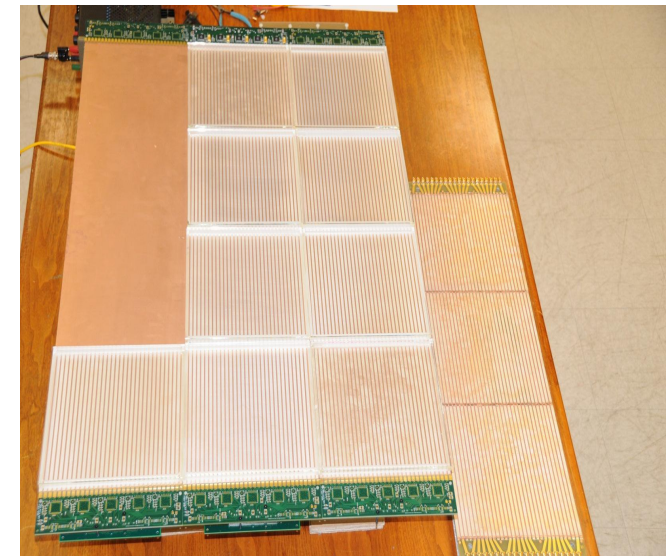
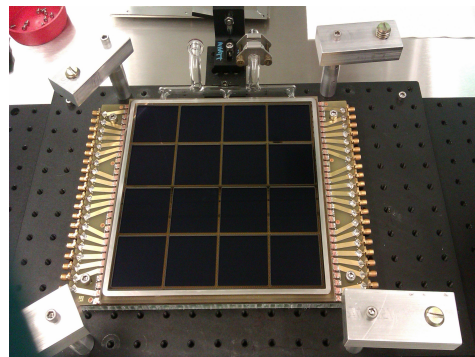
Now, I will talk about LAPPD and Hermetic Packaging

Large Area Picosecond Photo Detectors

Transformational Change

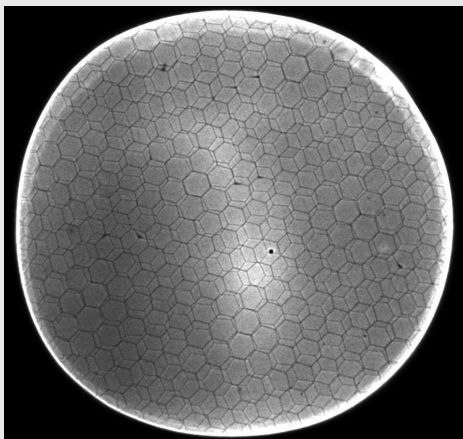


- *Large area*
- *Fast timing*
- *Inexpensive*

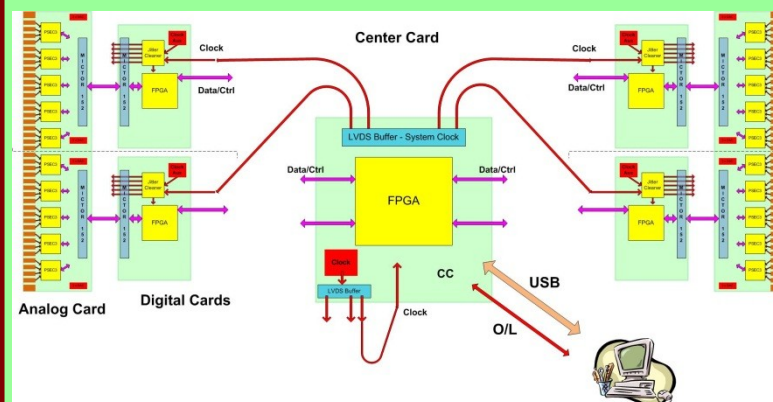


LAPPD Components

MicroChannel Plates



Electronics/Integration



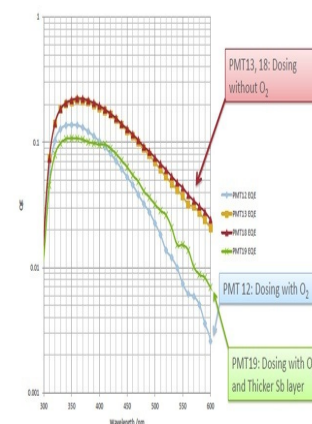
Hermetic Packaging



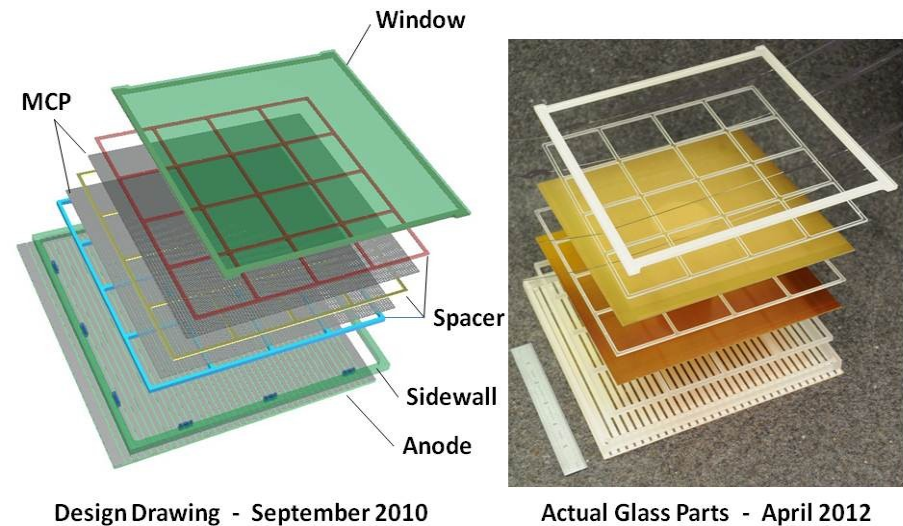
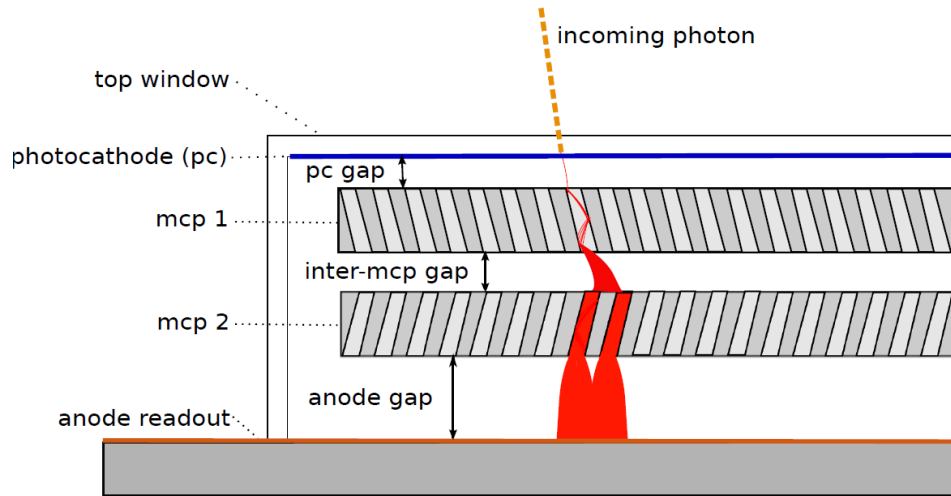
Photocathodes



Summary of cathodes grown by Burle Equip

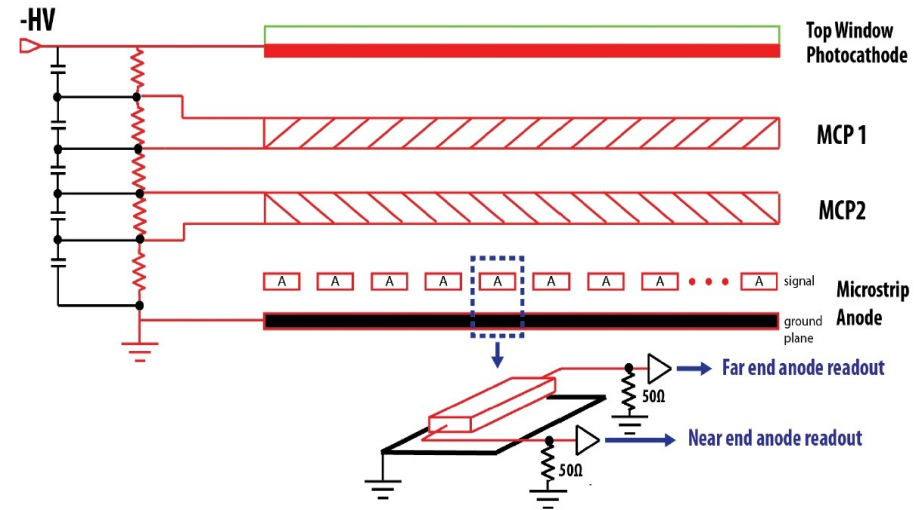


Glass Package (20x20cm²)

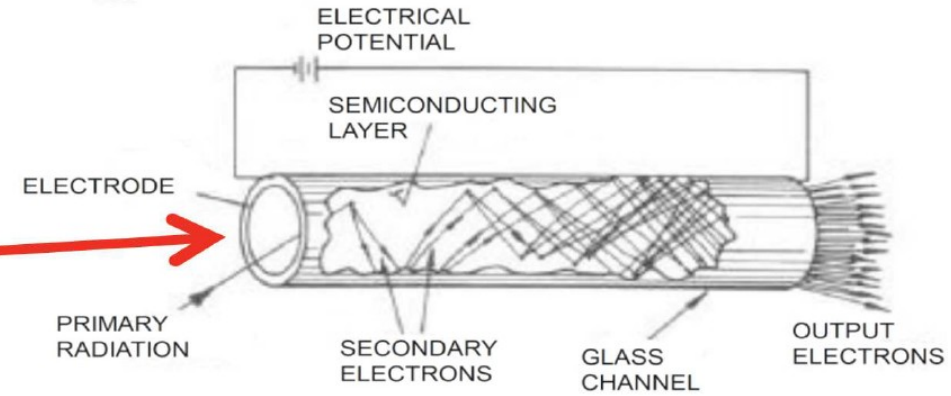
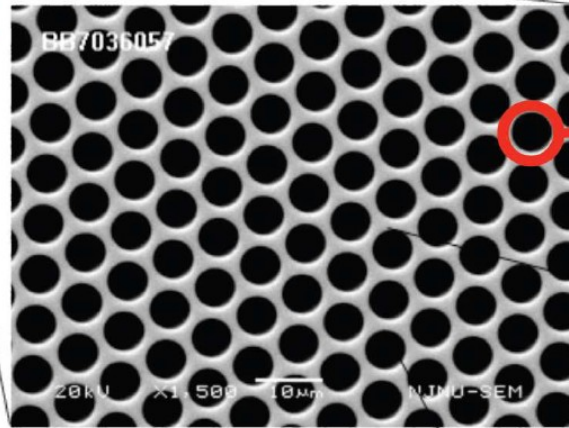


- Cheap, widely available float glass
- Anode is made by silk-screening
- Flat panel
- No pins, single HV cable
- Modular design
- High bandwidth 50 Ω object - designed for fast timing

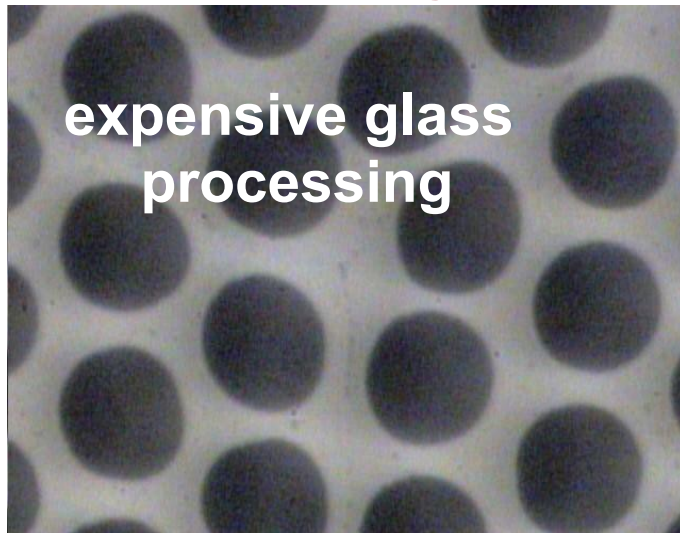
The Frugal Tile



MCP Fundamentals



Conventional Pb-glass MCP

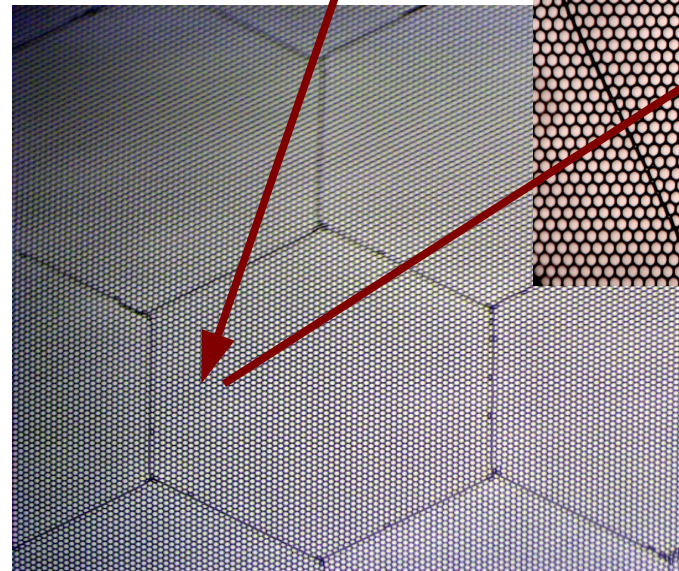
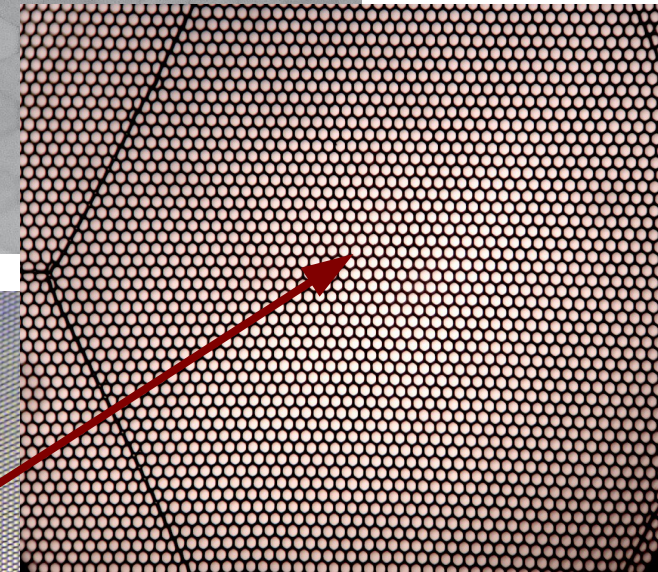
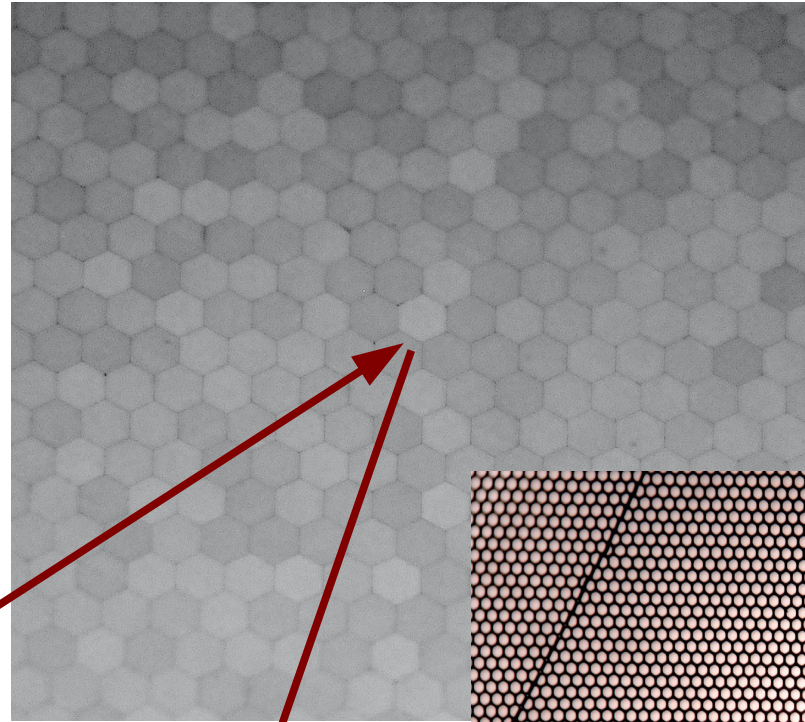
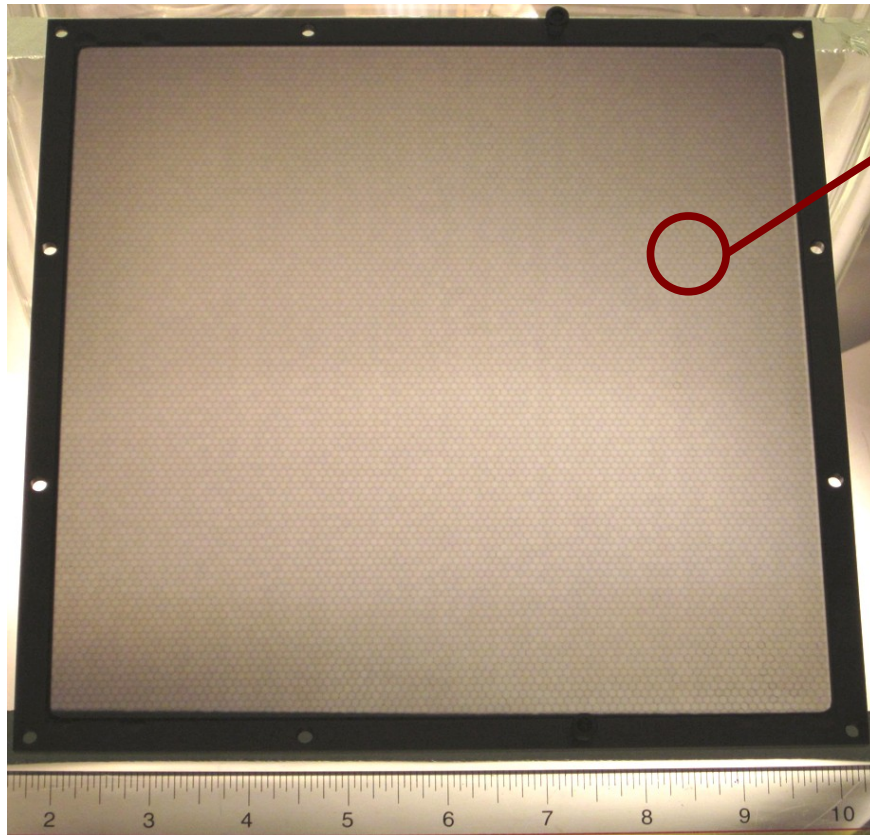


Incom glass substrate



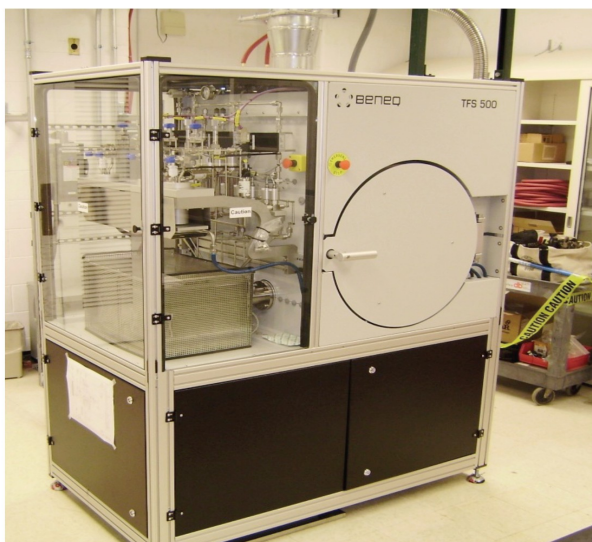
Micro-Capillary Arrays by Incom

- **Material:** borosilicate glass
- **Area:** 20x20cm²
- **Thickness:** 1.2mm
- **Pore size:** 20 μm
- **L/D ratio:** 60:1
- **Open area:** 65-83%

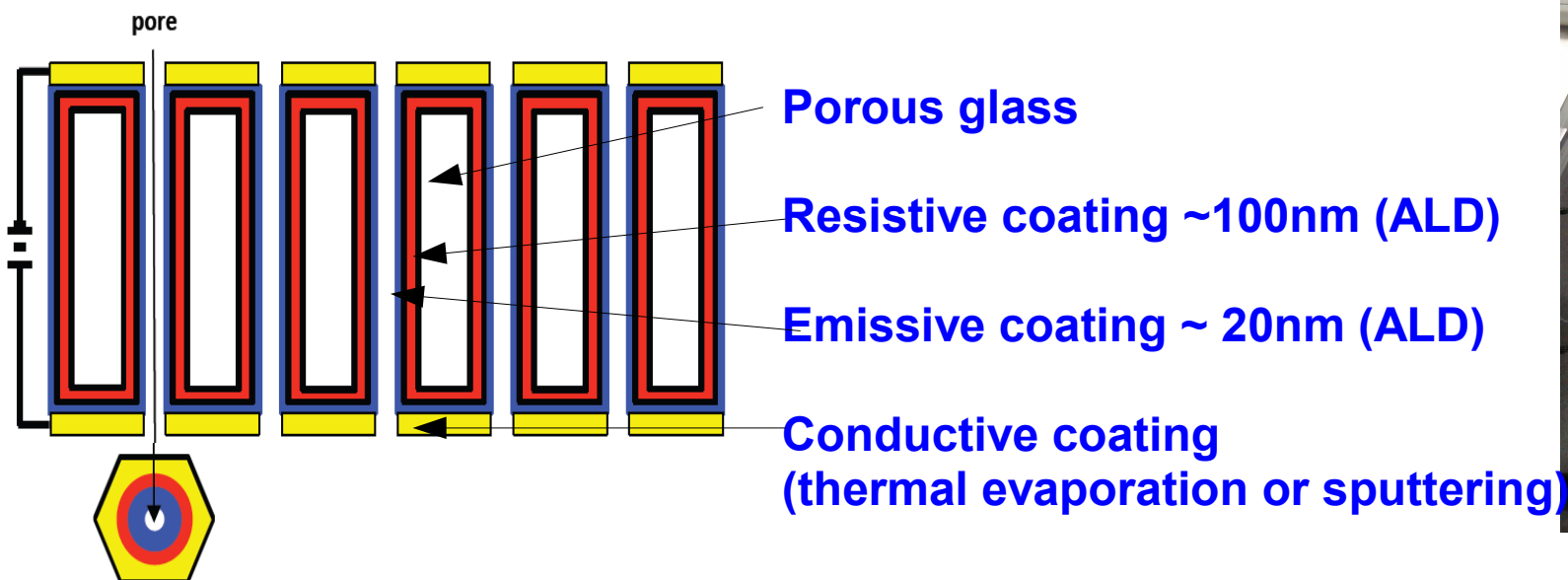
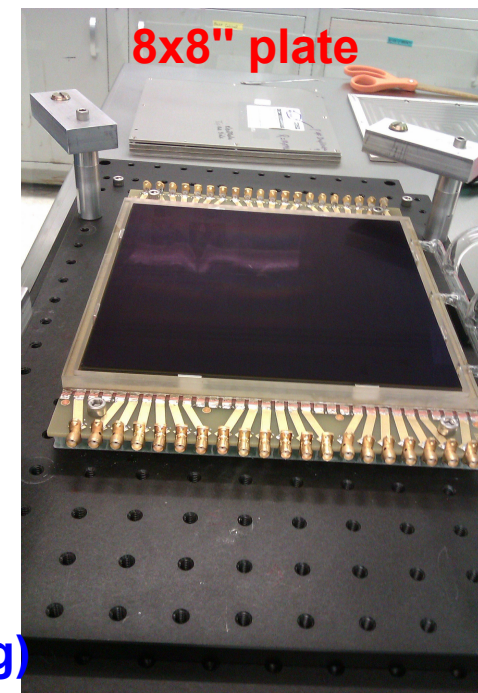
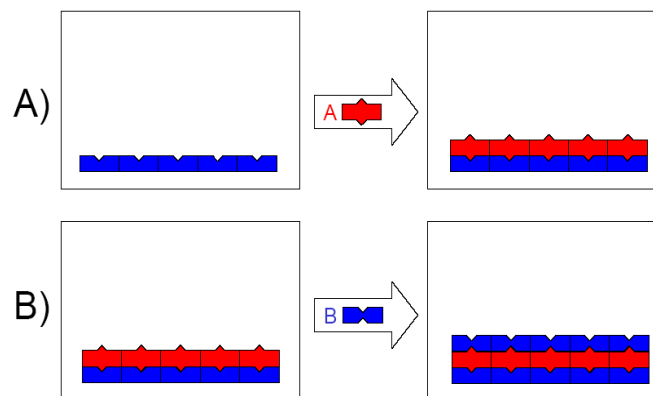


MCP by Atomic Layer Deposition (ALD)

Beneq reactor for ALD
@Argonne National Laboratory

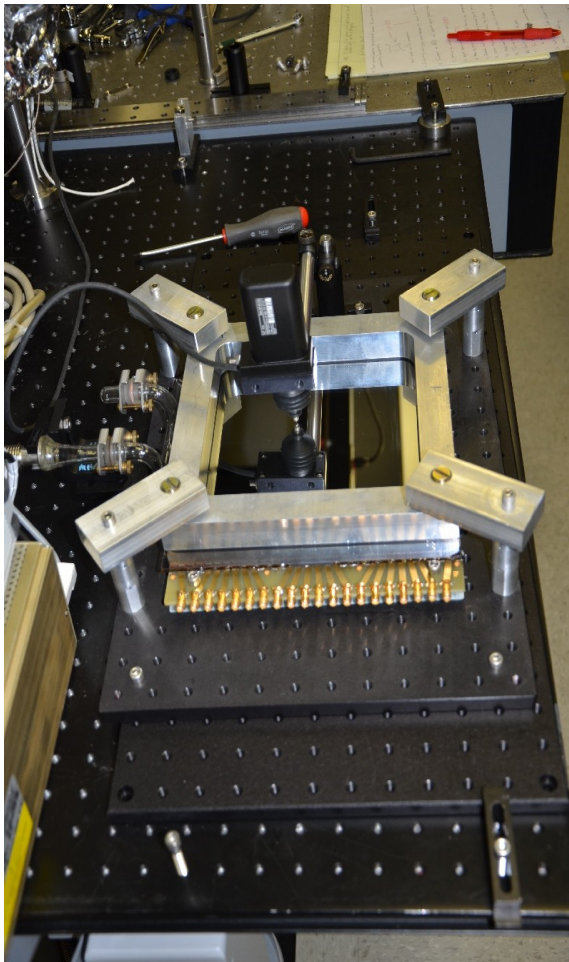


*ALD Process for MCP Coating
Developed by
A.Mane, J.Elam*

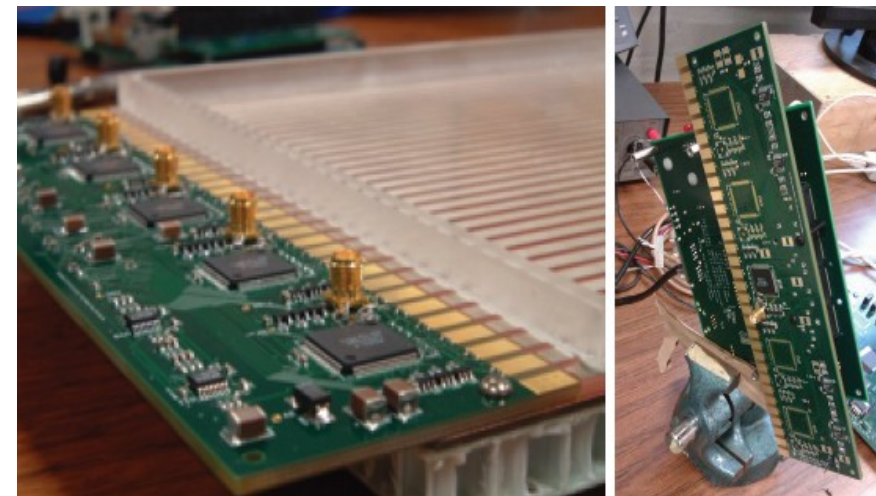
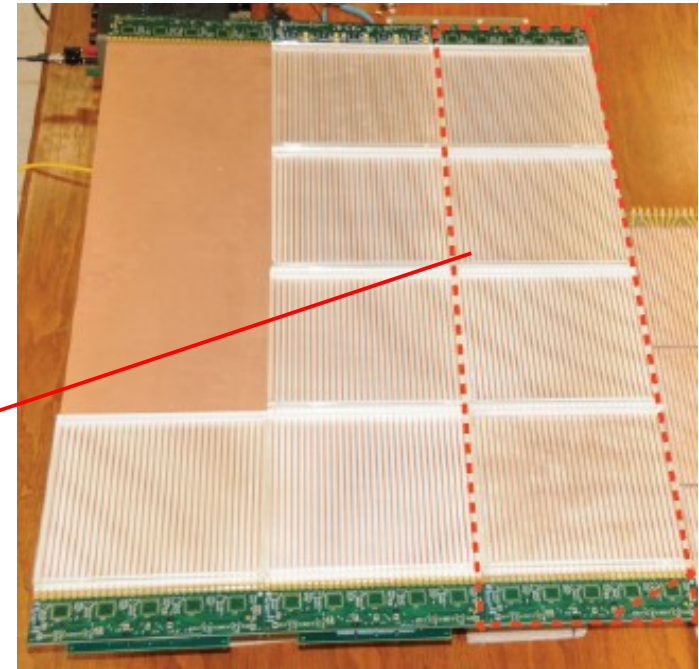
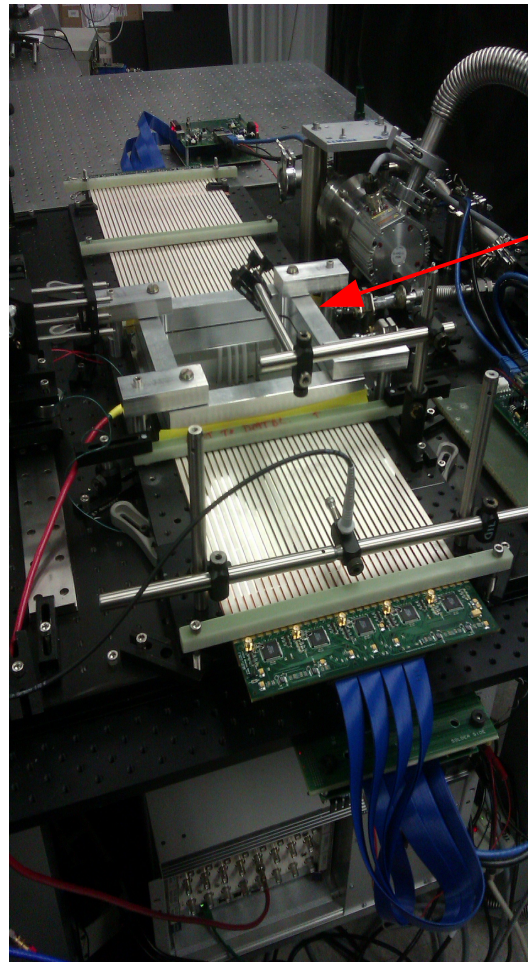


System Integration: "Demountable"

Demountable 1.0
(May 2012)



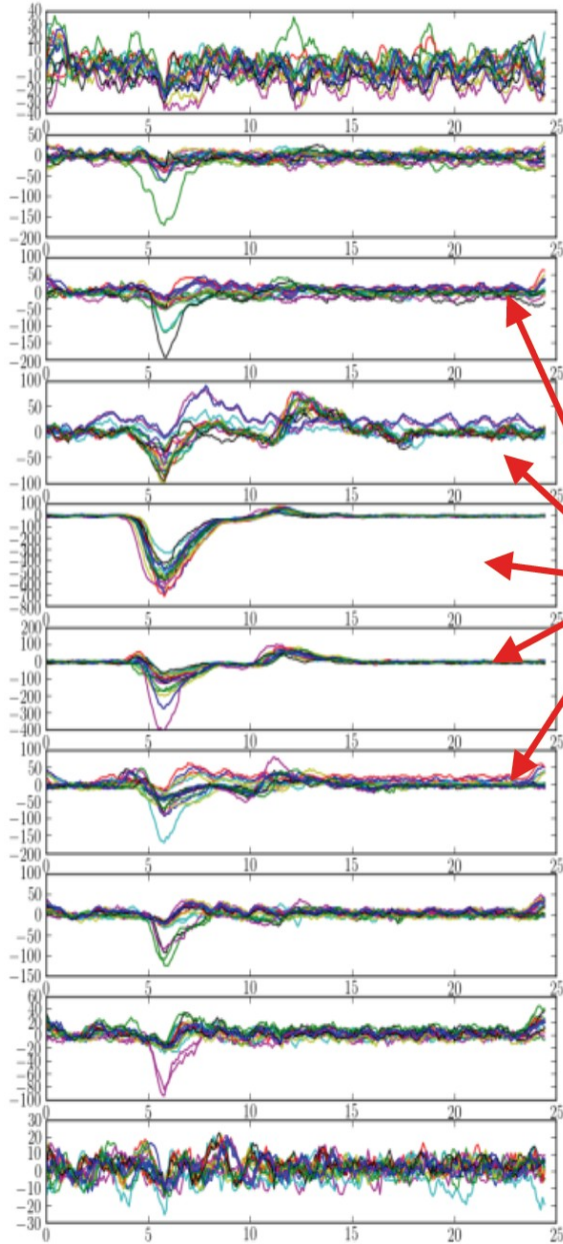
Demountable 3.0
(Sep-Dec 2012)



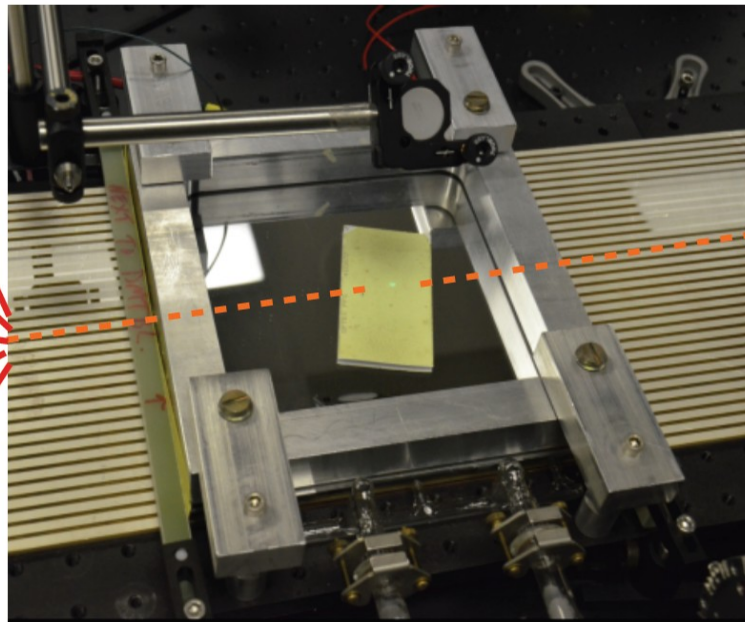
Performance

Pulses on 10 striplines

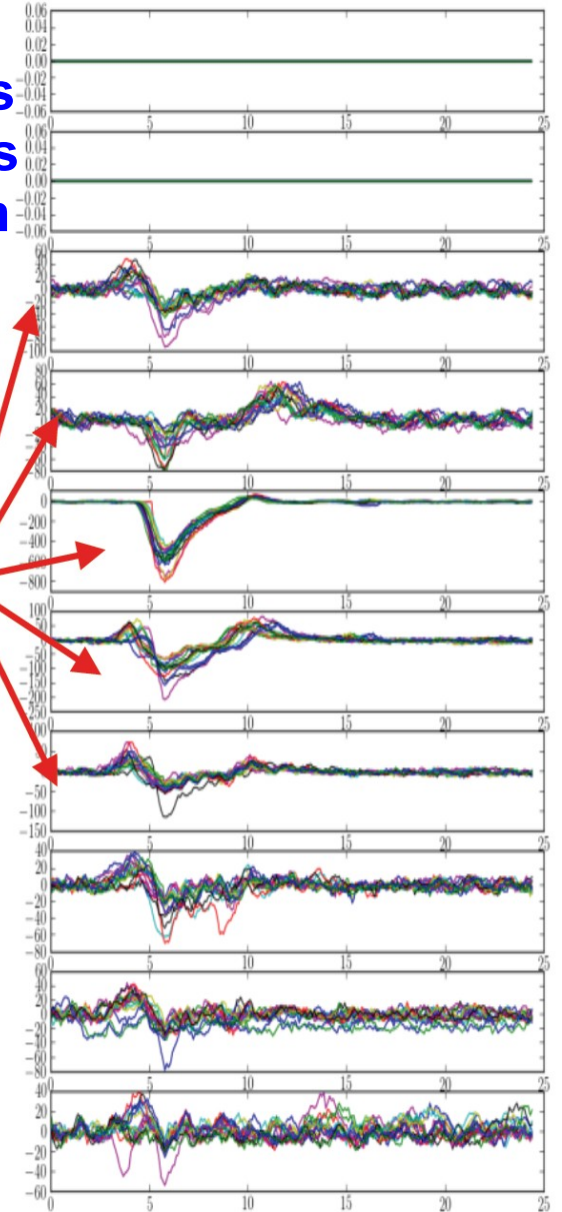
Left Side



Differential time resolution up to $\sim 5\text{ps}$
Time-of-flight resolution $\sim 60\text{ps}$
Spatial resolution $< 1\text{mm}$



90cm long anode!



Right Side
Pulses on 10 striplines

Rev. Sci. Instrum. 84, 061301 (2013)

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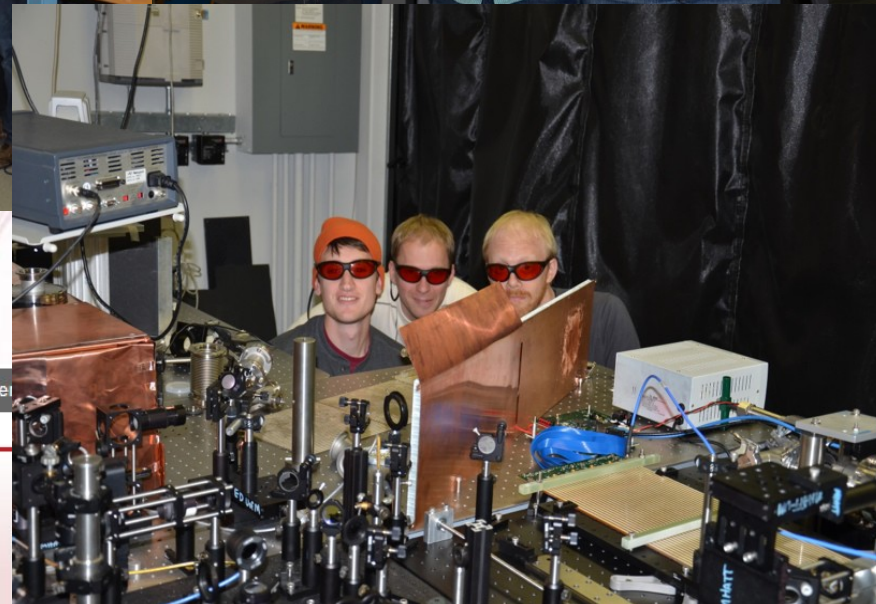
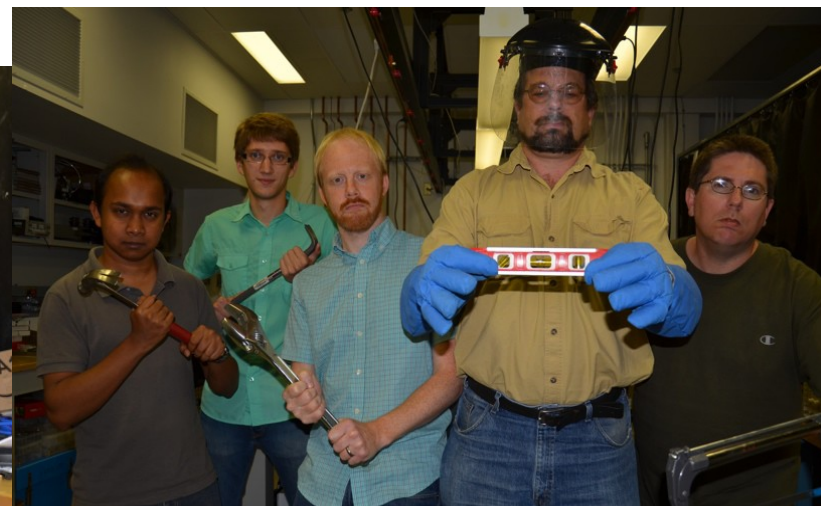
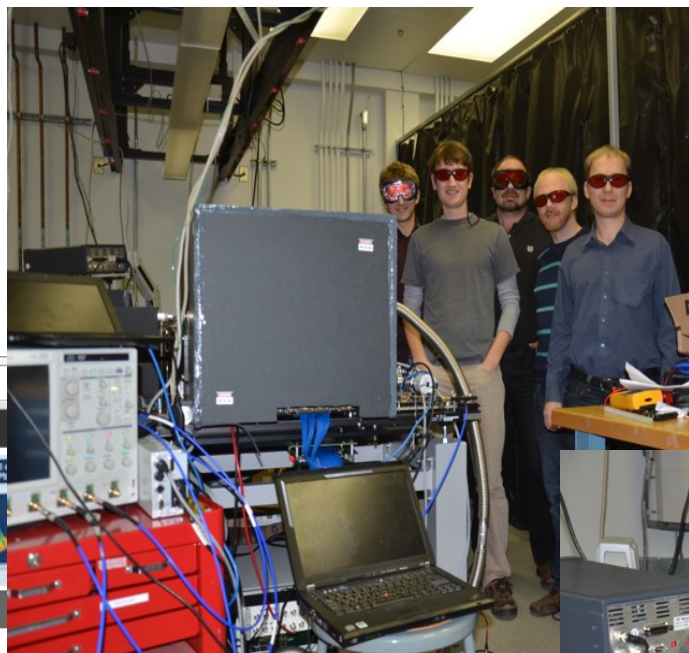
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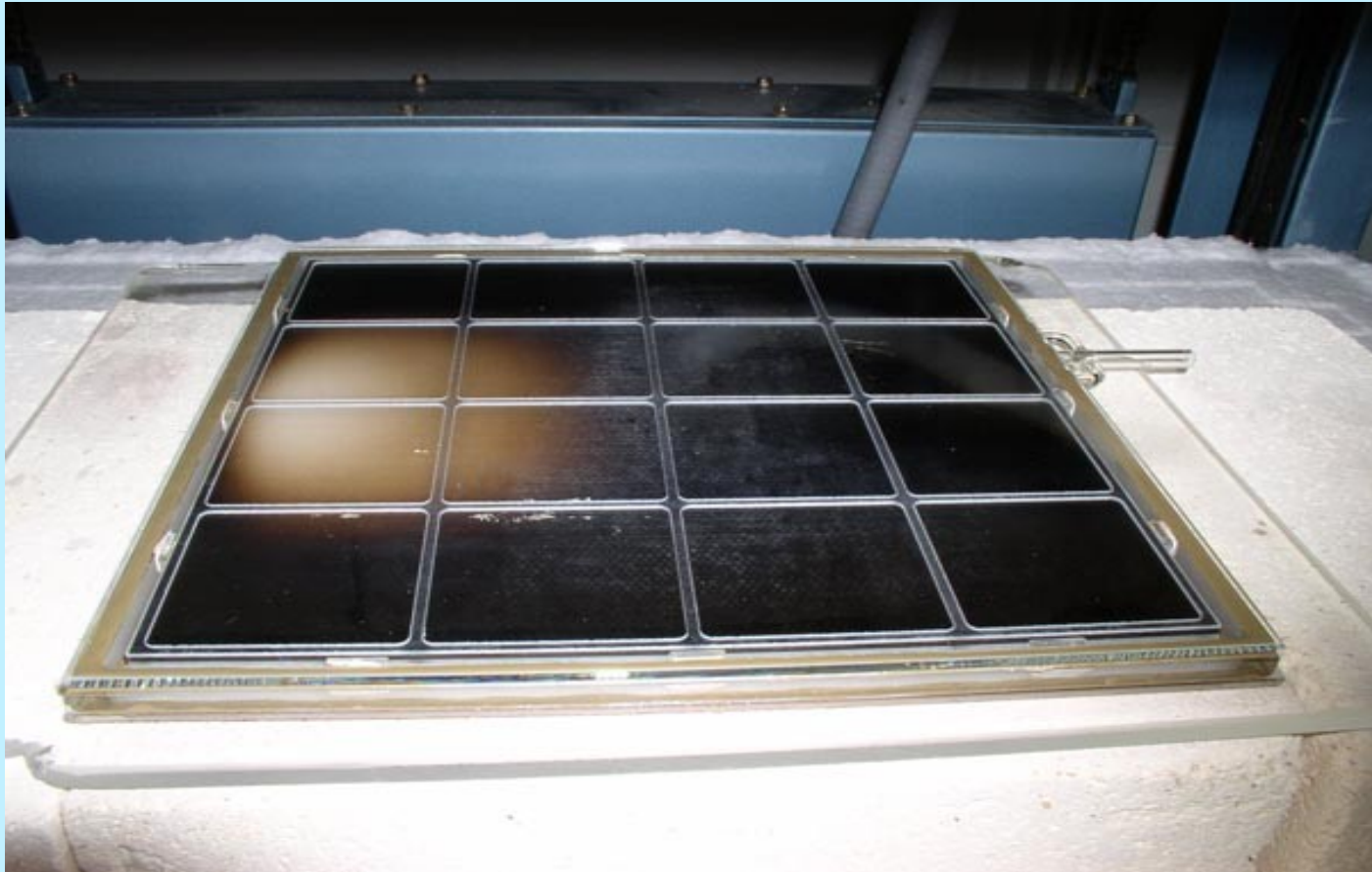
Research Highlights



Invited Article: Coherent imaging using seeded free-electron laser pulses with variable polarization: First results and research opportunities

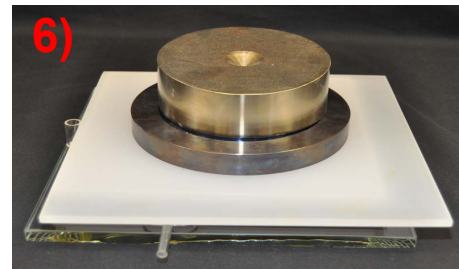
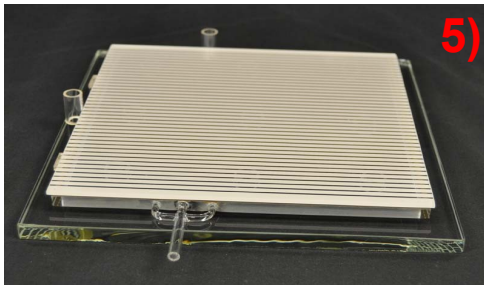
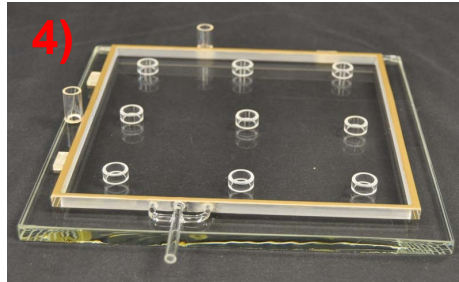
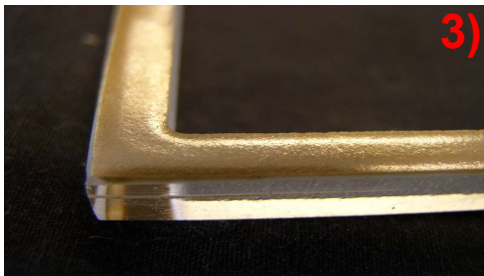
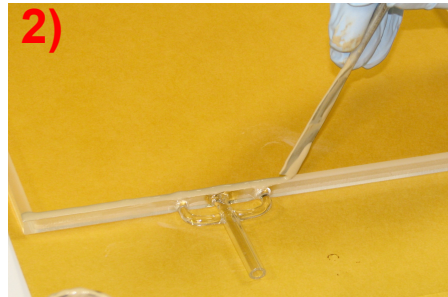
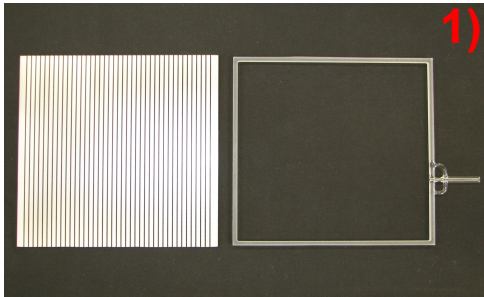


Hermetic Packaging



Frit Seal

J.Gregar, M.Minot



1) Attach pump out tube to 8.66x8.66" frame

2) Apply schott #G018-223 K3 frit paste to frame

3) Fire the frit
(many trials to optimize parameters)

4) Prepare for anode plate frit sealing

5) Position anode on top of the frame

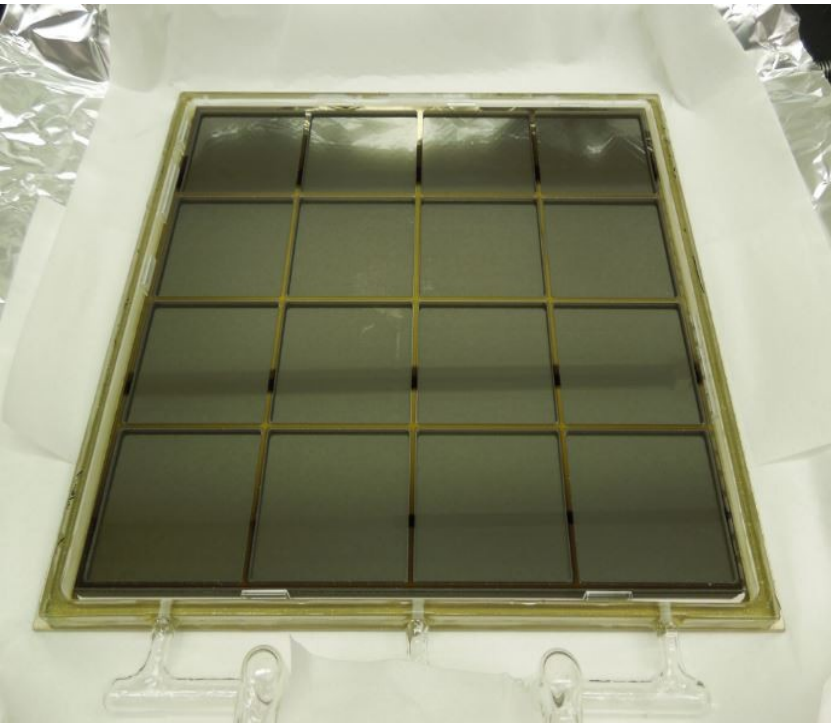
6) Add weight

- **Tile bases are reliably reproducible**
- **Mechanical and vacuum properties have been tested**



Top Seal

How to close frit sealed tile base at the top and stay at moderate temperatures? **"Top Seal" problem**



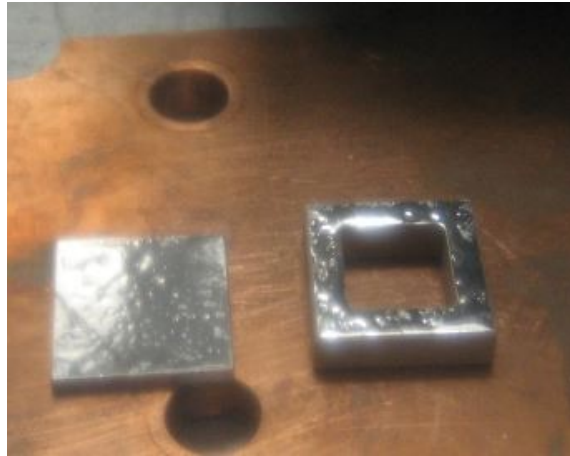
Use indium or indium alloys

- soft metal
- low melting point
(157C for pure In, 72C for InBi)
- essentially zero vapor pressure
- **indium-glass seals are successfully used by industry**

Parallel efforts: **„Hot Seal“** and „Cold Seal“ (or “Compression Seal“)

Hot Seal

Step 1:
*apply melted indium
onto the glass*



Step 3:
*pump from open side
and leak check*



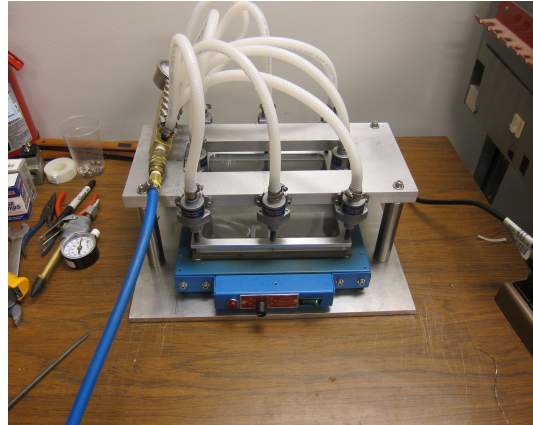
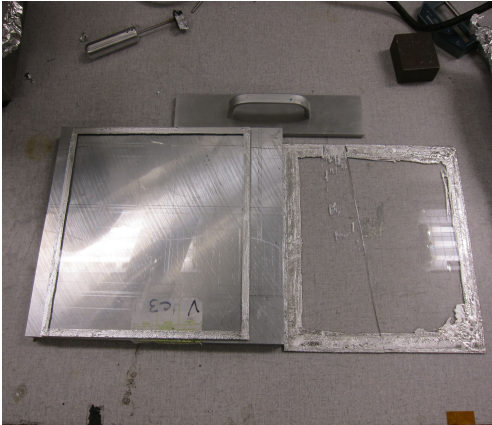
Step 2:
bring parts into contact and press



Prerequisites for leak-tight seal

- **Strong Indium-Glass bond**
- **Strong Indium-Indium bond**

Phase I (in air)



A.E., Henry Frisch, Mary Heintz,
Bob Metz, Richard Northrop,
Razib Obaid

- **Indium seal fundamentals**
 - interface (good adhesion of indium to the glass surface)
 - oxide formation
- **Proof of principle using 1x1" test samples**
 - little oxidation (assembly is fast)
 - many successful **reproducible** leak tight samples
- **Several (>10) attempts to make 8x8" seal**
 - oxide formation becomes limiting factor (slow assembly)
 - best result is a part with 10^{-6} cc/s leak at a single pinpoint

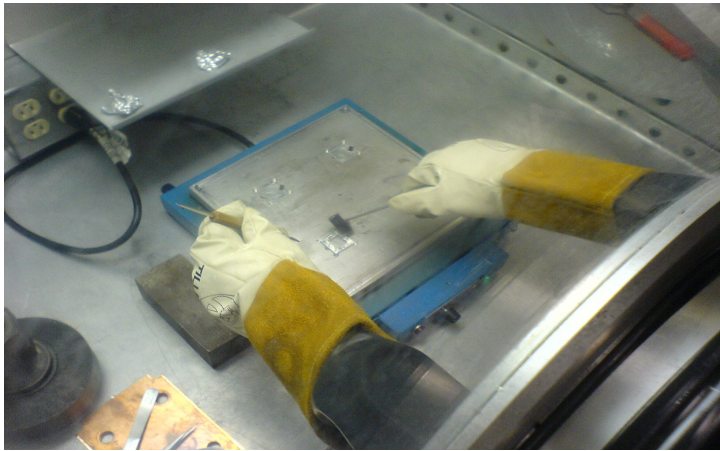
...indium oxidizes quickly...

Phase II (in inert atmosphere)

Phase IIa (in inert atmosphere)

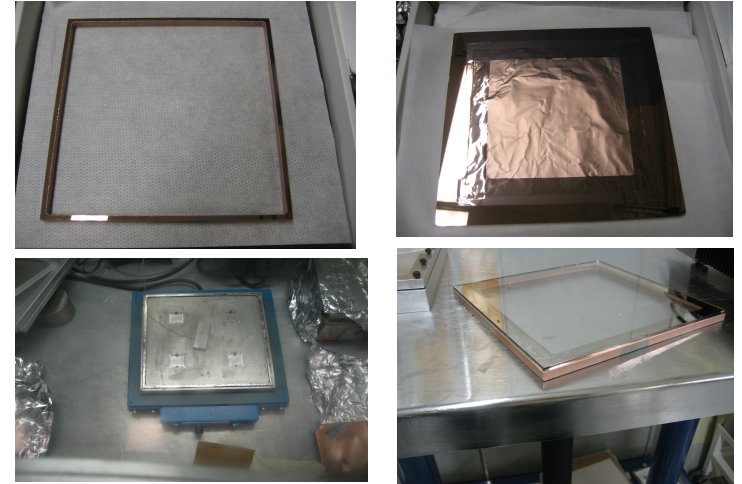
Nitrogen filled glove box:

O_2 and H_2O concentration ~ 5 ppm



...indium doesn't stick to glass if no O_2 ...

Phase IIb (add NiCr-Cu layer)



Borrowed from SSL seal
(200nm of NiCr+Cu)

Known facts:

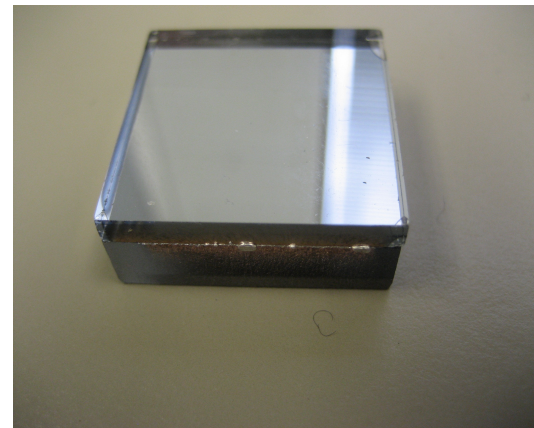
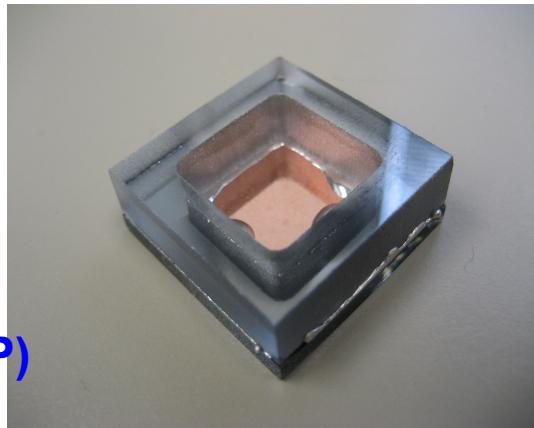
- **Indium wets copper surface**
 - Alloy is formed at the interface over time
- **NiCr interface to glass is essential**
 - NiCr is a good match to glass in terms of thermal coefficient
 - Cu would not stick to bare glass but³⁵ does so on NiCr

Testing NiCr-Cu-In Interface

Total 8 small size seals made: 5 are leak tight

3 have leaks (oxidation of Cu surface or electroding peeling off the glass)

NiCr-Cu coating of
1" samples done by
D. Walters (ANL)
M. Kupfer (UIC)
J. Williams (ANL-HEP)
C. Liu (ANL-APS)
Q. Guo (UChicago MRSEC)



Shear tests results

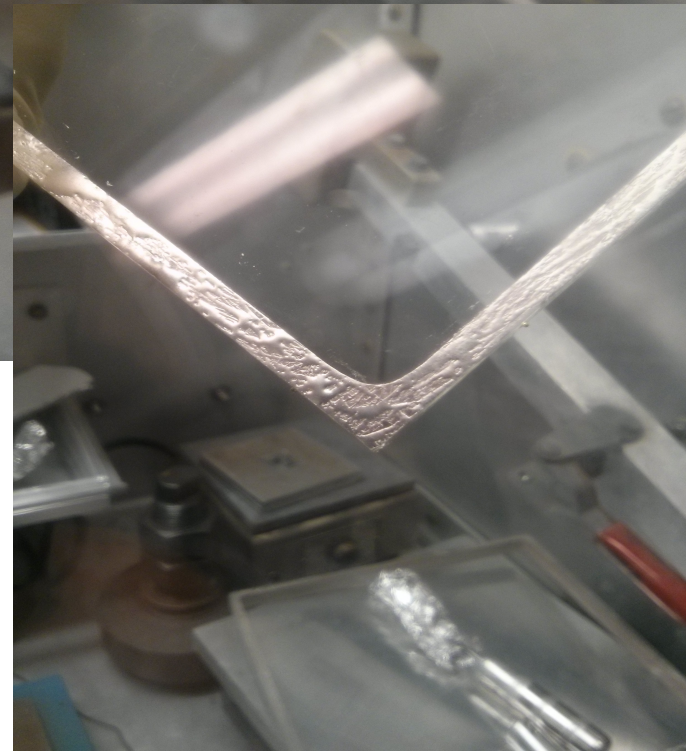
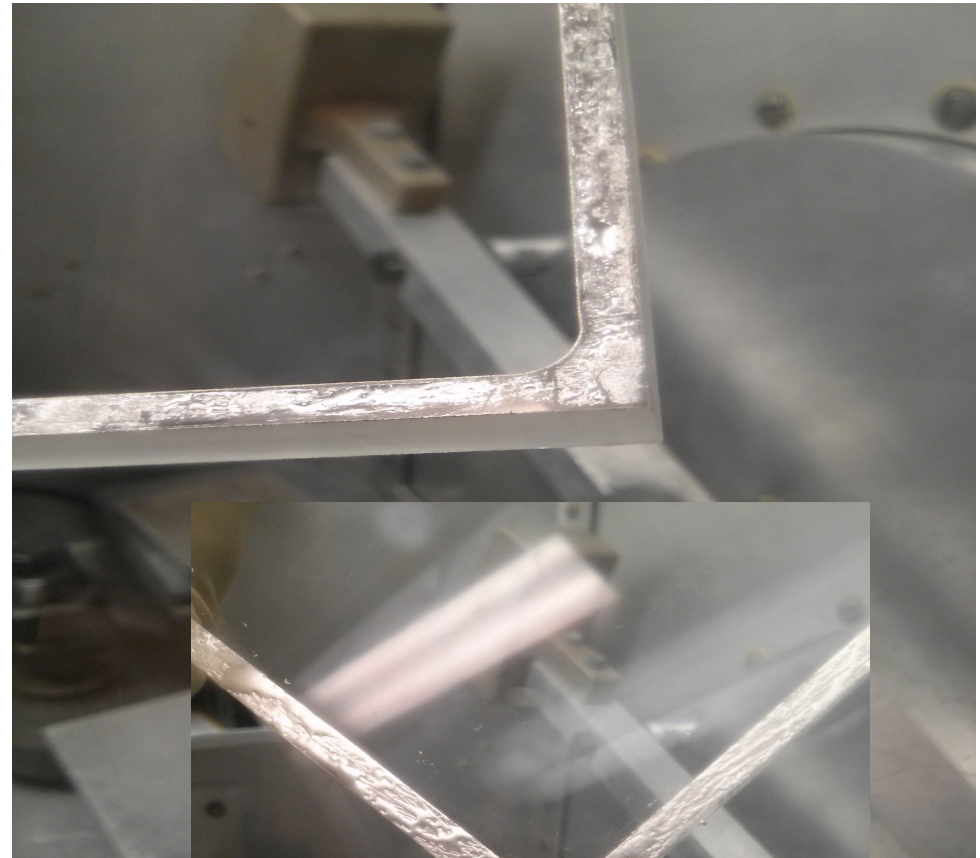
Leak tight samples:

Bare glass #1	190 lbs
Bare glass #2	278 lbs
Bare glass with groove	268 lbs
Cu coated glass #3	390 lbs
Cu coated glass #4	345 lbs

Samples with a leak:

Bare glass #4	47 lbs
Cu coated glass #1	213 lbs
Cu coated glass #2	221 lbs

Sealing 8x8" parts



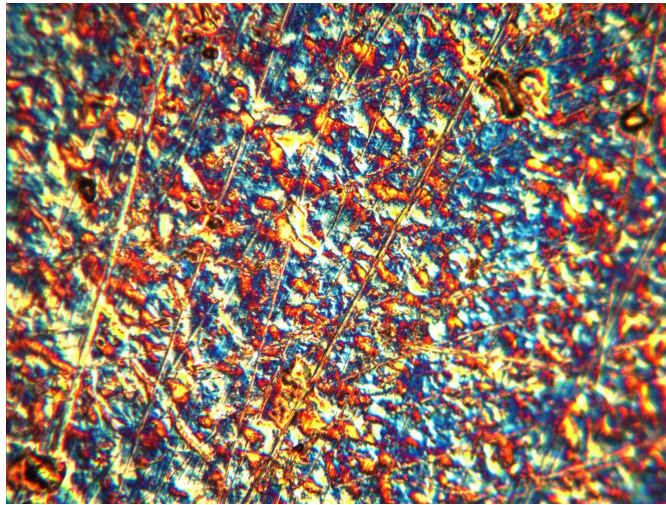
Testing NiCr-Cu-In Interface



Coating goes away if re-heated to 180C

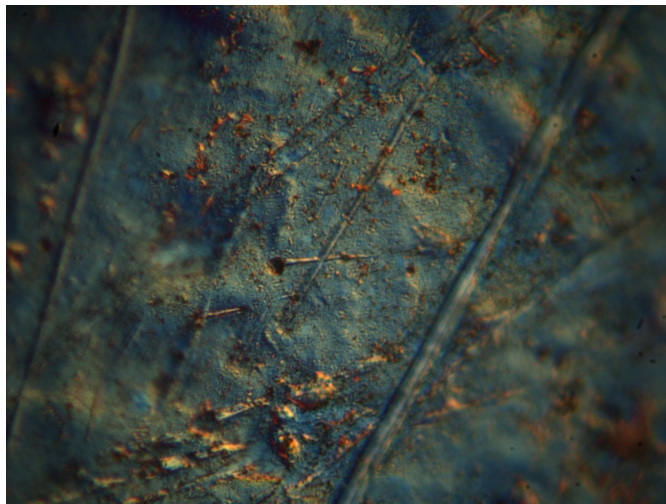
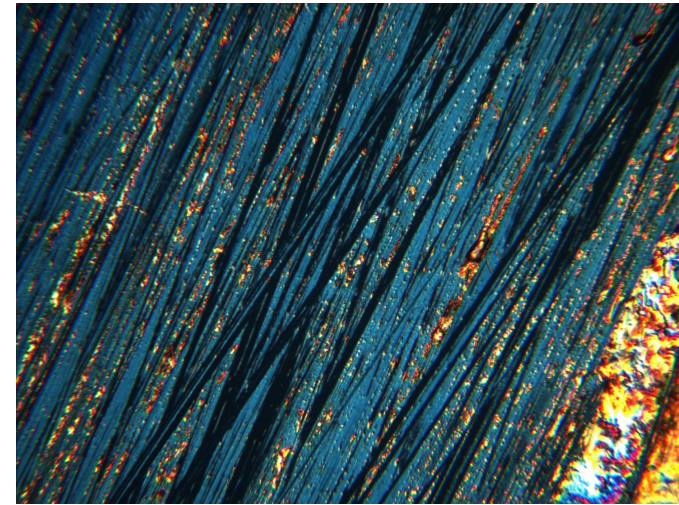
Photos from Metallurgical Microscope (by H.L.Clausing)

Thick side

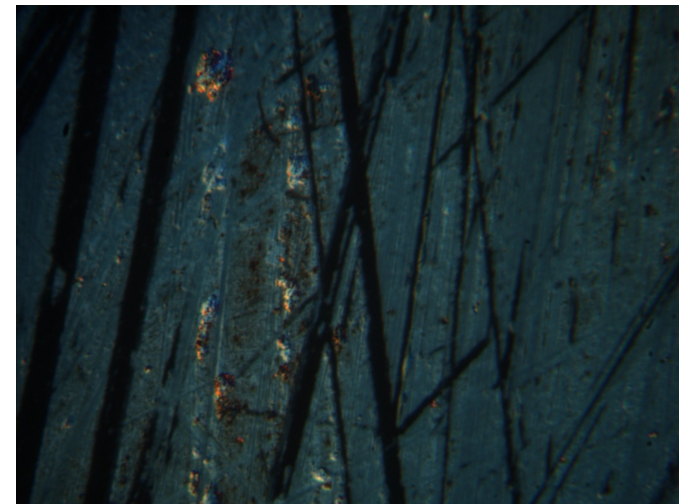


x200

Thin side



x1000



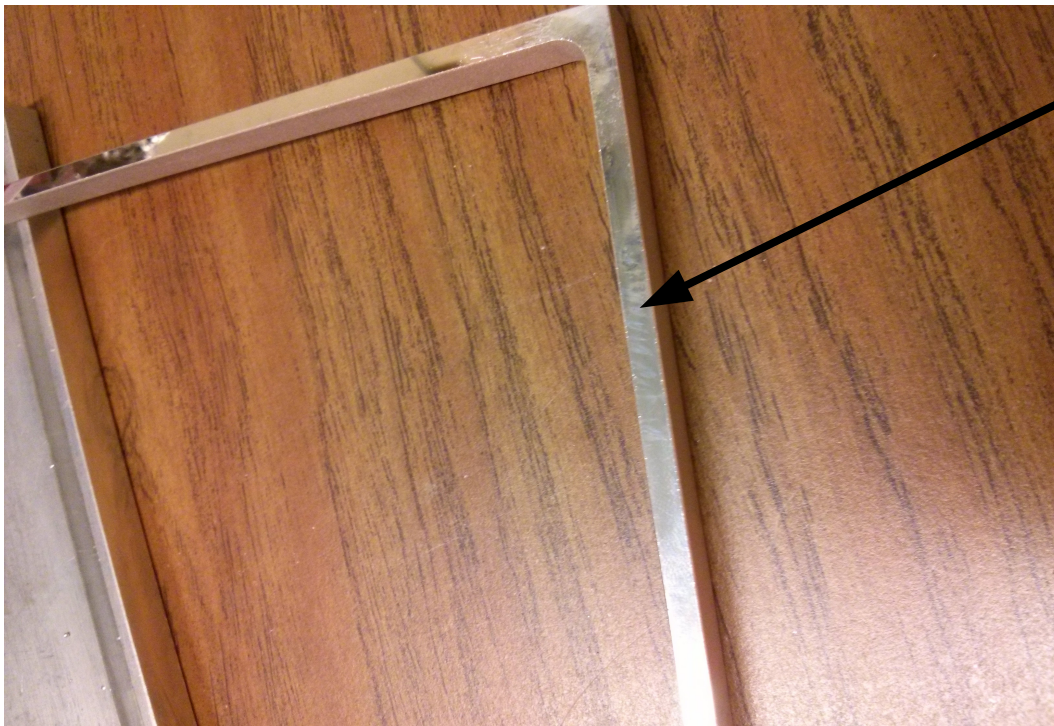
Cu layer dissolves in indium

Solution

**Switched to InBi alloy (melting point 72C)
No signs of Cu scavenging by InBi**

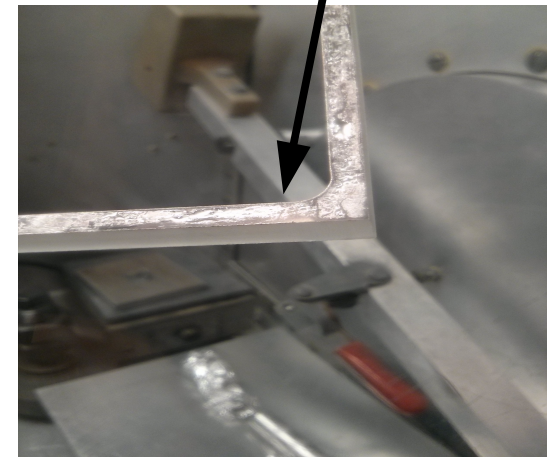
Why pure In works on 1x1" samples?

- *faster assembly (less time at high temperature)*
- *less mechanical rubbing*



InBi wetting test

Compare pure In



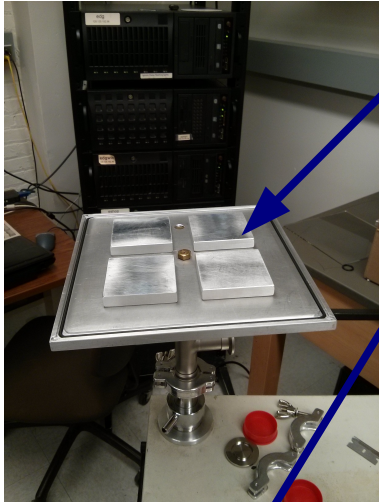
400cm² seal with 2.75 mm window



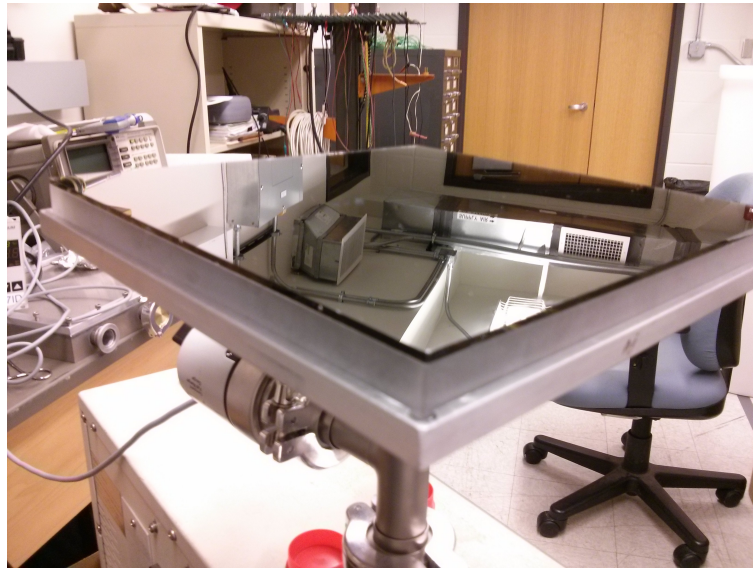
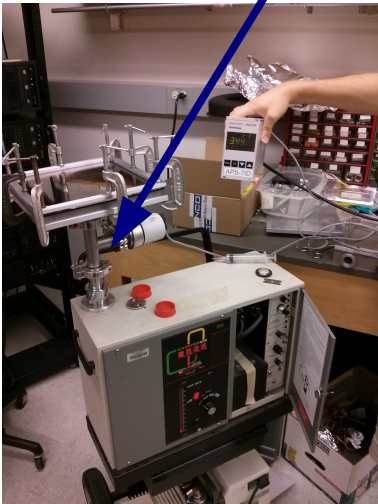
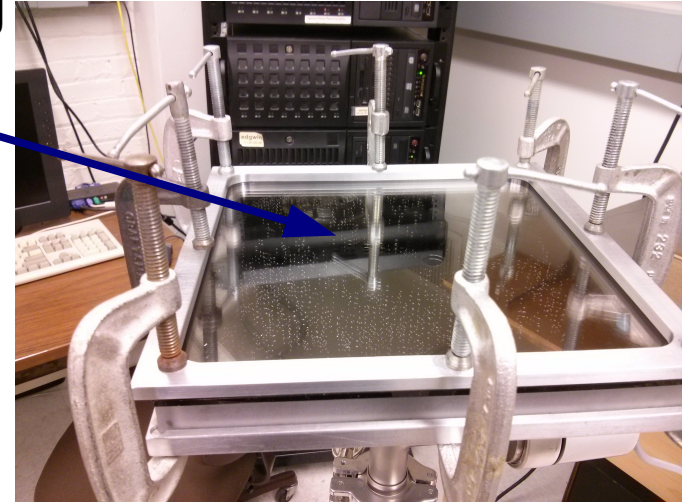
First leak tight 8x8" seal using LAPPD glass parts

Leak Test

Many tricks to avoid glass cracking when the pump starts



- Leak checker sensitivity 10^{-8} cc/s of He
- Leak test lasted for 1 hour
- **No leaks found!**



Next Steps

- **Test reproducibility of the “hot seal” recipe**
- **Move forward with tile assembly**
- **Try to make photo-cathode by In-Situ Photo-Cathode Synthesis** (possible lunch topic)

